

Consistent Kaluza-Klein reductions and susy AdS solutions

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based on work with

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Consistent KK reductions

- A very useful way to construct solutions to sugra theories in a higher dimension D is to uplift solutions of simpler sugras in lower dimension d .
- For this **uplift to be well defined**, there must exist a **consistent Kaluza-Klein (KK) reduction** from the sugra in dimension D to the sugra in dimension d .
- To determine if such KK reduction is consistent is an interesting problem by its own.

Consistent KK reductions

- Upon compactification on an internal manifold M_{D-d} , the D -dimensional fields give rise to d -dimensional fields: a finite set of *light* L and a KK tower of *heavy* H fields.
- The D -dimensional e.o.m.'s can be rewritten in terms of these:

$$\square L \sim a_{mn} L^m H^n$$

$$\square H \sim b_{mn} L^m H^n$$

- A truncation keeping L and discarding H ($H = 0$) will be consistent only if $b_{m0} = 0$.
- Then, the fields L satisfy *d -dimensional e.o.m.'s*: $\square L = 0$. Any solution to the *d -dimensional theory can then be uplifted to D dimensions*.

Consistent KK reduction

- Only in a few cases there is a group-theoretical argument behind the consistency of the truncation:
- (Toroidal) dimensional reductions,
- Compactifications on group manifolds.
- But in general, the compactification on arbitrary manifolds will be inconsistent.

$AdS \times Sphere$ compactifications

- Remarkable consistent compactifications of $D = 10, 11$ sugra are associated with the (maximally supersymmetric) solutions $AdS_7 \times S^4$, $AdS_4 \times S^7$ and $AdS_5 \times S^5$:
- The compactification of $D = 11$ sugra on S^4 can be consistently truncated to gauged maximal supergravity in $d = 7$ [Nastase, Vaman, van Nieuwenhuizen, hep-th/9905075, 9911238].
- The compactification of $D = 11$ sugra on S^7 can be consistently truncated to gauged maximal supergravity in $d = 4$ [De Wit, Nicolai, NPB 281 (1987) 211] .
- Similarly, the compactification of IIB sugra on S^5 is expected to consistently yield maximal gauged $d = 5$ sugra [Cvetic, Duff, Hoxha, Liu, Lu, Lu, Martinez-Acosta, Pope, Sati, Tran, hep-th/9903214; Lu, Pope, Tran, hep-th/9909203; Cvetic, Lu, Pope, Sadrzadeh, Tran, hep-th/0003103] .

A conjecture about consistency

- String/M-theory on all these backgrounds is dual, via the AdS/CFT correspondence, to a superconformal field theory (SCFT) in the boundary of AdS .
- Indeed, we would like to view the compactifications on those backgrounds as special cases of the following conjecture:
- For any supersymmetric $AdS_d \times_w M_{D-d}$ solution of $D = 10$ or $D = 11$ supergravity there is a consistent Kaluza-Klein truncation on M_{D-d} to a gauged supergravity theory in d -dimensions for which the fields are dual to those in the superconformal current multiplet of the $(d - 1)$ -dimensional dual SCFT [Gauntlett, OV, arXiv:0707.2315].

A conjecture about consistency

- Equivalently, the fields of the gauged supergravity are those that contain the d -dimensional graviton and fill out an irreducible representation of the superisometry algebra of the $D = 10$ or $D = 11$ supergravity solution $AdS_d \times_w M_{D-d}$.
- This is essentially a restricted version of the conjecture in [Duff, Pope, Nucl. Phys. **B255** (1985) 355].
- General arguments supporting it were subsequently put forward in [Pope, Stelle, Phys. Lett. **B198** (1987) 151].

A conjecture about consistency

- For example, the $AdS_5 \times S^5$ solution of type IIB, which has superisometry algebra $SU(2, 2|4)$, is dual to $N = 4$ superYang-Mills theory in $d = 4$.
- The superconformal current multiplet of the latter theory includes the energy momentum tensor, $SO(6)$ R-symmetry currents, along with scalars and fermions.
- These are dual to the metric, $SO(6)$ gauge fields along with scalar and fermion fields, and are precisely the fields of the maximally supersymmetric $SO(6)$ gauged supergravity in $d = 5$.

A conjecture about consistency

- Here we will give evidence of this conjecture for the case of $AdS_5 \times_w M_6$ solutions in $D = 11$, which are dual to $N = 1$ SCFTs in 4 dimensions.
- These SCFTs all have a $U(1)$ R-symmetry and so we expect that $D = 11$ sugra on M_6 gives a $d = 5$ sugra with
 - $N = 1$ supersymmetry
 - a metric ds_5^2 (dual to the energy-momentum tensor of the SCFT)
 - and a $U(1)$ gauge field A (dual to the R-symmetry current).
- This is precisely the content of minimal $d = 5$ gauged sugra.

AdS₅ ×_w M₆ solutions of D = 11 sugra

The most general solution of the D = 11 sugra equations containing an AdS₅ factor in the metric was analysed by [Gauntlett, Martelli, Sparks, Waldram, hep-th/0402153] using G-structure techniques [Gauntlett, Martelli, Pakis, Waldram, hep-th/0205050].

The most general form of the D = 11 bosonic fields ds_{11}^2 , $G_4 = dA_3$ containing AdS₅, compatible with SO(4, 2) symmetry is

$$ds_{11}^2 = e^{2\lambda} [ds^2(AdS_5) + ds^2(M_6)] ,$$

$$G_4 \in \Omega_4(M_6, \mathbb{R})$$

$$\lambda \in \Omega_0(M_6, \mathbb{R}) \quad (\text{warp factor})$$

and subject to the field equations.

AdS₅ ×_w M₆ solutions of D = 11 sugra: N = 1 supersymmetry

- In order to have **N = 1 supersymmetry**, the solution must admit a Killing spinor ϵ , solution to the **Killing spinor equation**

$$\mathcal{D}\epsilon = 0$$

where \mathcal{D} is the supercovariant derivative, involving the ordinary Riemannian covariant derivative and the D = 11 sugra fields.

AdS₅ ×_w M₆ solutions of D = 11 sugra: N = 1 supersymmetry

- The D = 11 Killing spinor splits as

$$\epsilon = \varepsilon \otimes e^{\lambda/2} \xi,$$

where ε is a Killing spinor on AdS₅ and ξ is a (non-chiral) spinor on M₆.

- The Killing spinor equation also splits into
 - an equation for ε on AdS₅, immediately satisfied, and
 - equations for ξ on M₆, which specify a particular G-structure on M₆ (with G = SU(2)).

AdS₅ ×_w M₆ solutions of D = 11 sugra: the G-structure on M₆

The existence of ξ defines a G -structure on M_6 , alternatively specified by a set of bilinears on ξ (e.g., $\tilde{K}_m^2 = \frac{1}{2}\bar{\xi}\gamma_m\gamma_7\xi$):

$$K^1, \tilde{K}^2 \in \Omega_1(M_6, \mathbb{R})$$

$$J \in \Omega_2(M_6, \mathbb{R})$$

$$\Omega \in \Omega_2(M_6, \mathbb{C})$$

$$\cos \zeta \in \Omega_0(M_6, \mathbb{R})$$

AdS₅ ×_w M₆ solutions of D = 11 sugra: the G-structure on M₆

- The Killing spinor equations for ξ translate into a set of differential and algebraic equations among these bilinear forms and the warp factor λ and four-form G_4 .
- e.g., $\nabla_{(\mu} \tilde{K}_{\nu)}^2 = 0$, i.e., $\tilde{K}^2 \equiv \cos \zeta K^2$ defines a Killing vector (related to the R-symmetry of the dual CFT)
- The equations among the bilinear forms constrain the internal geometry, i.e., the metric on M_6 , the flux G_4 and the warp factor λ .

Kaluza-Klein ansatz: the metric

- The ‘KK ansatz’ must express the $D = 11$ fields ds_{11}^2 , $G_4 = dA_3$ in terms of the $d = 5$ fields ds_5^2 , $F_2 = dA$. The KK ansatz must then satisfy the $D = 11$ field equations, provided ds_5^2 , F_2 satisfy the $d = 5$ equations.
- It is natural to think of the $d = 5$ $U(1)$ gauge field A as arising from the $U(1)$ isometry of M_6 generated by \tilde{K}^2 .
- Thus, we take the usual KK ansatz for the metric:

$$ds_{11}^2 = e^{2\lambda} [ds_5^2 + ds^2(\hat{M}_6)]$$

where \hat{M}_6 denotes the deformation of M_6 parametrised by A as

$$K^2 \longrightarrow \hat{K}^2 = K^2 + \frac{1}{3} \cos \zeta A$$

Kaluza-Klein ansatz: the four-form

KK ansatz for the four-form:

$$G'_4 = \hat{G}_4 + F_2 \wedge \hat{\beta}_2 + *_5 F_2 \wedge \hat{\beta}_1.$$

Here,

- $F_2 = dA$
- hatted quantities are forms on \hat{M}_6 (*i.e.* with $K^2 \longrightarrow \hat{K}^2 = K^2 + \frac{1}{3} \cos \zeta A$)
- G_4 is the four-form on M_6 corresponding to the undeformed background $AdS_5 \times_w M_6$
- β_1, β_2 are forms on M_6 to be determined.

Consistent truncation

- Direct substitution shows that the KK ansatz satisfies the $D = 11$ field equations provided that
 - the $d = 5$ fields satisfy the equations of minimal $d = 5$ gauged supergravity, and
 - a set of differential and algebraic equations among β_1 , β_2 and the warp factor λ and four-form G_4 is satisfied.
- These equations are actually of **the same** form than those among the bilinear forms defining the G -structure on M_6 .
- Indeed β_1, β_2 have a solution in terms of some of the spinor bilinears on M_6 :

$$\beta_1 = -\frac{1}{3}e^{3\lambda} \cos \zeta K^1$$

$$\beta_2 = \frac{1}{3}e^{3\lambda} (-\sin \zeta J + K^1 \wedge K^2) .$$

Consistent truncation

- To summarise, the $d = 5$ fields ds_5^2 , F_2 can be embedded into the $D = 11$ fields ds_{11}^2 , G_4 through the KK ansatz

$$ds_{11}^2 = e^{2\lambda}[ds_5^2 + ds^2(\hat{M}_6)] , \quad \tilde{K}^2 \longrightarrow \hat{K}^2 = \tilde{K}^2 + A$$

$$G'_4 = \hat{G}_4 + F_2 \wedge \frac{1}{3}e^{3\lambda}(-\sin\zeta J + K^1 \wedge \hat{K}^2) - *_5F_2 \wedge \frac{1}{3}e^{3\lambda} \cos\zeta K^1$$

- **This shows the consistency** of the truncation, at the level of the bosonic equations
[Gauntlett, O Colgain, OV, hep-th/0611219] .
- The $D = 11$ gravitino variations also reduce consistently to the $d = 5$ gravitino variation.

Further examples in $D = 11$

Other examples support our conjecture about consistent KK reductions.

- The $D = 11$ solutions of the form $AdS_4 \times SE_7$, where SE_7 is Sasaki-Einstein are dual to 3d $N = 2$ SCFTs, and the reduction of $D = 11$ on M_7 consistently truncates to $d = 4$, $N = 2$ gauged sugra [Gauntlett, OV, arXiv:0707.2315] .
- The $D = 11$ solutions of the form $AdS_4 \times_w M_7$, corresponding to M5-branes wrapping SLAG 3 cycles [Gauntlett, Mac Conamhna, Mateos, Waldram hep-th/0605146] , also allow for a consistent reduction of $D = 11$ sugra on M_7 to $d = 4$, $N = 2$ gauged sugra [Gauntlett, OV, arXiv:0707.2315] .

Further examples in IIB

- IIB sugra on $d = 5$ Sasaki-Einstein spaces is dual to a 4d $N = 1$ SCFT, and consistently truncates to minimal $d = 5$ gauged sugra [Buchel, Liu, hep-th/0608002] .
- The IIB solutions of the form $AdS_5 \times_w M_5$ with $N = 1$ susy and all fluxes active [Gauntlett, Martelli, Sparks, Waldram, hep-th/0510125] also allow for a consistent reduction of IIB on M_5 to $d = 5$ minimal gauged sugra [Gauntlett, OV, arXiv:0707.2315] .

Conclusions and outlook

- All known consistent KK truncations on spheres could be recast in this language.
- The consistency of the KK truncation makes it manifest from the gravity side that SCFTs with type a IIB or $D = 11$ dual share common sectors. E.g., black hole solutions in lower d sugra should be relevant to any of the SCFTs to which these sugras are related.
- It would be interesting to prove the conjecture from the CFT side.