

KIC 8462852: Will the Trojans return in 2021?

Fernando J. Ballesteros,^{1*} Pablo Arnalte-Mur,^{1,2}
 Alberto Fernández-Soto,^{3,4} and Vicent J. Martínez^{1,2,4}

¹*Observatori Astronòmic, Universitat de València, C/ Catedrático José Beltrán 2, E46980-Paterna (València), Spain*

²*Departament d'Astronomia i Astrofísica, Universitat de València, E46100-Burjassot (València), Spain*

³*Instituto de Física de Cantabria (CSIC-UC), E39005-Santander, Spain*

⁴*Unidad Asociada Observatorio Astronómico (IFCA-UV), E46980-Valencia, Spain*

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ABSTRACT

KIC 8462852 stood out among more than 100,000 stars in the *Kepler* catalogue because of the strange features of its light curve: a wide, asymmetric dimming taking up to 15 per cent of the light at D793 and a period of multiple, narrow dimmings happening approximately 700 days later. Several models have been proposed to account for this abnormal behaviour, most of which require either unlikely causes or a finely-tuned timing. We aim at offering a relatively natural solution, invoking only phenomena that have been previously observed, although perhaps in larger or more massive versions. We model the system using a large, ringed body whose transit produces the first dimming and a swarm of Trojan objects sharing its orbit that causes the second period of multiple dimmings. The resulting orbital period is $T \approx 12$ years, with a semi-major axis $a \approx 6$ au. Our model allows us to make two straightforward predictions: we expect the passage of a new swarm of Trojans in front of the star starting during the early months of 2021, and a new transit of the main object during the first half of 2023.

Key words: planets and satellites: dynamical evolution and stability – stars: individual: KIC 8462852 – stars: peculiar

1 INTRODUCTION

In September 2015 the discovery of an extraordinary object in the *Kepler* field was announced: the star KIC 8462852¹. Volunteer planet hunters (Fischer et al. 2012) have expressed that its light curve is probably the most bizarre among the more than 100,000 light curves in the *Kepler* field. It is unique because it presents brief, deep drops in flux, with non-periodic repetitions and asymmetric dips, and one particularly complex event around day D1500 that covers up to 20 per cent of the stellar flux². Boyajian et al. (2016) analyzed and discussed possible astrophysical scenarios that could account for it. They discarded instrumental problems or intrinsic variability of the star or its M star dwarf supposed companion as possible causes. The authors considered different scenarios including dust from collisions within objects in a possible asteroid belt, debris from a giant collision similar to the one that supposedly caused the creation of the Earth's Moon, and a swarm of comet fragments orbiting around the star, which could account for the

dips in the light curve. The latter possibility is the one that Boyajian et al. consider most likely, given that the other features are not expected in a main sequence star such as KIC 8462852. Subsequently other works by Schaefer (2016) and Montet & Simon (2016) added separate intriguing secular flux dimmings with timescales of ~ 100 and 4 years, respectively. Note, however, that the century-scale dimming has been disputed by some authors (Hippke et al. 2016, 2017).

In this work we introduce an alternative scenario where stable debris would be expected: the Trojan regions around a giant, ringed object orbiting KIC 8462852. Most of the scenarios that have already been discussed by other authors invoke the presence of astronomical objects that range from uncommon to never directly observed, from the relatively mundane comet clouds in Boyajian et al. (2016) to the alien Dyson sphere in Wright et al. (2016). Our model requires the presence of relatively familiar objects, namely a large planet with orbiting rings and a cloud of Trojan asteroids. Moreover, our model allows us to make a definite prediction: the leading Trojan cloud should induce a new period of irregularities in the light curve approximately in 2021.

* E-mail: fernando.ballesteros@uv.es

¹ Also known in the literature as *Boyajian's star*, *Tabby's star* or *the WTF star*.

² All through this work we will refer to dates using the *Kepler* mission dating system, *i.e.* BJD-2454833.

were obtained from the ‘NASA Exoplanet Archive’³ web service (Basri et al. 2005). This Letter is organized as follows: Section 2 discusses the Trojan hypothesis, we make a brief discussion of the implication of the results in Section 3, and give our conclusions in Section 4.

2 TROJAN HYPOTHESIS

We present our interpretation of the features observed in the light curve of KIC 8462852 as due to the transit of a large orbiting body and its Trojan cohort. The detailed properties of the body that produced the D793 event, not critical to the Trojan hypothesis we introduce in this manuscript, will be presented elsewhere (Ballesteros et al., in prep.). We mention here that the observed shape of the first transit can be reproduced by means of a large body with an extensive ring system, transiting the star at a relatively large impact parameter. The ring system is slightly tilted with respect to the orbital plane, a fact that combined with its impact parameter produces the observed temporal asymmetry.

One of our preferred hypothesis among those considered in Boyajian et al. (2016) to explain the peculiarities in the light curve of KIC 8462852 is that of a giant impact. According to it the second event at D1500 would be produced by the same material observed at D793, dispersed by a large impact and seen at a time corresponding to the following orbit (from where an orbital period $T \sim 700$ days would be deduced). Nevertheless, Boyajian and collaborators objected to this model based on the low probability of witnessing such an event and the non-repetition of the dips that appeared early in the Kepler mission coverage, and they preferred the hypothesis of a group of objects in a highly eccentric orbit around the star. We offer here an alternative scenario which does not require fine-tuned time dependence, so our witnessing it does not render it improbable—it is a recurrent event. It may represent an extreme scenario, but this is to be expected given the *a priori* fact that we are trying to explain a light curve selected because of its extreme rarity.

We interpret the cluttered behaviour of the second epoch, observed from \sim D1500, as caused by a cohort of objects close to the L5 Lagrangian point associated to the main orbiting body. Such a stable and large swarm of bodies, debris and dust, gravitationally confined around the L5 region, can explain several of the observed features. We remark that the presence of relatively large amounts of dust is not excluded; if dust is farther away than 0.2–0.3 au from the star it wouldn’t be readily detectable by WISE or *Spitzer* (see Boyajian et al. 2016; Marengo et al. 2015).

In our Solar System several worlds have gathered bodies at their Trojan regions. It is the case of the Earth, Mars, Neptune and especially Jupiter, but also of moons like Tethys and Dione; all of them sharing their orbits with objects in their Trojan points. Hydrodynamic simulations of protoplanetary disks (Laughlin & Chambers 2002) show that dust and disk material linger in the Trojan stable regions of a planet, remaining there after the planetary formation process is complete. Orbital instabilities such as proposed by the Nice model (Gomes et al. 2005; Tsiganis et al.

2005; Morbidelli et al. 2005) could cause large numbers of small bodies to be in unusually excited orbits, which could create both Trojans and rings. Additionally, violent events in the past could also have led to the capture of objects in these regions (Morbidelli et al. 2005). Thus, stable Trojan bodies related to exoplanets should be commonplace. In fact, Hippke & Angerhausen (2015) already found a hint of the presence of Trojans in long-period exoplanetary orbits using *Kepler* data.

In order to estimate the extension of the Trojan swarm along the orbit we can examine the case of Jupiter. Jewitt et al. (2000) measured an apparent FWHM of the L4 swarm along the ecliptic of $26.4^\circ \pm 2.1^\circ$, corresponding to a linear size of 2.4 au. Taking twice this value to encompass the whole swarm, we can estimate that Jupiter’s L4 Trojan cloud roughly covers an angular extent of $\sim 50^\circ$ along its orbit, which is also in accordance with Karlsson (2010). Regarding the strange structure of the second transit epoch, we remark that studies of stability for planets more massive than Jupiter with elliptic orbits (Erdi et al. 2007) show that the spatial distribution of the stability regions around L4 and L5 could have a very complex structure.

In the Solar System we observe that the ensemble mass of the Jupiter Trojans is $\sim 0.0001M_\oplus$, with a total cross section similar to a disk of radius ~ 2000 km (i.e. larger than that of the Moon), but a mass which is two orders of magnitude smaller (Jewitt et al. 2000; Sheppard & Trujillo 2010). Even though in our Solar System the Trojan-to-planet mass ratio is low ($\leq 10^{-8}$), extrasolar planets may have more massive Trojans, perhaps even with ratios as large as one-to-one (Ford & Gaudi 2006).

Taking as an example a hypothetical amount of mass $M = M_{\text{Jup}}$ trapped in the Trojan regions of KIC 8462852 following the same size distribution as the Jupiter Trojans (Wong et al. 2014; Fernández et al. 2003; Harris & Harris 1997), we reach an effective cross section as high as ~ 200 times the star cross section when considering bodies with diameter larger than 1 km. Allowing for Trojan bodies as small as 50 meters the figure can rise up to 500 times the cross section of the star. Nevertheless, large amounts of dust can remarkably relax any mass requirement: dust accumulated in the Trojan region can be a major factor for opacity, in fact dominating the cross section. Boyajian et al. (2016) estimate the mass of dust necessary to produce the observed opacity to be 6.7×10^{18} g. As we will discuss in the next Section, using this same model we obtain that a cloud mass well below $10^{-4} M_\oplus$ within the Trojan regions would be enough to produce the observed results. Collisions among Trojans or catastrophic events in the past could have generated or trapped large enough amounts of dust in these regions.

One problem remains: the largest individual dip observed close to D1500, which shows substructure in the light curve, would correspond to a large event which covered a significant fraction of the stellar cross section. The presence of substructure points to the possibility that we may be observing a clustering of smaller bodies and dust, probably gravitationally linked. This is indeed one of the most difficult aspects that any model of this fascinating observation ought to address, but, as in Boyajian’s cometary scenario, clumps seem natural for the Trojan model and the existence of collisional families (that can produce an additional amount of dust, thus increasing the cross section) should be

³ <http://exoplanetarchive.ipac.caltech.edu/>

expected. In our case we need to postulate the presence of one such cluster close to the Lagrange L5 point of the main planet. As mentioned above, some studies point that this may not be particularly strange (Ford & Gaudi 2006).

All in all, according to the hypothesis that the features observed at \sim D1500 correspond to the passage of Trojans close to the trailing L5 point of the transiting planet, one should expect a similar signature in the symmetric point before the planetary transit. Unfortunately *Kepler* observations of KIC 8462852 started at D120, just after the epoch when the putative symmetric L4 Trojans would have transited (ending around D65). Notice, however, that at the beginning of the time series there are some small features that could be produced by the last Trojans inhabiting the L4 zone. In fact, reversing the time series around D793 (see Figure 1, top), a certain symmetry with the original time series appears. Beyond a distance of \sim 300 days before and after this date the variability of the time series seems to increase, while it remains quieter in the region in between. Calculation of the time series standard deviation in weekly periods seems to confirm this symmetry around the presumed planet, strengthening the Trojan hypothesis. We have confirmed that during the same period other stars in the nearby region of the CCD where KIC 8462852 was detected did not show this kind of fluctuations, ruling out the hypothesis of variations in the CCD sensitivity with time, changes in the illumination on the focal plane, or other instrumental effects that should have affected all targets approximately in the same manner. One thing we cannot exclude, however, is that some of the minor dips that we are assuming could be caused by either trailing or leading Trojans could in reality be caused by dimming episodes in nearby, blended objects, as discussed by Makarov & Goldin (2016). We remark, though, that because of their aperiodic nature we would be somehow transferring the problem of explaining their nature from KIC 8462852 to the other objects.

3 IMPLICATIONS OF THE MODEL

If the Trojan hypothesis is correct we can obtain a direct estimate of the orbital period. The time interval between the first and second main events in the *Kepler* data is approximately two years. Assuming an orbital separation of $\sim 60^\circ$ between them, the orbital period would be ~ 12 years, which given the mass estimated for the star ($1.43 M_\odot$, Boyajian et al. 2016) would imply a semi-major axis of 5.9 au. With these parameters and the stellar radius $R = 1.58 R_\odot$ estimated by Boyajian et al. (2016), an object in a circular orbit would move at 15 km s^{-1} and an equatorial transit would last ~ 1.7 days. Considering a ~ 20 per cent uncertainty in the mass and radius of the star as Boyajian et al. (2016) quote in their paper, and the fact that an elliptical orbit would induce an extra uncertainty in the transit speed⁴, the transit speed and duration could accommodate changes of up to a factor of 2. These figures move the expected duration of the main planetary transit closer to the one observed,

although tensions persist: Boyajian et al. (2016) estimate a velocity for the transiting material of up to 50 km s^{-1} , and a particular transiting event close to D1568 had a duration as short as 0.4 days. The above-mentioned uncertainties together with the effect of a large impact parameter and the possible libration of the Trojans around their equilibrium points could diminish the tension between both values.

If we consider that in our model the hypothetical Trojan cloud started at the epoch when the light curve of KIC 8462852 began to increase its variability (shaded region in Fig. 1), the duration until the end of the light curve is ~ 500 days, which would represent an angular extent of $\sim 40^\circ$ for the assumed orbital period of ~ 12 years. Given that the irregular behaviour very probably extended beyond the end of the observed light curve, the extension could perfectly reach a value comparable to $\sim 50^\circ$, the orbital extension of the Trojan swarm around Jupiter's L4 point presented in the previous Section.

We have used a very simple model, a bi-Gaussian distribution of material around the L5 point with a FWHM extension of 26.4 degrees along the orbit and the same spread in the vertical direction, to estimate that only $\approx 1/350$ of the cloud material will cross in front of the projected area of the star. This implies the amount of dust and material in the whole proposed Trojan region could be ~ 350 times bigger than that producing the transit, or 700 times bigger considering both Lagrange regions. As the infrared luminosity limits from WISE and Spitzer presented in Boyajian et al. (2016, see their Figure 12) only refer to dust passing directly in front of the star, the luminosity estimate at 6 au should then be increased by a factor 700. This is still well below the observational WISE limits, that could accommodate a factor up to 2000. Moreover the dust model in Boyajian et al. (2016) rescaled to 6 au and distributed according to the former bi-Gaussian model yields an mass estimate $M = 4 \times 10^{23} \text{ g}$ if we allow particle sizes up to 1 cm—a mass of dust well below $10^{-4} M_\oplus$ within both Trojan regions.

Boyajian et al. (2016) performed an estimate of the range of planetary masses and orbital periods that could be compatible with the observed absence of radial velocity variations in the available spectroscopic observations. We remark that these are scarce (four runs), of not very high precision ($\sigma_v \sim 400 \text{ m s}^{-1}$), and cover a relatively small period of time (less than 1.5 years). We have performed a similar calculation, fixing the orbital period at $T = 12$ years and assuming that the D793 transit happened close to the periastron. Under this assumptions, and allowing for orbital ellipticities between 0 and 0.6, we derive a 1σ upper limit to the mass of the planetary object $M_p < 130 - 170 M_{\text{Jup}}$, with the highest limit corresponding to circular orbits. This limit does hardly constrain the model at all. In fact, just requiring the stability of the Trojans near the Lagrangian points imposes a ratio $M_{\text{star}}/M_p > 25$ which, given the estimated mass of the star, translates into $M_p \lesssim 60 M_{\text{Jup}}$, a more restrictive limit. New determinations of the radial velocity, even if they were of similar precision, would greatly reduce the uncertainty which is mostly caused by the degeneracy between the systemic radial velocity and the maximum Δv induced by the orbiting body. A stringent limit on the mass of the orbiting body would not be critical for the Trojan model here presented, as our estimates depend on apparent sizes derived from observed opacities, large values of which

⁴ A large orbital eccentricity would also introduce instability in the Trojan regions. Previous works about this indicate that an eccentricity as high as 0.3 (Chanut et al. 2004; Robutel & Gabern 2006) or even 0.6 (Erdi et al. 2007) could be allowed for.

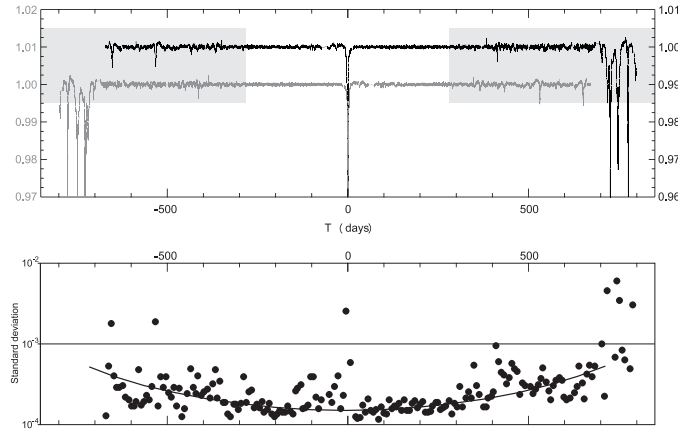


Figure 1. Top: Light curve of KIC 8462852, inverted around D793 (gray) and superimposed on the original (black). The time-inverted curve is shifted downwards for clarity. There is a quiescent period lasting ≈ 600 days centered around D793. Outside this period we detect an increase in the variability of the light curve (shaded regions). Bottom: Standard deviation of the data. Each point represents the standard deviation of the light intensities during a one-week period, and seems to increase symmetrically as one moves away from D793. The solid line is a guide to the eye and does not represent a fit to the data. Points with large deviations correspond to light curve peaks.

can be produced by relatively small masses of dust and/or small bodies, and a probable highly opaque ring system.

3.1 The dimming event at D3060

On May 19, 2017 an alert was launched by T. Boyajian reporting the observation of a possible low-intensity dip in the lightcurve of KIC 8462852 (Boyajian 2017; Boyajian et al. 2017). Observers interpreted this dip as the possible beginning of a new passage of the debris associated to the catastrophic event that may have happened to the large planetary object that transited at D793, causing it to be observed as a swarm at D1500, confirming in such case an orbital period of approximately two years.

Within our model the timing of this event would correspond approximately to the opposition of the main body, therefore the corresponding L3 point would be passing in front of the star by this time. In this scenario, this dimming event could be explained by the effect of objects akin to the Hildian asteroids in the Solar System passing the L3 area⁵. If what we have identified as Trojan asteroid regions are really so densely populated, in principle one could expect also a rather high density of Hildas, particularly if some catastrophic event in the (relatively) near past has populated both orbits. In this case, we would expect a few dimming events concentrated around the time when L3 transits in front of the star, as well as the possibility that some of the events that we have generally classified as Trojans could in reality be due to Hildas. This hypothesis could help reconcile some of the fastest ones, as the Hilda orbits have a shorter period. The concentration of Hilda events close to L3 would, in any case, be much smaller than the complex event at D1500 caused by the passage of the Trojans.

⁵ The Hildas are a dynamic group of Solar System asteroids that occupy a 3:2 orbital resonance with Jupiter. The resonance makes their aphelia alternatively close to Jupiter’s Lagrange points L3, L4, and L5; in such a way that at any time they are preferentially found close to those three points.

4 CONCLUSIONS

We have presented an alternative model to explain the odd appearance of the light curve of KIC 8462852. We propose that a grazing transit of a large, ringed object could produce the asymmetric transit observed at D793, whereas a huge swarm of Trojan objects inhabiting its L5 orbital point could have caused the irregular transit at D1500. We deduce an orbital period $T \approx 12$ years. In principle this would imply a transit speed slower than observed, although a highly eccentric orbit would make them close to compatible. Full details of the modeling of the main planetary transit at D793 will be presented elsewhere. We estimate its mass to be $\lesssim 150 M_{\text{Jup}}$ (stellar radial velocity) and $\lesssim 60 M_{\text{Jup}}$ (Trojan cloud dynamic stability), but it can be much smaller as the observed large cross section does not necessarily imply a large mass. We show in Figure 2 a diagram representing the main parameters and properties of our model, together with an idealized vision of the light curve. Note that our model does not explain the secular dimmings observed by Schaefer (2016) and Montet & Simon (2016). Such effects would call for a completely unrelated explanation.

Given the exceptional behaviour of this light curve, our explanation is also somehow exceptional –as are all the other proposed hypotheses– but not too unconventional. It is nourished by the evidence of similar (although obviously not identical) existing objects in our Solar System and beyond, which have been previously studied in detail. A key advantage of our model is the necessary repeatability of the phenomenon. This fact puts aside any coincidence or temporal fine-tuning quandary, as our observation does not imply a particular moment in the history of the observed system.

This repeatability also allows us to carry out a testable prediction: under the assumptions of our hypothesis, considering an orbital period of ~ 12 years, and taking into account that the region of deepest dimming lasted three months, we predict the onset of a new epoch of irregular transits at \sim D4430, i.e., February 2021. In other words, during the early months of 2021 the swarm of objects at the symmetric L4 Lagrangian point will transit the star, starting an epoch of

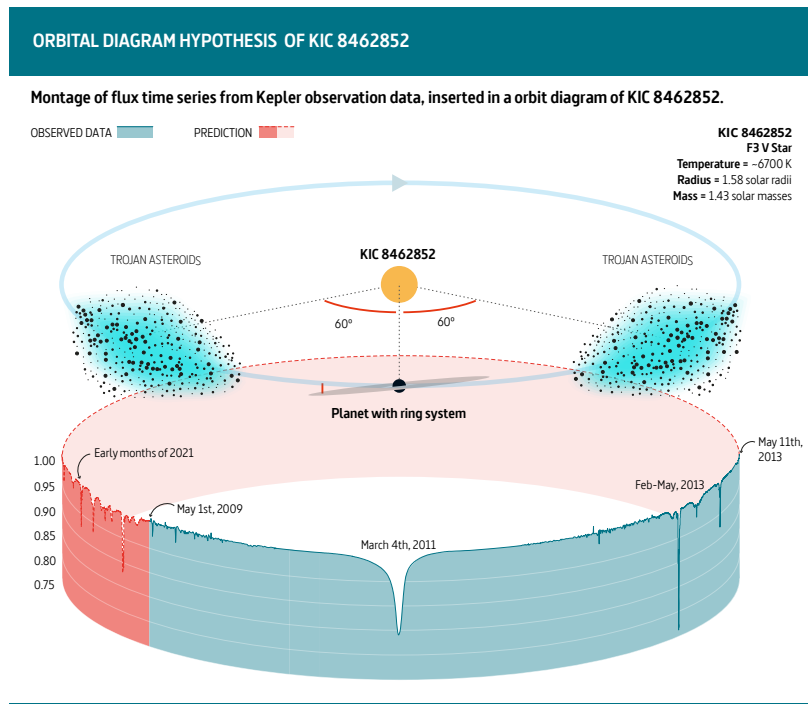


Figure 2. Diagram showing a hypothetical ringed giant body orbiting the star, together with its dense populations of Trojan bodies and dust around the L4 and L5 Lagrange points. We also present below it the observed (blue) and expected (red) light curve.

cluttered dimmings. Two years later, during the first half of 2023, we expect a new transit of the ringed planetary body.

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References

- Basri G., Borucki W. J., Koch D., 2005, *New Astron. Rev.*, **49**, 478
- Boyajian T. (@tsboyajian), 2017, "#TabbysStar IS DIPPING! OBSERVE!!". 19 May 2017, 4:32 AM. Tweet
- Boyajian T. S., et al., 2016, *MNRAS*, **457**, 3988
- Boyajian T., et al., 2017, The Astronomer's Telegram, 10405
- Chanut T. G. G., Tsuchida M., Winter O. C., 2004, in *Advances in space dynamics 4: Celestial mechanics and astronautics*. pp 57–63
- Erdi B., Frohlich G., Nagy I., Zs. S., 2007, in *Proc. of the 4th Austrian Hungarian Workshop on celestial mechanics*. p. 85, doi:10.1111/j.1749-6632.1980.tb15927.x
- Fernández Y. R., Sheppard S. S., Jewitt D. C., 2003, *AJ*, **126**, 1563
- Fischer D. A., et al., 2012, *MNRAS*, **419**, 2900
- Ford E. B., Gaudi B. S., 2006, *ApJ*, **652**, L137
- Gomes R., Levison H. F., Tsiganis K., Morbidelli A., 2005, *Nature*, **435**, 466
- Harris A. W., Harris A. W., 1997, *Icarus*, **126**, 450
- Hippke M., Angerhausen D., 2015, *ApJ*, **811**, 1
- Hippke M., Angerhausen D., Lund M. B., Pepper J., Stassun K. G., 2016, *ApJ*, **825**, 73
- Hippke M., et al., 2017, *ApJ*, **837**, 85
- Jewitt D. C., Trujillo C. A., Luu J. X., 2000, *AJ*, **120**, 1140
- Karlsso O., 2010, *A&A*, **516**, A22
- Laughlin G., Chambers J. E., 2002, *AJ*, **124**, 592
- Makarov V. V., Goldin A., 2016, *ApJ*, **833**, 78
- Marengo M., Hulsebus A., Willis S., 2015, *ApJ*, **814**, L15
- Montet B. T., Simon J. D., 2016, *ApJ*, **830**, L39
- Morbidelli A., Levison H. F., Tsiganis K., Gomes R., 2005, *Nature*, **435**, 462
- Robutel P., Gabern F., 2006, *MNRAS*, **372**, 1463
- Schaefer B. E., 2016, *ApJ*, **822**, L34
- Sheppard S. S., Trujillo C. A., 2010, *ApJ*, **723**, L233
- Tsiganis K., Gomes R., Morbidelli A., Levison H. F., 2005, *Nature*, **435**, 459
- Wong I., Brown M. E., Emery J. P., 2014, *AJ*, **148**, 112
- Wright J. T., Cartier K. M. S., Zhao M., Jontof-Hutter D., Ford E. B., 2016, *ApJ*, **816**, 17