

# Higher-spin dynamics and Chern-Simons theories

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Based on:

- Johan Engquist, O.H.,  
Higher-spin Chern-Simons theories in odd dimensions, [arXiv:0705.3714 \[hep-th\]](#)
- Johan Engquist, O.H.,  
Geometry and dynamics of higher-spin frame fields, [arXiv:0708.1391 \[hep-th\]](#)

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## The higher-spin problem:

Consistent interactions of Higher-Spin (HS) fields?

→ String/M-theory [Gross (1988), Witten (1988), Sundborg (2001), Sagnotti (2004)]

Free HS theories [Fronsdal (1978), de Wit, Freedman (1980)]

Massless spin- $s$  field: totally symmetric  $h_{\mu_1 \dots \mu_s}$  with gauge symmetry

$$\delta_\epsilon h_{\mu_1 \dots \mu_s} = \partial_{(\mu_1} \epsilon_{\mu_2 \dots \mu_s)}$$

Gravity coupling  $\partial \rightarrow \nabla$  violates HS symmetry → Inconsistency!

Vasiliev's proposal:

non-vanishing cosmological constant  $\Lambda \rightarrow$  no flat-space limit

Gauging of HS algebra  $\mathfrak{h}\mathfrak{o}(D-1, 2) \supset \mathfrak{so}(D-1, 2)$

$$A_\mu = \underbrace{e_\mu^a P_a + \frac{1}{2} \omega_\mu^{ab} M_{ab}}_{\mathfrak{so}(D-1, 2)} + \text{infinite HS sum}$$

Dynamics? → 'unfolded formulation' [Vasiliev (1994)]

# Chern-Simons theory

Chern-Simons theory in  $D = 3$ :

$$S = \int \text{Tr} \left( A \wedge dA + \frac{2}{3} A \wedge A \wedge A \right)$$

→ Gravity [Witten (1988)], SUGRA [Achucarro & Townsend (1986)], HS [Blencowe (1989)]

e.o.m.  $F_{\mu\nu} = 0$  → topological, no local d.o.f.

Chern-Simons theory in  $D \geq 5$ :

$$S = \int g_{abc} A^a \wedge dA^b \wedge dA^c + \dots$$

Gauge algebra  $\mathfrak{g} = \mathfrak{so}(2n, 2)$  → Lovelock gravities in  $D = 2n + 1$ .

e.o.m.:

$$g_{abc} F^a \wedge F^b = 0$$

Background  $\bar{F}^a = 0$  → no propagator, topological phase

KK background in  $D = 5$ :  $AdS_4 \times S^1$  → non-topological phase

[Chamseddine (1990)]

# Higher-spin Chern-Simons theory

Gauge algebra: Infinite-dimensional extension of  $\mathfrak{so}(4, 2)$  [Vasiliev (2003)]:

HS generator in definite Young tableau (spin- $s$ ):

$$Q_{A_1 \dots A_{s-1}, B_1 \dots B_{s-1}} \quad \vdots \quad \underbrace{\begin{array}{|c|c|c|} \hline & & \\ \hline & & \\ \hline \end{array} \cdots \begin{array}{|c|c|c|} \hline & & \\ \hline & & \\ \hline \end{array}}_{s-1}$$

Linearisation around gravitational background for spin-3:

$$Q_{AB,CD} : \begin{array}{|c|c|} \hline & \\ \hline & \\ \hline \end{array} \longrightarrow Q_{ab,cd} : \begin{array}{|c|c|} \hline & \\ \hline & \\ \hline \end{array} \quad Q_{ab,c} : \begin{array}{|c|c|} \hline & \\ \hline & \\ \hline \end{array} \quad Q_{ab} : \begin{array}{|c|c|} \hline & \\ \hline & \\ \hline \end{array}$$

gauge field in  $\mathfrak{so}(4, 1)$  basis

$$\mathcal{W}_\mu = \bar{e}_\mu^a P_a + \frac{1}{2} \bar{\omega}_\mu^{ab} M_{ab} + \kappa \left( \frac{1}{2} e_\mu^{ab} Q_{ab} + \frac{1}{3} \omega_\mu^{ab,c} Q_{ab,c} + \frac{1}{12} \omega_\mu^{ab,cd} Q_{ab,cd} \right)$$

Spin-3 field and symmetries (after gauge-fixing  $\begin{array}{|c|c|} \hline & \\ \hline & \\ \hline \end{array}$  to zero):

$$h_{\mu\nu\rho} = \bar{e}_{(\mu}^a \bar{e}_\nu^b e_{\rho)ab}, \quad \delta_\epsilon h_{\mu\nu\rho} = \bar{D}_{(\mu} \epsilon_{\nu\rho)}$$

# HS geometry: connections, curvature and torsion

Linearized field strength:

$$\mathcal{F}_{\mu\nu} = \bar{T}_{\mu\nu}{}^a P_a + \frac{1}{2} \bar{\mathcal{R}}_{\mu\nu}{}^{ab} M_{ab} + \kappa \left( \frac{1}{2} T_{\mu\nu}{}^{ab} Q_{ab} + \frac{1}{3} T_{\mu\nu}{}^{ab,c} Q_{ab,c} + \frac{1}{12} R_{\mu\nu}{}^{ab,cd} Q_{ab,cd} \right) + \dots$$

with

$$\begin{aligned} T^{ab} &= \bar{D}e^{ab} + \omega^{ab,c} \wedge \bar{e}_c \\ T^{ab,c} &= \bar{D}\omega^{ab,c} + 3e^{\langle ab} \wedge \bar{e}^{c\rangle} + \omega^{ab,cd} \wedge \bar{e}_d, \\ R^{ab,cd} &= \bar{D}\omega^{ab,cd} + 4\omega^{\langle ab,c} \wedge \bar{e}^{d\rangle} \end{aligned}$$

Torsion constraints:

$$\begin{aligned} T^{ab} = 0 &\Rightarrow \omega^{ab,c} = \omega^{ab,c}(\partial e) \\ T^{ab,c} = 0 &\Rightarrow \omega^{ab,cd} = \omega^{ab,cd}(\partial\omega^{ab,c}) = \omega^{ab,cd}(\partial^2 e) \end{aligned}$$

[  $\rightarrow$  hierarchy of de Wit–Freedman connections in metric-like formulation]

Spin-3 Riemann tensor:  $R_{\mu\nu}{}^{ab,cd}$  3<sup>rd</sup> derivative-order!

## HS Dynamics:

How are e.o.m. of  $s$ -th derivative order consistent with Fronsdal eq.?

Damour-Deser identity:

$$(\text{Ric})_{\mu\nu|\rho\sigma} = 2\partial_{[\mu}\mathcal{F}_{\nu]\rho\sigma} = 0 ,$$

$\mathcal{F}_{\mu\nu\rho} = \square h_{\mu\nu\rho} + \dots$  HS-invariant 2<sup>nd</sup> order Fronsdal operator.

$\Rightarrow$  locally solvable (Poincaré lemma)

$$\mathcal{F}_{\mu\nu\rho} = \partial_\mu\partial_\nu\partial_\rho\alpha .$$

$\Rightarrow$  invariant under unconstrained HS transformations, provided

[Francia & Sagnotti (2002)]

$$\delta_\epsilon\alpha = \eta^{\mu\nu}\epsilon_{\mu\nu} .$$

$\Rightarrow$  gauge-fixing  $\alpha = 0 \Rightarrow$  Fronsdal equations recovered!

[Bekaert & Boulanger (2003), Sagnotti, Sezgin & Sundell (2005)]

Generalization to arbitrary spin on  $AdS$ ,

in particular: correct free field limit around  $AdS_4$ .

## Summary & Outlook:

- We constructed consistent HS-gravity coupling  $\longleftrightarrow$  consistent deformation of free HS symmetry  $\delta_\epsilon h_{\mu_1 \dots \mu_s} = \partial_{(\mu_1} \epsilon_{\mu_2 \dots \mu_s)}$
- $AdS_4 \times S^1$  background: Correct free field equations despite higher-derivative nature of Ricci tensor  
 $\longleftrightarrow$  contact with compensator formulation (trace-full generator)
- Dynamics for non-trivial HS torsion?
- Meaning of higher derivatives at non-linear level?
- Closed form of HS Lie algebra  $\longleftrightarrow$  interactions?
- Organization into  $\mathfrak{h}_0(D-1, 2)$  multiplets?  
No dynamics around  $AdS_5$ ,  
but on  $AdS_4$  off-diagonal KK components (scalar?)
- M theory as Chern-Simons theory of  $Osp(1|32)$ ? [Horava (1999)]  
Recently: 'Non-critical M-theory' in  $D=3$   
 $\longleftrightarrow$  HS Chern-Simons theory on  $AdS_2 \times S^1$  [Horava & Keeler (2007)]