Topological Strings and Donaldson-Thomas invariants

Michele Cirafici

University of Patras $\Pi \alpha \nu \epsilon \pi \iota \sigma \tau \dot{\eta} \mu \iota \sigma \Pi \alpha \tau \rho \dot{\omega} \nu$

RTN07 Valencia - Short Presentation

work in progress with A. Sinkovics and R.J. Szabo

- Topological Strings play a very important role in modern mathematical physics
- They compute F—terms in supersymmetric theories and capture microscopic properties of Black Holes.
- Mathematically they count enumerative invariants of Calabi-Yau geometries such as the Gromov–Witten, the Gopakumar–Vafa and the recently introduced Donaldson–Thomas.
- Roughly, the physical interpretation of the Donaldson–Thomas invariant $D_{n,\beta}$ is that it gives the number of bound states of n D0 branes and D2 branes wrapping the cycle β with a single D6 brane wrapping the full Calabi–Yau
- They are the underlying objects in the "melting crystal" formulation of topological strings



- Topological Strings play a very important role in modern mathematical physics
- They compute F-terms in supersymmetric theories and capture microscopic properties of Black Holes.
- Mathematically they count enumerative invariants of Calabi-Yau geometries such as the Gromov-Witten, the Gopakumar-Vafa and the recently introduced Donaldson-Thomas.
- Roughly, the physical interpretation of the Donaldson–Thomas invariant $D_{n,\beta}$ is that it gives the number of bound states of n D0 branes and D2 branes wrapping the cycle β with a single D6 brane wrapping the full Calabi–Yau
- They are the underlying objects in the "melting crystal" formulation of topological strings

- Topological Strings play a very important role in modern mathematical physics
- They compute F-terms in supersymmetric theories and capture microscopic properties of Black Holes.
- Mathematically they count enumerative invariants of Calabi-Yau geometries such as the Gromov–Witten, the Gopakumar–Vafa and the recently introduced Donaldson–Thomas.
- Roughly, the physical interpretation of the Donaldson–Thomas invariant $D_{n,\beta}$ is that it gives the number of bound states of n D0 branes and D2 branes wrapping the cycle β with a single D6 brane wrapping the full Calabi–Yau
- They are the underlying objects in the "melting crystal" formulation of topological strings

- Topological Strings play a very important role in modern mathematical physics
- They compute F-terms in supersymmetric theories and capture microscopic properties of Black Holes.
- Mathematically they count enumerative invariants of Calabi-Yau geometries such as the Gromov–Witten, the Gopakumar–Vafa and the recently introduced Donaldson–Thomas.
- Roughly, the physical interpretation of the Donaldson–Thomas invariant $D_{n,\beta}$ is that it gives the number of bound states of n D0 branes and D2 branes wrapping the cycle β with a single D6 brane wrapping the full Calabi–Yau
- They are the underlying objects in the "melting crystal" formulation of topological strings

- Topological Strings play a very important role in modern mathematical physics
- They compute F-terms in supersymmetric theories and capture microscopic properties of Black Holes.
- Mathematically they count enumerative invariants of Calabi-Yau geometries such as the Gromov–Witten, the Gopakumar–Vafa and the recently introduced Donaldson–Thomas.
- Roughly, the physical interpretation of the Donaldson–Thomas invariant $D_{n,\beta}$ is that it gives the number of bound states of n D0 branes and D2 branes wrapping the cycle β with a single D6 brane wrapping the full Calabi–Yau
- They are the underlying objects in the "melting crystal" formulation of topological strings

- We will focus on local toric Calabi-Yaus
- In this framework the DT problem can be rephrased in terms of an auxiliary topological Yang-Mills theory that lives on the D6 branes worldvolume
- DT invariants are "generalized instantons" of this theory an can be computed by the techniques of equivariant cohomology and localization used by Nekrasov in the context of Seiberg—Witten theory
- They are classified by 3d Young tableaux: the generalized instantons reproduce the states of the melting crystal

- We will focus on local toric Calabi-Yaus
- In this framework the DT problem can be rephrased in terms of an auxiliary topological Yang-Mills theory that lives on the D6 branes worldvolume
- DT invariants are "generalized instantons" of this theory an can be computed by the techniques of equivariant cohomology and localization used by Nekrasov in the context of Seiberg-Witten theory
- They are classified by 3d Young tableaux: the generalized instantons reproduce the states of the melting crystal

- We will focus on local toric Calabi-Yaus
- In this framework the DT problem can be rephrased in terms of an auxiliary topological Yang-Mills theory that lives on the D6 branes worldvolume
- DT invariants are "generalized instantons" of this theory an can be computed by the techniques of equivariant cohomology and localization used by Nekrasov in the context of Seiberg-Witten theory
- They are classified by 3d Young tableaux: the generalized instantons reproduce the states of the melting crystal

- We will focus on local toric Calabi-Yaus
- In this framework the DT problem can be rephrased in terms of an auxiliary topological Yang-Mills theory that lives on the D6 branes worldvolume
- DT invariants are "generalized instantons" of this theory an can be computed by the techniques of equivariant cohomology and localization used by Nekrasov in the context of Seiberg-Witten theory
- They are classified by 3d Young tableaux: the generalized instantons reproduce the states of the melting crystal

- This is a generalization of the DT invariants with an arbitrary number of D6 branes. We are using two independent techniques to compute the non abelian invariants on any toric geometry.
- A noncommutative deformation of the theory resolves the singularities of the moduli space. The invariants can be computed by direct localization of the path integral.
- We can introduce an auxiliary topological matrix quantum mechanics directly on the moduli space. This reformulates the problem in an "ADHM-like" fashion.
- But many physical open problems and potential developments further: OSV conjecture, S-duality in topological strings, mirror symmetry, wallcrossings...



- This is a generalization of the DT invariants with an arbitrary number of D6 branes. We are using two independent techniques to compute the non abelian invariants on any toric geometry.
- A noncommutative deformation of the theory resolves the singularities of the moduli space. The invariants can be computed by direct localization of the path integral.
- We can introduce an auxiliary topological matrix quantum mechanics directly on the moduli space. This reformulates the problem in an "ADHM-like" fashion.
- But many physical open problems and potential developments further: OSV conjecture, S-duality in topological strings, mirror symmetry, wallcrossings...



- This is a generalization of the DT invariants with an arbitrary number of D6 branes. We are using two independent techniques to compute the non abelian invariants on any toric geometry.
- A noncommutative deformation of the theory resolves the singularities of the moduli space. The invariants can be computed by direct localization of the path integral.
- We can introduce an auxiliary topological matrix quantum mechanics directly on the moduli space. This reformulates the problem in an "ADHM-like" fashion.
- But many physical open problems and potential developments further: OSV conjecture, S-duality in topological strings, mirror symmetry, wallcrossings...



- This is a generalization of the DT invariants with an arbitrary number of D6 branes. We are using two independent techniques to compute the non abelian invariants on any toric geometry.
- A noncommutative deformation of the theory resolves the singularities of the moduli space. The invariants can be computed by direct localization of the path integral.
- We can introduce an auxiliary topological matrix quantum mechanics directly on the moduli space. This reformulates the problem in an "ADHM-like" fashion.
- But many physical open problems and potential developments further: OSV conjecture, S-duality in topological strings, mirror symmetry, wallcrossings...



- In the recent years it has been proven possible to introduce Gromov–Witten invariants for orbifolds
- What about DT invariants? Open problem: we choose a pragmatic approach and try to compute them directly as an instanton problem.
- We define the invariants in terms of "colored" partition: each box of the 3D Young tableau has a different color according to the transformation properties of the instanton under the orbifold action
- Comparison with orbifold Gromov—Witten invariants is technically challenging: not yet understood!

- In the recent years it has been proven possible to introduce Gromov–Witten invariants for orbifolds
- What about DT invariants? Open problem: we choose a pragmatic approach and try to compute them directly as an instanton problem.
- We define the invariants in terms of "colored" partition: each box of the 3D Young tableau has a different color according to the transformation properties of the instanton under the orbifold action
- Comparison with orbifold Gromov—Witten invariants is technically challenging: not yet understood!

- In the recent years it has been proven possible to introduce Gromov–Witten invariants for orbifolds
- What about DT invariants? Open problem: we choose a pragmatic approach and try to compute them directly as an instanton problem.
- We define the invariants in terms of "colored" partition: each box of the 3D Young tableau has a different color according to the transformation properties of the instanton under the orbifold action
- Comparison with orbifold Gromov—Witten invariants is technically challenging: not yet understood!

- In the recent years it has been proven possible to introduce Gromov–Witten invariants for orbifolds
- What about DT invariants? Open problem: we choose a pragmatic approach and try to compute them directly as an instanton problem.
- We define the invariants in terms of "colored" partition: each box of the 3D Young tableau has a different color according to the transformation properties of the instanton under the orbifold action
- Comparison with orbifold Gromov–Witten invariants is technically challenging: not yet understood!