

N=1 4d Supermembrane from 11D

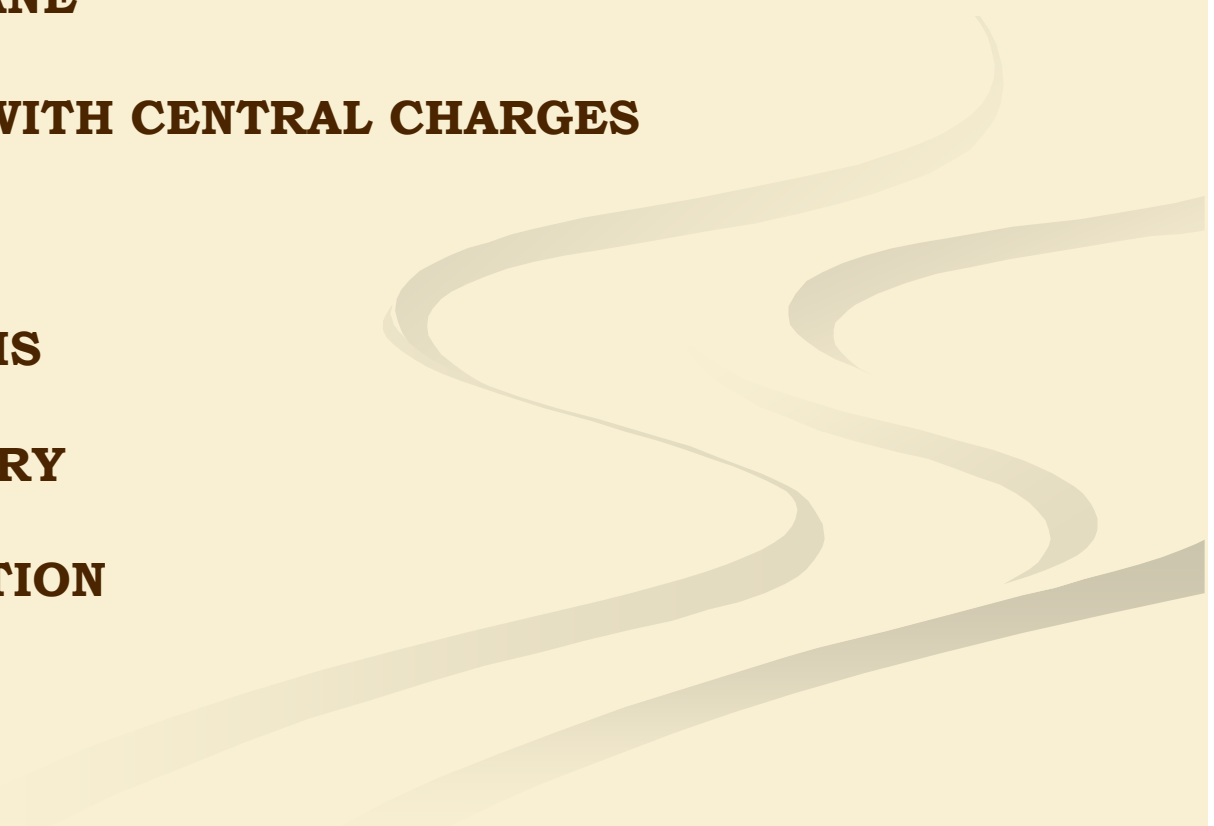
M.P.GARCIA DEL MORAL, J.M. PENA, A. RESTUCCIA

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U. of Torino & INFN, Univ. Central de Venezuela, Potsdam Institute & USB

OUTLINE OF THE TALK

- **MOTIVATION**
 - **11D SUPERMEMBRANE**
 - **SUPERMEMBRANE WITH CENTRAL CHARGES**
 - **4D ACTION**
 - **SPECTRUM ANALYSIS**
 - **N=1 SUPERSYMMETRY**
 - **MODULI STABILIZATION**
- 

1. MOTIVATION

•**GOAL:** Perform the quantization of M2 in 4d with $N=1$

No Supergravity approach. Top- Down Approach.

• **OPEN PROBLEM:** NONPERTURBATIVE QUANTIZATION OF STRING TH.



M–theory Quantization: M2, M5

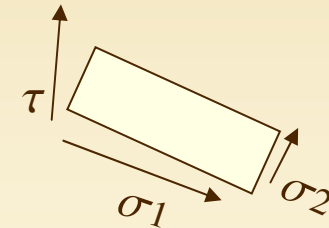
Bergshoeff, Sezgin, Townsend 87

•**TOOLS:** Semiclassical Analysis: *Duff, Inami, Pope, Sezgin, Stelle 88; Restuccia et al. 96*
Matrix Formulation: *de Witt, Hoppe, Nicolai 88'; ..., GM, Restuccia 00' + ...*
Exact Formulation *Boulton, G.M. Restuccia 06'*

2. 11D SUPERMEMBRANE IN THE L.C.G

Bergshoeff, Sezgin, Townsend; De Witt, Hoppe, Nicolai; De Witt, Luscher, Nicolai, De Witt, Marquard, Nicolai, De Witt, Peeters, Plefka. 87-96

2+1D CLOSED RIEMANN SURFACE IN M_{11}



$$\mathbf{H} = \int_{\Sigma} \sqrt{W} \left(\frac{1}{2} \left(\frac{P_M}{\sqrt{W}} \right)^2 + \frac{1}{4} \{X^M, X^N\}^2 + \text{Fermionic terms} \right)$$

with

$$\{X^M, X^N\} = \frac{\varepsilon^{ab}}{\sqrt{W(\sigma)}} \partial_a X^M \partial_b X^N$$

Subject to the following constraints:

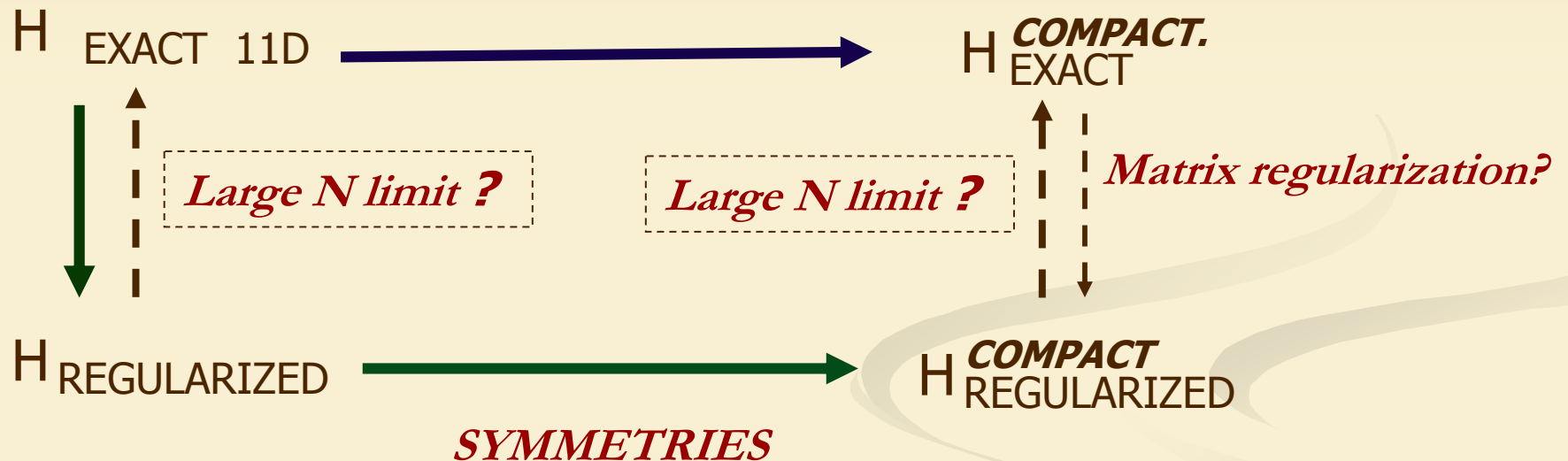
$$\phi_1 := d \left(\frac{P_M}{\sqrt{W}} dX^M \right) = 0$$

$$\phi_2 := \oint_{C_s} \frac{P_M}{\sqrt{W}} dX^M = 0,$$

MATRIX MODELS

Original Point of View: *Halpern; Hoppe, De Wit+Hoppe+Nicolai*

Integrate the spatial dependence of the original action and obtain a susy Quantum mechanics of matrices.



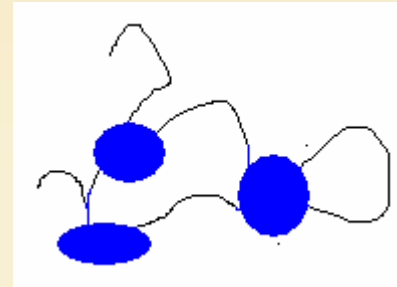
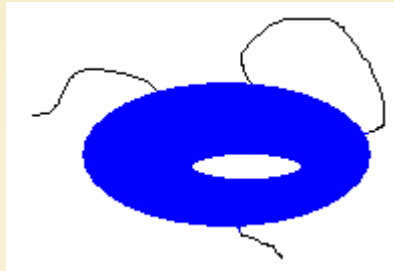
Present Point of View

BFSS/IKKT CONJECTURE: D0 or D-1 action is considered **FUNDAMENTAL**
Impose **SYMMETRIES:** BMN MODEL ETC..

11D Quantization Supermembrane

Historical Approach

CLASSICALY IS UNSTABLE. Topology nor number of Particles is preserved



String-like Spikes

•QUANTUM ANALYSIS : MATRIX MODELS

De Wit, Marquard, Nicolai

Bosonic sector: Purely **DISCRETE**

Luscher 86; G.M, Navarro, Perez, Restuccia; 06

SUSY Sector: **CONTINUOUS**

De Wit+Luscher+Nicolai

•CONCLUSIONS: 2nd QUANTIZED THEORY

Compactifications generically do not change behaviour *De Wit, Peeters, Plefka 96*

Picture of interacting D0's or Strings. Former interpretation as fundamental object seemed to be discarded

4. THE SUPERMEMBRANE WITH CENTRAL CHARGES

*Martin, Restuccia, Torrealba; Martin, Ovalle, Restuccia;MPGM, Restuccia(1); Bellorin, Restuccia, 96-07
Boulton,MPGM., Restuccia(3), Boulton,MPGM., Martin, Restuccia, Boulton(2),MPGM+R,*

$$M_9 \times S^1 \times S^1 \quad \& \quad \text{TOPOLOGICAL CONDITION}$$

The 11D hamiltonian of M2 in the L.C.G. Indices (+ , - , M, N =1,..,9)

$$\mathbf{H} = \int_{\Sigma} \sqrt{W} \left(\frac{1}{2} \left(\frac{P_M}{\sqrt{W}} \right)^2 + \frac{1}{4} \{X^M, X^N\}^2 + \text{Fermionic terms} \right)$$

Subject to the previous constrains + compactification conditions

$$\oint_{C_i} dX^r = 2\pi S_i^r R^r, \quad r = 1, 2,$$

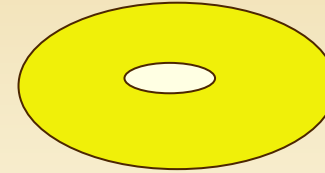
$$\oint_{C_i} dx^m = 0 \quad m = 3, \dots, 9$$

CENTRAL CHARGE CONDITION:

$$Z_{rs} = \int_{\Sigma} dX^r \wedge dX^s = \epsilon^{rs} (2\pi^2 R_1 R_2) n,$$

QUALITATIVE PROPERTIES

SU(N) Spectrum:



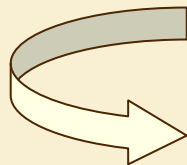
CLASSICALLY: **NOT** STRING-LIKE SPIKES (1)

QUANTUM: BOSONIC \longrightarrow PURELY DISCRETE SPECTRUM (2)

FERMIONIC \longrightarrow PURELY DISCRETE SPECTRUM (3)

BOSONIC EXACT LEVEL *Boulton, MPGM, Restuccia 06*

Infinite number of Configuration Space fields: Canonical quantization Scheme



Pure discreteness of Spectrum of the M2 with central Charges in the 11D algebra.

1 Step: Compactifying to 5D

Bellorin+Restuccia

We perform the following decomposition on the multivalued fields

$$dX^r = 2\pi R^r l^r d\hat{X}^r + \delta_s^r dA_s. \quad \text{MPGM+Restuccia 00', Martin+Restuccia '00}$$

The only dynamical degrees of freedom to quantized are (X's, A's): are univalued

$$H = \int \sqrt{W} d\sigma^1 \wedge d\sigma^2 \left[\frac{1}{2} \left(\frac{P_m}{\sqrt{W}} \right)^2 + \frac{1}{2} \left(\frac{\Pi^r}{\sqrt{W}} \right)^2 + \frac{1}{4} \{X^m, X^n\}^2 + \frac{1}{2} (\mathcal{D}_r X^m)^2 + \frac{1}{4} (\mathcal{F}_{rs})^2 + \Lambda(\{P_m, X^m\} + \mathcal{D}_r \Pi^r) \right] + \text{Ferm. term.}$$

$$\mathcal{D}_r X^m = \partial_r X^m - A_r^s \partial_s X^m, \quad \mathcal{F}_{rs} = \partial_r A_s - \partial_s A_r - [A_r, A_s]$$

With

$$D_r = 2\pi R^r \frac{\varepsilon^{ab}}{\sqrt{W}} \partial_a \hat{X}^r \partial_b$$

A N=1 2+1 symplectic NCSYM coupled to scalars proceeding from NCSYM 10D reduction=N=1 Supermembrane with Z

2 Step: Compactifying from 5D to 4D

*MPGM+Pena
+Restuccia*

Due to the condition:

$$\oint_{C_5} dX = RL_s$$

We decompose:

$$dX^7 = RL_s d\hat{X}^s + d\hat{\phi}$$

Where the decomposition is performed on the harmonic forms of the circle

The 4D hamiltonian of the M2 :

$$\begin{aligned} H_d = \int \sqrt{w} d\sigma^1 \wedge d\sigma^2 & \left[\frac{1}{2} \left(\frac{P_m}{\sqrt{W}} \right)^2 + \frac{1}{2} \left(\frac{\Pi^r}{\sqrt{W}} \right)^2 + \frac{1}{4} \{X^m, X^n\}^2 + \frac{1}{2} (\mathcal{D}_r X^m)^2 \right. \\ & + \frac{1}{4} (\mathcal{F}_{rs})^2 + \frac{1}{2} \left(F_{ab} \frac{\varepsilon^{ab}}{\sqrt{W}} \right)^2 + \frac{1}{8} \left(\frac{\Pi^c}{\sqrt{W}} \partial_c X^m \right)^2 + \frac{1}{8} [\Pi^c \partial_c (\hat{X}_r + A_r)]^2 \Big] + \\ & \Lambda \left(\left\{ \frac{P_m}{\sqrt{W}}, X^m \right\} - \mathcal{D}_r \Pi^r - \frac{1}{2} \Pi^c \partial_c \left(F_{ab} \frac{\varepsilon^{ab}}{\sqrt{W}} \right) \right) + \lambda \partial_c \Pi^c \end{aligned}$$

4. Spectral Analysis

5D CASE:

$$V = \langle \mathcal{D}_r X^m \mathcal{D}_r X^m + \frac{1}{4} \mathcal{F}_{rs} \mathcal{F}_{rs} \rangle.$$

CLASSICALLY: NO-STRING-LIKE SPIKES *MPGM+ A. Restuccia*

$$V(X^m, A_r) = 0 \Leftrightarrow X^{mB} = 0, A_r^B = 0$$

QUANTUM LEVEL: BOSONIC SECTOR *Boulton+MPGM+Martin+Restuccia*

$$V(X^m, A_r) \rightarrow \infty \quad \text{whenever} \quad (X, A) \rightarrow \infty$$

QUANTUM LEVEL FERMIONIC SECTOR *Boulton+MPGM+Restuccia*

$$\forall v_k(x) \rightarrow \infty \quad \text{when} \quad |x| \rightarrow \infty \quad \Rightarrow \quad H_{susy}^N \quad \text{discrete}$$

Eigenvalues of $V(X)$

$$\mu \equiv \Delta + V_B \mathcal{I} + V_F$$

$$X^{mB} \equiv R\phi^{mB} \quad A_r^B \equiv R\Psi^{rB}$$

$$\det\left(\left(\frac{\lambda - V_B}{R}\mathcal{I}\right) - M(\Psi, \Phi)\right) = 0$$

$$\lambda = V_B(R\phi, R\phi) + R\hat{\lambda}$$

4D CASE:

MPGM+Pena+Restuccia

The piece of the potential added is:

$$V_7 = \langle (\mathcal{D}_r X^7)^2 + \{X^m, X^7\}^2 \rangle = \langle (L_s \mathcal{D}_r \tilde{X}^s)^2 + (\mathcal{D}_r \phi)^2 + \{X^m, X^7\}^2 \rangle$$

Classically : NO string –like Spikes

$$\mathbf{V}(\mathbf{X}, \mathbf{A}, \Phi) = 0 \quad \longleftrightarrow \quad \mathbf{X} = \mathbf{A} = \Phi = \mathbf{0} \quad \text{No Flat directions}$$

Quantically: Spectrum Purely Discrete

$$H_N^B \geq c_N H_{sc,N}^B \quad \& \quad \text{Fermion contribution subdominant}$$

N=1 SUPERSYMMETRY

The hamiltonian preserves the whole supersymmetry when compactified in T^7

However the central condition charge: **THE VACUUM SOLUTION: BPS**
Dirac monopole solutions

$$\begin{aligned}\Psi &= 0 \\ X^m &= 0 \\ X_i^r &= \hat{X}_i^r\end{aligned}$$

The invariance under the spinor implies the breaking of supersymmetry. $(1/2)^3$

The action is invariant under the whole Susy and Kappa symmetry when the variations on $\delta X^r = 0$ are imposed. The central charge condition is invariant under Susy transformations. Classically there is a whole class of minima connected through susy transformations.

$$\begin{aligned}\Psi &= \varepsilon_1 + \varepsilon_2 \\ X^r &= \hat{X}^r + i\bar{\varepsilon}_2 \Gamma \varepsilon_1 \\ X^M &= i\bar{\varepsilon}_2 \Gamma^m \varepsilon_1.\end{aligned}$$



Quantically the system chooses one: **SPONTANEOUS BREAKING OF SUSY**

MODULI

There are two types of moduli in this set-up:

a. **Moduli associated to the scalar fields** $V(X) \implies$ No Flat Directions

b. **Moduli associated to the parameters R_i** of the tori.

Hip: **Tori Isotropic $R_i=R$** \implies The dependence of the potential with the radius

$$V = A + BR + CR^2 + DR^4$$

With the coefficients:

$$A = \int_{\Sigma} \sqrt{W} \left[\frac{1}{4} \{X^m, X^n\}^2 + \frac{1}{4} \{X^m, \phi\}^2 + \frac{1}{2} \{A_r, X^m\}^2 + \frac{1}{2} \{A_r, \phi\}^2 + \{A_r, A_s\}^2 \right]$$

$$C = \frac{1}{2} \int_{\Sigma} \sqrt{W} \left[(D_r X^m)^2 + (D_r \phi)^2 + (D_r A_s)^2 + \{X^m, \hat{X}^s L_s\}^2 \right]$$

$$D = \frac{1}{4} \int_{\Sigma} \sqrt{W} \{ \hat{X}^r, \hat{X}^s \}^2$$

There exist a global minimum

$$\frac{d^2 V}{dR^2} = 2C + 12DR^2 > 0$$

CONCLUSIONS

- It has been constructed the 4d hamiltonian of the supermembrane.
- The 4d supermembrane is classically stable.
- Its regularized quantum supersymmetric spectrum is discrete.
- The theory has $N=1$ supersymmetry
- There are no flat directions associated to the scalar fields parametrizing the transverse position of the supermembrane.
- For the case of the isotropic 7-tori, the potential has a global minimum with respect to the radius.

The End



MATRIX MODELS

D=11 SUPERMEMBRANES

10D SYM

$$H = T \int \sqrt{W} d\sigma^2 \left[\frac{1}{2} \left(\frac{P_m}{\sqrt{W}} \right)^2 + \frac{1}{4} \{X^m, X^n\}^2 + \Psi \Gamma_- \Gamma_M \{X^M, \Psi\} \right]$$

$$H_{ym} = \int_{\Sigma_2} \frac{g^2}{2} (\Pi^i)^2 + \frac{1}{4g^2} F_{ij}^a F_{ij}^a - \mathcal{D}_i \Pi^{ai} A_0^{ai}$$

MATRIX REGULARIZATION

$$T^{2/3} H$$

$$\frac{g^{2/3}}{V_m^{1/3}} H$$

DIMEN. REDUCCION 0+1

$H^{\text{SU}(N)}$

$$tr \left[\frac{1}{2} (\Pi^i)^2 + \frac{1}{4} (f^{abc} A_i^b A_j^c)^2 - [A_i, \Pi^i]^a A_0^{ia} \right] (t)$$

LARGE N LIMIT ?

MATRIX REGULARIZATION IN COMPACTIFIED SPACES?

PROBLEM OF CLOSED HARMONIC FORMS

TOPOLOGICAL INFORMATION?

De Wit+Peeters+Plefka