## World-sheet description of $A$ and $B$ branes revisited

## Alexander Wijns (University of Iceland)

in collaboration with A. Sevrin and W. Staessens (V.U. Brussels)
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## Outline

(1) Motivation

- Nonlinear $\sigma$-models in superspace
- D-brane effective actions


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- Formalism
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(4) Type B branes
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- Derivative corrections


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## Nonlinear $\sigma$-models

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- Geometry of target space clarified by using $N=(2,2)$ superspace


## Closed string superspace: $N=(2,2)$

- $N=(2,2)$ superspace action determined by Kähler potential

$$
\mathcal{S}_{(2,2)}=\int d^{2} \sigma D_{+} D_{-} \hat{D}_{+} \hat{D}_{-} V(X, \bar{X})
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BUT one needs constrained superfields!

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- Most general $N=(2,2)$ superspace description in terms of chiral, twisted chiral and semi-chiral fields [LindSTRÖm, Roček, von Unge, Zabzine '05]


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$\bullet \rightarrow$ target space is Kähler manifold with metric $g_{\alpha \bar{\beta}}= \pm V_{\alpha \bar{\beta}}$


## Boundaries?

## First motivation

What happens when we include boundaries?
[Ooguri, Oz, Yin; Albertsson, Lindström, Zabzine, ...]
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Completely local $N=2$ superspace formulation still incomplete
$\downarrow$
World-sheet description of $A$ and $B$ branes on Kähler manifolds
B branes: - holomorphic cycles

- holomorphic line bundle connection

A branes: - lagrangian cycles with flat connection

- coisotropic cycles with non-vanishing flux


## D-brane effective actions

- Effective action for flat Dp-brane for slowly varying fields 10d Born-Infeld action reduced to $p+1$ dimensions

$$
\mathcal{S}_{B I}=-\tau_{9} \int d^{10} x \sqrt{-\operatorname{det} h_{a b}^{ \pm}}, \quad h_{a b}^{ \pm}=\eta_{a b} \pm F_{a b}
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- For holomorphic line bundle connection: $F_{\alpha \beta}=F_{\bar{\alpha} \bar{\beta}}=0$

$$
g^{\alpha \bar{\beta}}(\operatorname{arcth} F)_{\alpha \bar{\beta}}=0
$$

solves the BI equations of motion $\rightarrow$ Deformation of DUY stability condition

$$
g^{\alpha \bar{\beta}} F_{\alpha \bar{\beta}}=0
$$

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- Multiple coinciding D-branes $\rightarrow$ non-abelian gauge theory


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$$
[D, D] F=[F, F]
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|  | $\alpha^{\prime}=0$ | $\alpha^{\prime} \neq 0$ | $\alpha^{\prime} \neq 0$ |
| :---: | :---: | :---: | :---: |
|  | $\partial F=0$ | $\partial F=0$ | $\partial F \neq 0$ |
| $[]=0$, | Maxwell | Born-Infeld | $!$ |
| $[] \neq 0$, | Yang-Mills | up to 4th order in $\alpha^{\prime}$ |  |

$\rightarrow$ look at derivative corrections to BI

## Bl action + derivatives?

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How to compute derivative corrections efficiently?

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Simple $N=2$ superspace description of space-filling B brane in flat space with holomorphic line bundle connection

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Simple $N=2$ superspace description of space-filling B brane in flat space with holomorphic line bundle connection
$\square$
Derivative corrections from $\beta$-function calculation in $N=2$ superspace using supergraph techniques

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## $\mathrm{N}=1$ superspace

- Introduce a boundary $\rightarrow$ from $N=(1,1)$ to $N=1$ [Koerber, Nevens, Sevrin '03]

$$
\begin{aligned}
D & =D_{+}+D_{-} \\
D^{\prime} & =D_{+}-D_{-} \\
D^{2}=D^{\prime 2} & =-\frac{i}{2} \partial_{\tau}, \quad\left\{D, D^{\prime}\right\}=-i \partial_{\sigma}
\end{aligned}
$$

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$$
\left.\begin{array}{rl}
\begin{array}{c}
D
\end{array}=D_{+}+D_{-} & \text {unbroken } \\
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\text {broken }
\end{array}\right] \begin{array}{cc}
D^{2}=D^{\prime 2}=-\frac{i}{2} \partial_{\tau}, \quad\left\{D, D^{\prime}\right\}=-i \partial_{\sigma}
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$$

- $N=1$ action

$$
\mathcal{S}_{N=1}=-\int d^{2} \sigma D\left[D^{\prime}\left(D_{+} X^{a} D_{-} X^{b}\left(g_{a b}+b_{a b}\right)\right)\right]
$$

Equivalent to bulk action up to boundary term:
$D_{+} D_{-}=-1 / 2 D D^{\prime}-i / 2 \partial_{\sigma}$

## $\mathrm{N}=1$ superspace: boundary conditions

- Boundary conditions
- Dirichlet: $\delta X^{a}=0$
- Neumann: $D^{\prime} X^{a}=b^{a}{ }_{b} D X^{b}$


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- Boundary conditions
- Dirichlet: $\delta X^{a}=0$
- Neumann: $D^{\prime} X^{a}=b^{a}{ }_{b} D X^{b}$
- Mixed conditions: introduce projection operators

$$
\mathcal{P}_{ \pm}=\frac{1}{2}(1 \pm \mathcal{R}), \quad \mathcal{R}^{2}=1
$$

$\mathcal{P}_{+}$: Neumann, $\mathcal{P}_{-}$: Dirichlet

$$
\mathcal{P}_{-} \partial_{\tau} X=\mathcal{P}_{-} D X=0 \rightarrow \mathcal{P}_{+b}^{d} \mathcal{P}_{+c}^{e} \mathcal{R}_{[d, e]}^{a}=0
$$

$\rightarrow \mathcal{P}_{+}$is integrable
$\rightarrow$ Brane wraps integrable submanifold

## $\mathrm{N}=2$ superspace

- From $N=(2,2)$ to $N=2$ : two choices B-type:

$$
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D & =D_{+}+D_{-}, & & \hat{D}=\hat{D}_{+}+\hat{D}_{-}, \\
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A-type:

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- But chiral with A-type $=$ twisted chiral with B-type


## $\mathrm{N}=2$ superspace: Action

$N=2$ action

$$
\mathcal{S}_{N=2}=\int d^{2} \sigma D \hat{D}\left[D^{\prime} \hat{D}^{\prime} V(X, \bar{X})\right]+i \int d \tau D \hat{D} W(X, \bar{X})
$$

with

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\begin{gathered}
D^{2}=\hat{D}^{2}=D^{\prime 2}=\hat{D}^{\prime 2}=-\frac{i}{2} \partial_{\tau}, \\
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| Chiral with B-type | $=$ | type B branes |
| ---: | :---: | :--- |
| $\uparrow$ | mirror symmetry | $\downarrow$ |
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## Formalism

- Twisted chiral fields with B-type boundary

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- Impose Dirichlet conditions

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- Neumann conditions

$$
\begin{aligned}
& \left(\pi_{+} \mathcal{P}_{+} D^{\prime} X\right)^{\mu}=\mathcal{R}^{\mu}{ }_{\nu} D^{\prime} X^{\nu} \\
& \left(\pi_{-} \mathcal{P}_{+} D^{\prime} X\right)^{\mu}=0
\end{aligned}
$$

$\rightarrow$ non-degenerate $F$ along $\pi_{+}$and $F=0$ along $\pi_{-}$

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$\rightarrow \operatorname{Im} \pi_{+}$is $4 m$-dimensional, $m \in \mathbb{N}$
- In general this thus describes an $(n+2 m)$-brane where $d=2 n$


## Geometric interpretation

- If $\sigma^{a}$ are coordinates along the brane $\mathcal{N}$ in $\pi_{-}$directions

$$
T_{\mathcal{N}}^{\perp}=\left\{\partial / \partial \sigma^{a}\right\} \subset T_{\mathcal{N}}=\left\{\partial / \partial \sigma^{a}\right\} \oplus \operatorname{Im} \pi_{+}
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Brane wraps coisotropic submanifold

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- Pullback of $\omega, F$ and $K=\omega^{-1} F$ only non-vanishing and non-degenerate on $T_{\mathcal{N}} / T_{\mathcal{N}}^{\perp}=\operatorname{Im} \pi_{+}$
[Kapustin, Orlov '01]
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- Pullback of $\omega, F$ and $K=\omega^{-1} F$ only non-vanishing and non-degenerate on $T_{\mathcal{N}} / T_{\mathcal{N}}^{\perp}=\operatorname{Im} \pi_{+}$
- When $\pi_{-}=1$, we find $T_{\mathcal{N}}^{\perp}=T_{\mathcal{N}} \rightarrow$ brane wraps lagrangian submanifold with flat connection
[Kapustin, Orlov '01]
[Lindström, Zabzine '02]


## Examples:

- Lagrangian brane of any shape on $T^{2}$
$\rightarrow W$ determines the shape of the brane


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$\rightarrow W$ determines the shape of the brane
- Maximally coisotropic brane on $T^{4}$ and 4d hyperkähler manifold $\mathcal{K}$ with $V \equiv V(z-\bar{z}, w+\bar{w})$

In both cases:

$$
\begin{gathered}
F_{z w}=i \\
W=\frac{i}{2}\left(z V_{z}+w V_{w}-\bar{z} V_{\bar{z}}-\bar{w} V_{\bar{w}}\right)
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In general: possible corrections from higher derivatives on $V$

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- 5-brane on $T^{6}$ or $T^{2} \times \mathcal{K}$


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## Formalism

- Chiral fields with B-type boundary

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- Carries holomorphic flux: $F_{\alpha \bar{\beta}}=-i W_{\alpha \bar{\beta}}, \quad F_{\alpha \beta}=F_{\bar{\alpha} \bar{\beta}}=0$


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- B brane wraps holomorphic cycle on Kähler manifold
- Carries holomorphic flux: $F_{\alpha \bar{\beta}}=-i W_{\alpha \bar{\beta}}, \quad F_{\alpha \beta}=F_{\bar{\alpha} \bar{\beta}}=0$
- Application: space filling brane in flat space

$$
D^{\prime} X^{\alpha}=F_{\beta}^{\alpha} D X^{\beta}
$$

$\rightarrow$ study stability by computing $\beta$-function

## Conformal invariance

- Conformal invariance $\rightarrow$ vanishing $\beta$-function $\rightarrow$ stability condition


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- Example: closed strings (no b-field)
$N=(2,2)$ supersymmetry $\quad \rightarrow$ Kähler manifold
One-loop conformal invariance $\rightarrow$ Ricci flat
Four-loop conformal invariance $\rightarrow R^{4}$ term


## Conformal invariance

- Conformal invariance $\rightarrow$ vanishing $\beta$-function $\rightarrow$ stability condition
- Example: closed strings (no b-field)
$\mathrm{N}=(2,2)$ supersymmetry
$\rightarrow$ Kähler manifold
One-loop conformal invariance $\rightarrow$ Ricci flat
Four-loop conformal invariance $\rightarrow R^{4}$ term
- Four-loop calculation is greatly simplified by using $N=(2,2)$ superspace techniques [Grisaru, van de Ven, Zanon '86]


## Deformed stability condition from $\beta$-function

- Open strings on flat space-filling B brane coupled to boundary potential W
[Nevens, Sevrin, Troost, AW '06]


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- At one loop: $\beta(W) \propto g^{\alpha \bar{\beta}}(\operatorname{arcth} F)_{\alpha \bar{\beta}} \rightarrow \mathrm{BI}$ action
- Two loops: no contribution to the $\beta$-function $\rightarrow$ no two-derivative corrections


## [Nevens, Sevrin, Troost, AW '06]

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- Open strings on flat space-filling B brane coupled to boundary potential W
- At one loop: $\beta(W) \propto g^{\alpha \bar{\beta}}(\operatorname{arcth} F)_{\alpha \bar{\beta}} \rightarrow \mathrm{BI}$ action
- Two loops: no contribution to the $\beta$-function $\rightarrow$ no two-derivative corrections
- Three loops: four-derivative correction to the BI action


## [Nevens, Sevrin, Troost, AW '06]

## Deformed stability condition from $\beta$-function

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- Three loops: four-derivative correction to the BI action
- All results in complete agreement with renormalization group equations


## [Nevens, Sevrin, Troost, AW '06]

## Result: $\mathrm{BI}+$ derivative corrections

Abelian effective action up to 4 derivatives [WylLard '01]

$$
\begin{aligned}
\mathcal{S}= & -\tau_{9} \int d^{10} \times \sqrt{-h_{+}}\left[1+\frac{1}{96}\left(\frac{1}{2} h_{+}^{\mu \nu} h_{+}^{\rho \sigma} S_{\nu \rho} S_{\sigma \mu}\right.\right. \\
& \left.\left.-h_{+}^{\rho_{2} \mu_{1}} h_{+}^{\mu_{2} \rho_{1}} h_{+}^{\sigma_{2} \nu_{1}} h_{+}^{\nu_{2} \sigma_{1}} S_{\mu_{1} \mu_{2} \nu_{1} \nu_{2}} S_{\rho_{1} \rho_{2} \sigma_{1} \sigma_{2}}\right)\right]
\end{aligned}
$$

With stability condition [Koerber '04]
$g^{\alpha \bar{\beta}}(\operatorname{arcth} F)_{\alpha \bar{\beta}}+\frac{1}{96} S_{a b \alpha \bar{\beta}} S_{c d \gamma \bar{\delta}} h_{+}^{b c} h_{+}^{d a}\left(h_{+}^{\alpha \bar{\delta}} h_{+}^{\gamma \bar{\beta}}-h_{-}^{\alpha \bar{\delta}} h_{-}^{\gamma \bar{\beta}}\right)=0$
where $S_{a b c d}=\partial_{a} \partial_{b} F_{c d}+2 h_{+}^{e f} \partial_{a} F_{[c \mid e} \partial_{b} F_{\mid d] f}, \quad h_{\alpha \bar{\beta}}^{ \pm}=\eta_{\alpha \bar{\beta}} \pm F_{\alpha \bar{\beta}}$

## Outline

(1) Motivation

- Nonlinear $\sigma$-models in superspace
- D-brane effective actions
(2) Boundary superspace
- $N=1$
- $N=2$
(3) Type A branes
- Formalism
- Coisotropic branes
(4) Type B branes
- Formalism
- Derivative corrections
(5) Duality transformations
(6) Remarks


## Duality transformations

- Closed strings
- In the presence of an isometry: chiral field dualized to twisted chiral field
- Explicit transformation by gauging the isometry and passing through a first order action


## Duality transformations

- Closed strings
- In the presence of an isometry: chiral field dualized to twisted chiral field
- Explicit transformation by gauging the isometry and passing through a first order action
- Open strings
- Analogous to closed string case, but boundary conditions of first order action should be treated with care
- Example 1: space filling B brane on $d=2 n$-dimensional Kähler manifold parameterized by $n$ chiral fields is dual to $n$-dimensional Lagrangian A brane on dual Kähler manifold parameterized by $n$ twisted chiral fields
- Example 2: One of the two isometries of maximal coisotropic brane on hyperkähler manifold can be dualized $\rightarrow 3$-brane on generalized Kähler manifold


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## Conclusion and outlook

## Conclusions

- $N=2$ superspace description of $A$ and $B$ branes on Kähler manifolds
- Explicit examples of coisotropic branes
- Duality transformations
- Application: D-brane effective action


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## Outlook

- Geometries parameterized by chiral + twisted chiral (+ semi-chiral) $\rightarrow$ branes on generalized complex geometries
- Conformal invariance $\rightarrow$ stability conditions [Marino, Minasian, Moore, Strominger '00; Kapustin, Li '03; Koerber '05]

