

# Induced GW from Warm Inflation

Cold inflation/Warm inflation

Slow-roll warm inflation

Primordial spectrum: PBH & GW

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[arXiv:2105.08045]



# Expanding Universe

Flatness problem

$$\Omega_T = 1 \rightarrow \Omega_T(t_{\text{nucl}}) - 1 \approx 10^{-16}$$

Horizon problem

The observable Universe was larger than the **particle horizon** at LSS

Inflation

Early period of accelerated expansion

$$\ddot{a} > 0 : P < -\rho/3$$

Primordial spectrum?

Too small sub-horizon  
**(causal)** perturbations

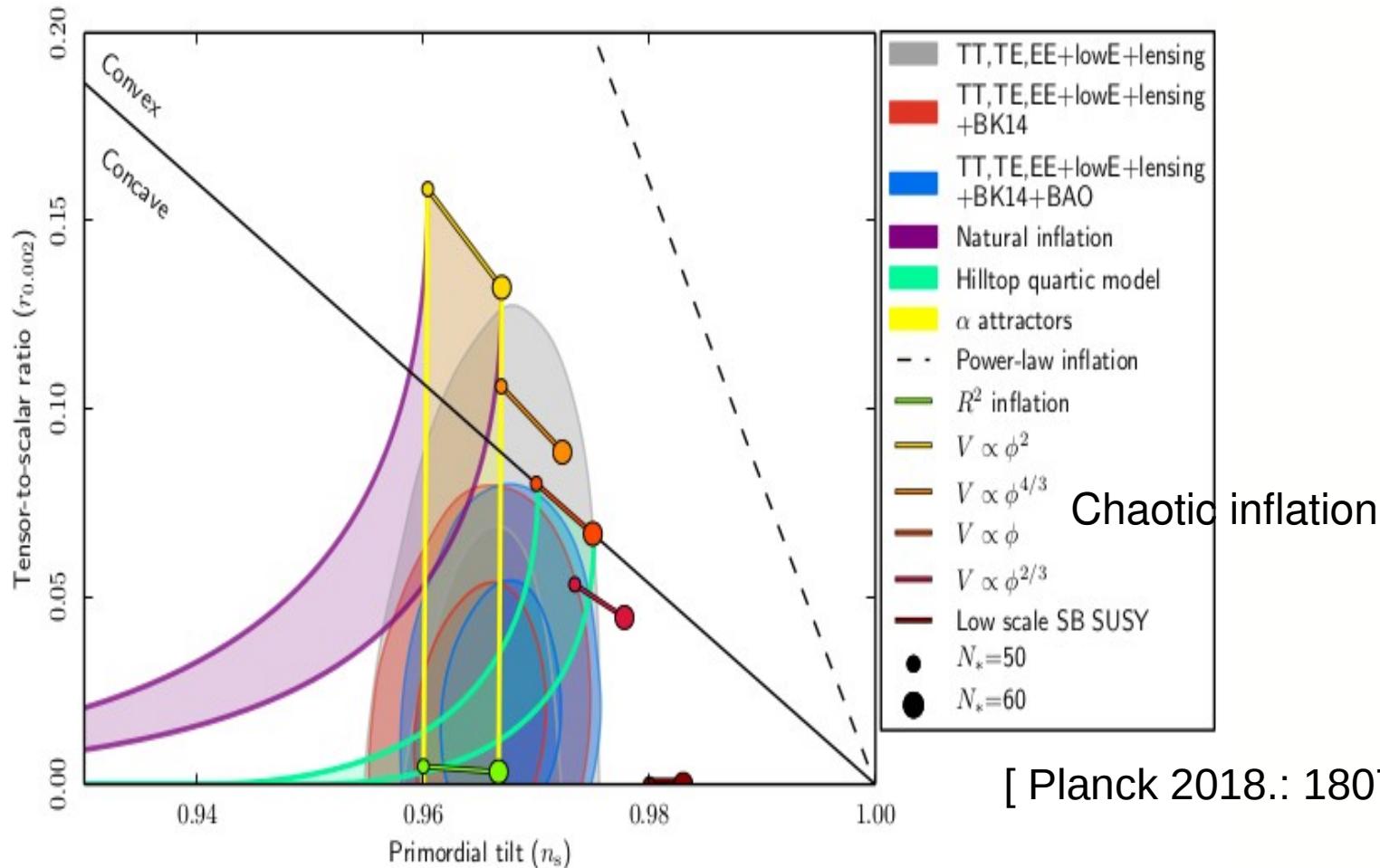
Unwanted relics...

**monopoles**, moduli, gravitinos,...

# Primordial spectrum: ~adiabatic, ~scale-invariant, gaussian?, tensors?

**Primordial spectrum:**  $P_R = P_R(k_0)(k/k_0)^{n_s-1}$      $k_0 = 0.002 \text{ Mpc}^{-1}$

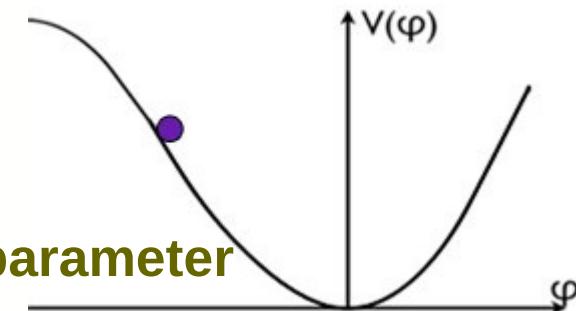
**Tensor-to-scalar Ratio:**  $r = P_T/P_R$      $P_R = 2.2 \times 10^{-9}$



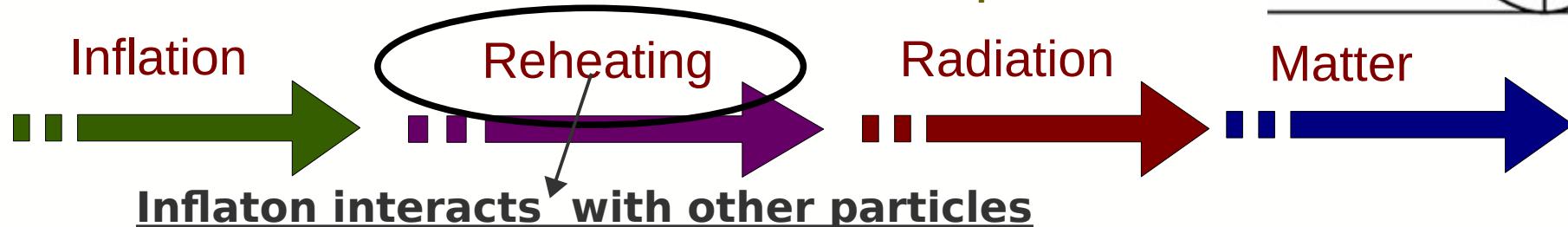
**CMB constraints ~10 efolds of inflation....**

## Slow Roll Inflation

Scalar field rolling down its (flat) potential

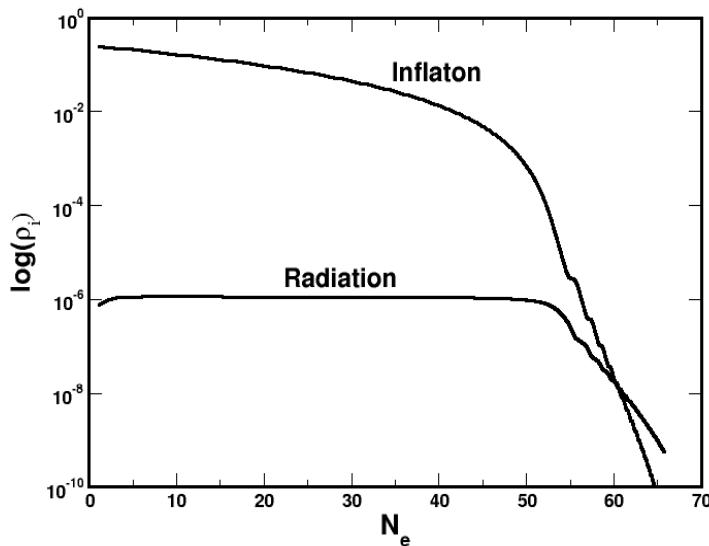
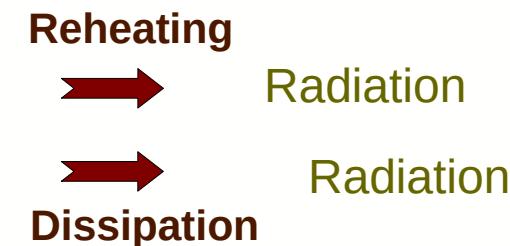


Kinetic energy << potential energy  $H^2 \sim V/3m_P^2$  Hubble parameter



"Cold" inflation: Interactions negligible during Inflation

"Warm" inflation: Inflaton "decay" into light dof



A (small) fraction of the vacuum energy is converted into radiation during inflation

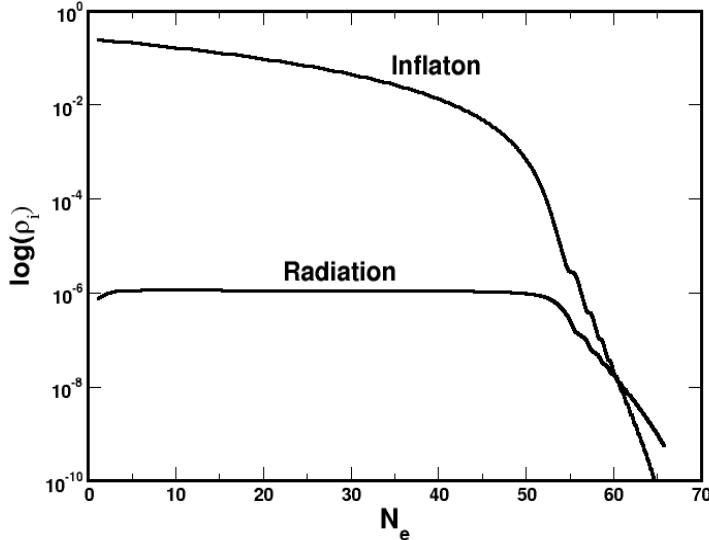
$$\ddot{\phi} + (3H + Y)\dot{\phi} + V_{\phi} = 0$$

$$\dot{\rho}_R + 4H\rho_R = Y\dot{\phi}^2 \quad \text{"Source term"}$$

"Decay" into light dof= dissipation= extra friction

# Slow Roll Warm Inflation

Scalar field rolling down its (flat) potential+ dissipation  $\Upsilon$



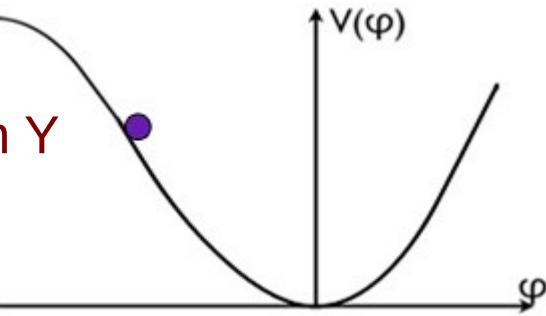
$$Q = \Upsilon / 3H$$

$$\left. \begin{aligned} 3H(1+Q)\dot{\varphi} &\simeq -V_{\varphi} \\ 4\rho_R &\simeq 3Q\dot{\varphi}^2 \end{aligned} \right\}$$

$$\epsilon_{\varphi} = \frac{m_P^2}{2} \left( \frac{V_{\varphi}}{V} \right)^2 \ll 1+Q$$

$$\eta_{\varphi} = \frac{V_{\varphi\varphi}}{3H^2} \ll 1+Q$$

Slow-roll parameters



Inflation lasts longer due to extra friction  $\rightarrow H_{60}^{\text{warm}} < H_{60}^{\text{cold}}$   $\rightarrow r < 0.1$

Slow-roll parameters can be larger than 1, when  $Q > 1$   $\leftarrow$

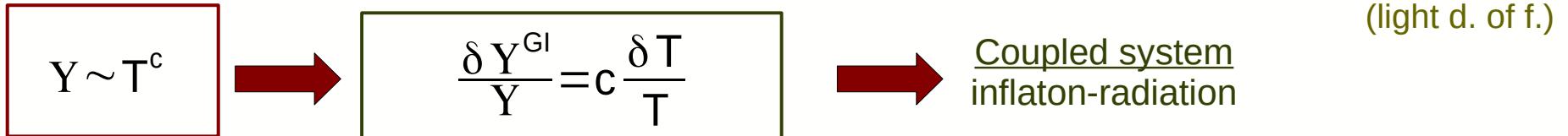
- $\left\{ \begin{array}{l} \underline{\eta \text{ problem}} : \text{inflaton mass can be larger than } H \text{ (sugra models)} \\ \underline{\text{Swampland conjecture}} : \text{both slow-roll parameters can be larger than one} \\ \text{(inflation & string theory)} \end{array} \right.$

# Fluctuations & primordial spectrum: coupled system

Field EOM:

$$\delta \ddot{\varphi}_k + (3H + Y) \delta \dot{\varphi}_k + \dot{\varphi} \delta Y + \left( \frac{k^2}{a^2} + V_{\varphi\varphi} \right) \delta \varphi_k \simeq (2 YT)^{1/2} \hat{\xi}_k$$

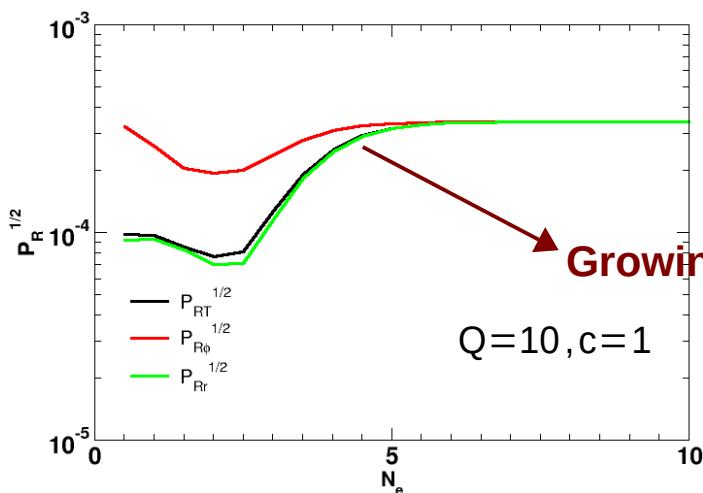
fluctuation force  $\hat{\xi}$



(light d. of f.)

[Ramos, da Silva, 1302.3544; BG, Berera, Moss & Ramos, 1401.1149]

[ $R$  is constant after horizon crossing (freeze-out)]



$$P_R = \frac{h_\varphi}{h_T} P_{R\varphi} + \frac{h_r}{h_T} P_{Rr} \simeq P_{Rr} \simeq P_{R\varphi}, \quad (h_i = \rho_i + p_i)$$

When  $c > 0$ ,  $G[Q] \sim Q^\alpha$

Blue tilted spectrum  
Amplification of the  
primordial spectrum

We want to have  $Y \sim 1/T$  at CMB scales to avoid the effect on the spectral index of the "growing mode", but to allow  $Y \sim T^c$  later on to amplify the primordial spectrum

[Green & Liddle PRD56 '97]



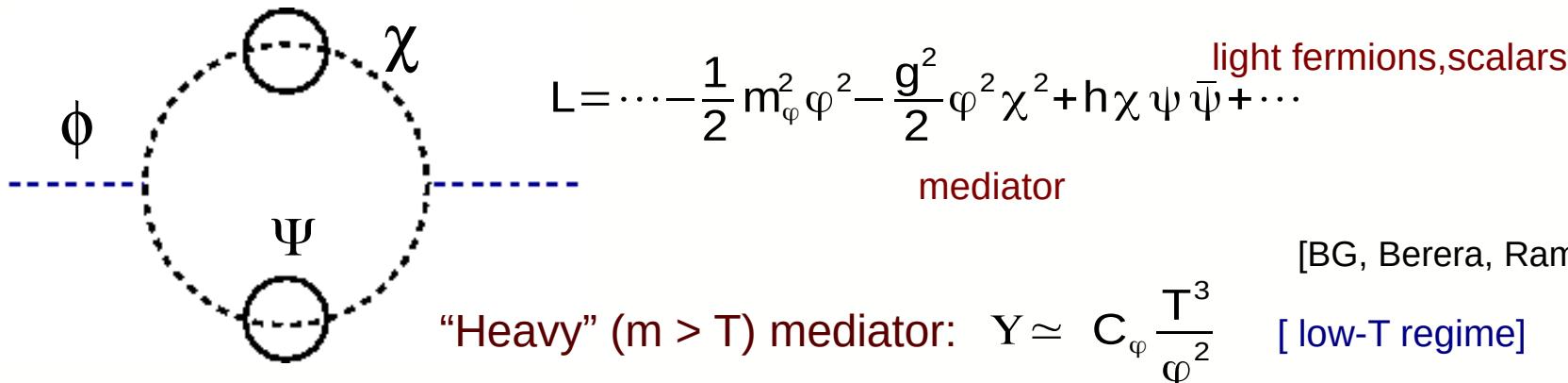
PBHs??

[R. Arya 1910.05238]

Second order source of primordial tensors??

[Mollerach, Harari, Matarrese PRD69 '04; Ananda, Clarkson, Wands PRD75 '07]

## Interactions & (T-dependent) Dissipative coefficient



[BG, Berera, Ramos & Rosa 2012]

"Heavy" ( $m > T$ ) mediator:  $Y \simeq C_\varphi \frac{T^3}{\varphi^2}$  [low-T regime]

Thermal corrections under control (inflaton coupled to heavy fields), but may required large no. of mediator fields

"Light" ( $m < T$ ) mediator:  $Y \simeq C_\varphi T, C_\varphi/T$  [high-T regime]

Thermal corrections may spoil inflation!  $\Delta V_T = -\frac{\pi^2}{90} g_R T^4 + \frac{g^2 \varphi^2}{12} T^2 + \dots$

Berera, Gleiser & Ramos PRD'98; Yokoyama & Linde PRD '98

**Solution:** use symmetries to protect the inflaton mass

Inflaton a PNGB of a broken U(1) symmetry + pair of scalars + exchange sym.

[BG, Berera, Ramos & Rosa 2016]

[BG, Berera, Ramos & Rosa 2019]

# Scalar Little Inflaton Model

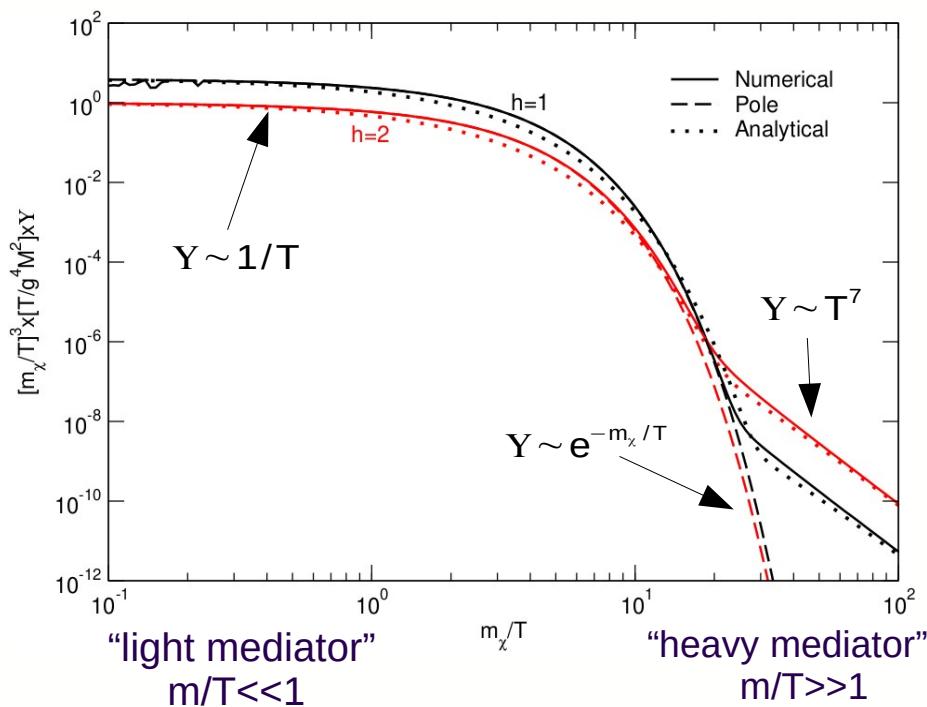
$$L = \dots - g^2 M^2 \cos^2(\varphi/M) \chi_1^2 + g^2 M^2 \sin^2(\varphi/M) \chi_2^2 + h \sigma \sum_{i=1,2} \chi_i \bar{\psi}_j \psi_k + \dots$$

No thermal mass for the inflaton:  $\Delta V_T = -\frac{\pi^2}{90} g_R T^4 + \frac{g^2 M^2}{12} T^2 + \frac{g^4(\varphi) M^4}{16\pi^2} (\log \frac{u^2}{T^2} - c_s)$

no inflaton thermal mass

[  $M$  = symmetry breaking scale]

**Dissipative coefficient:**



$$Y \approx \frac{4g^2}{h^2} \frac{g^2 M^2}{T} F[m_\chi/T] \begin{cases} M^2/T & \text{"High-T"} \\ T^7/M^6 & \text{"Low-T"} \end{cases}$$

During inflation  $T$  always decreases:

$$gM/T \ll 1 \quad \xrightarrow{\hspace{2cm}} \quad gM/T \gg 1$$

"High-T"

"Low-T"

"No growing mode"

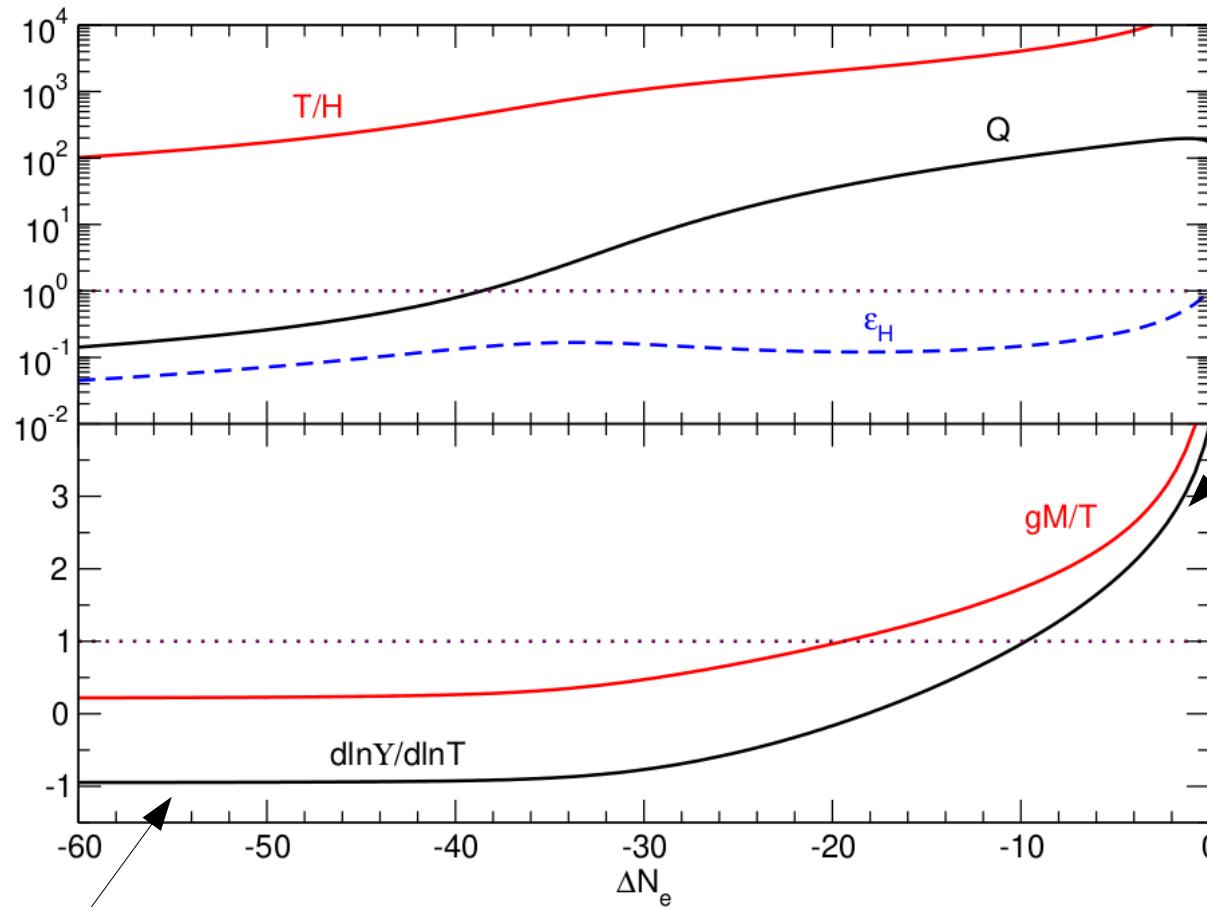
"growing mode"

$$[m_\chi^2 = \frac{g^2 M^2}{2} + \frac{h^2 T^2}{12} + \frac{h_s^2 T^2}{12}]$$

# Scalar Little Inflaton Model + quartic chaotic

$$V = \lambda \varphi^4 / 4$$

$[M = 10^{-4} m_P, g = 1, h = 2.5, \lambda = 10^{-14}]$



CMB scales: no growing mode

$$Y \simeq M^2 / T$$

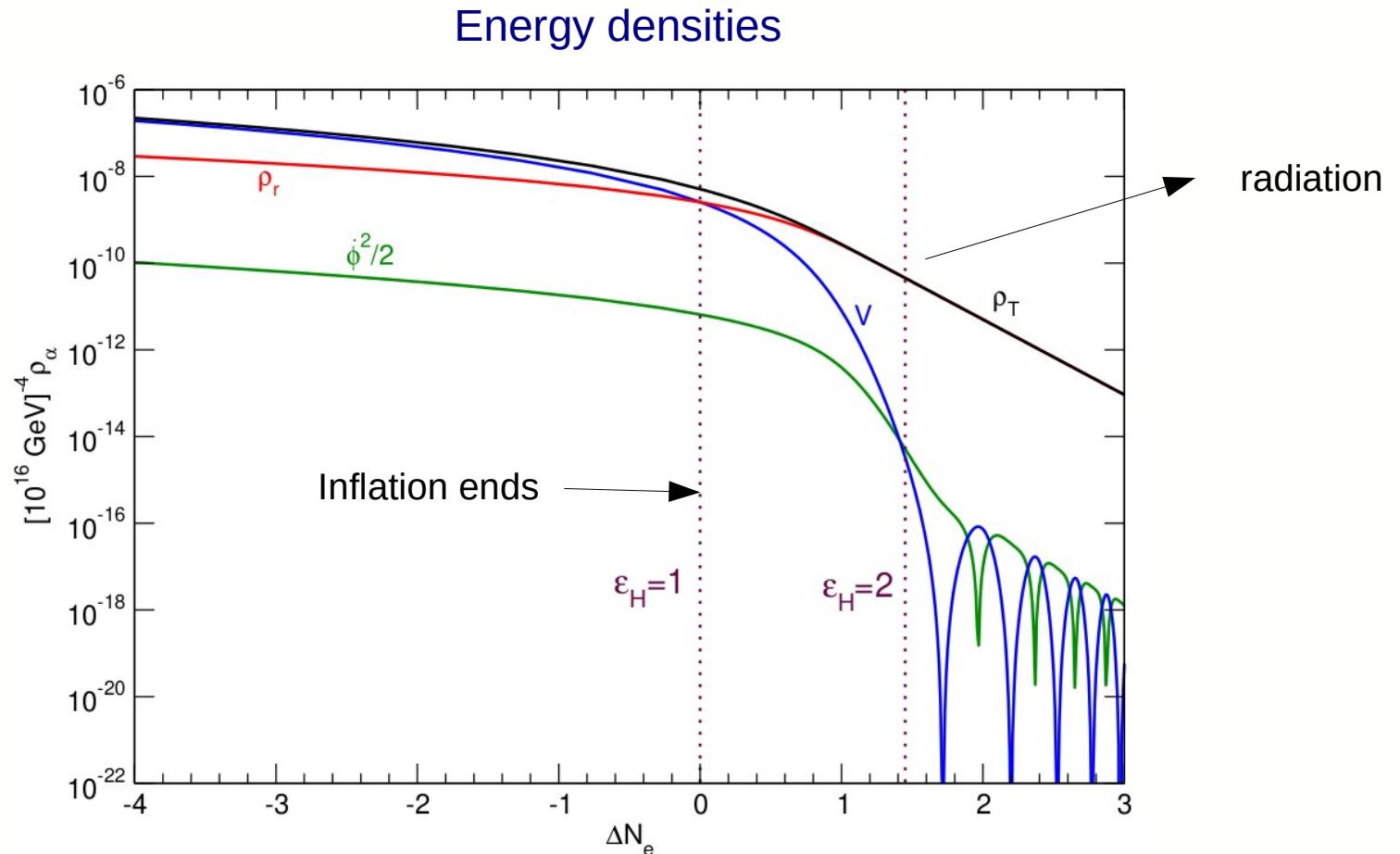
End of inflation:  
amplification of the  
primordial spectrum

$$Y \simeq T^3 / M^2$$

# Scalar Little Inflaton Model + quartic chaotic

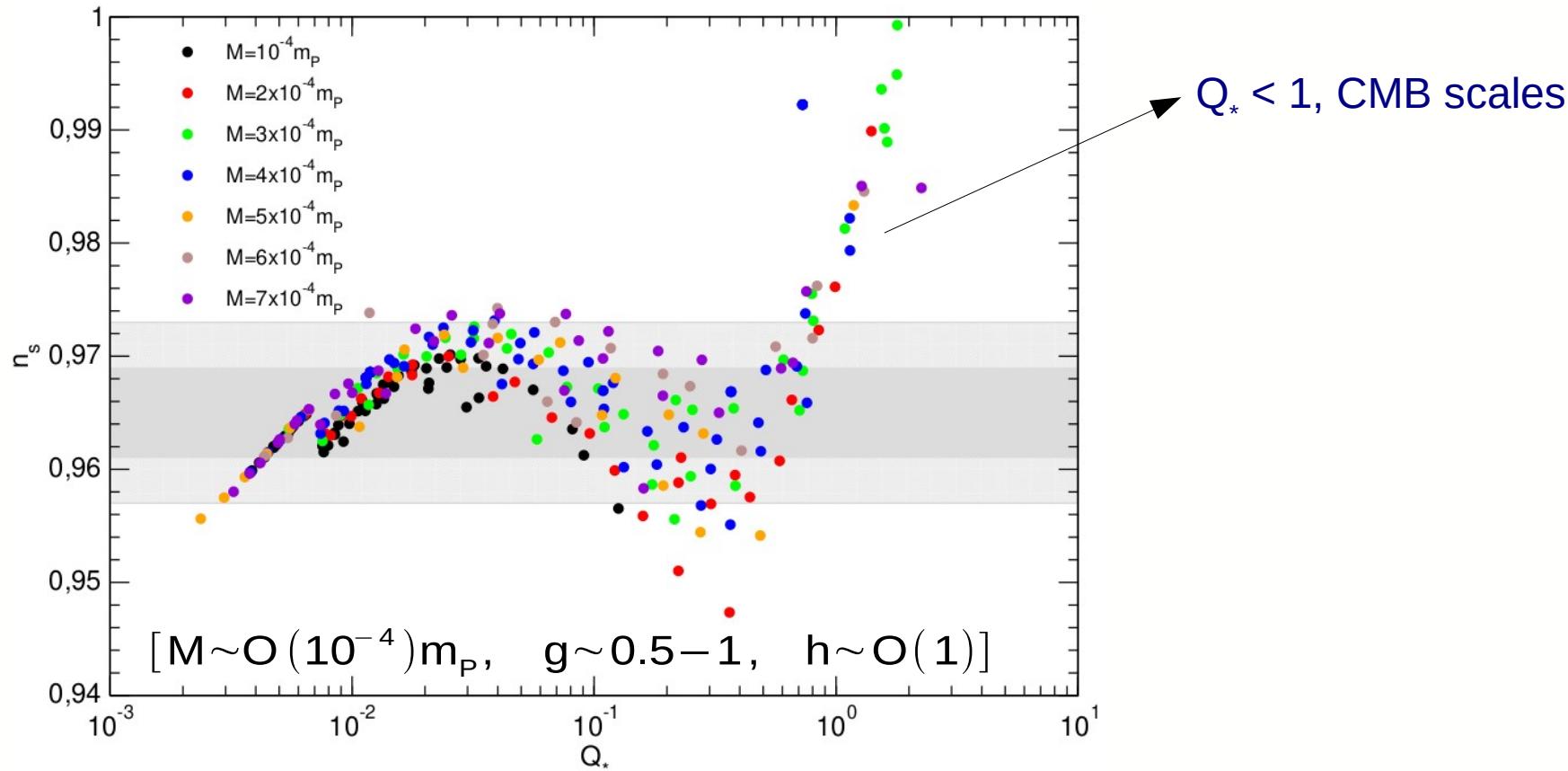
$$V = \lambda \varphi^4 / 4$$

$[M = 10^{-4} m_P, g = 1, h = 2.5, \lambda = 10^{-14}]$



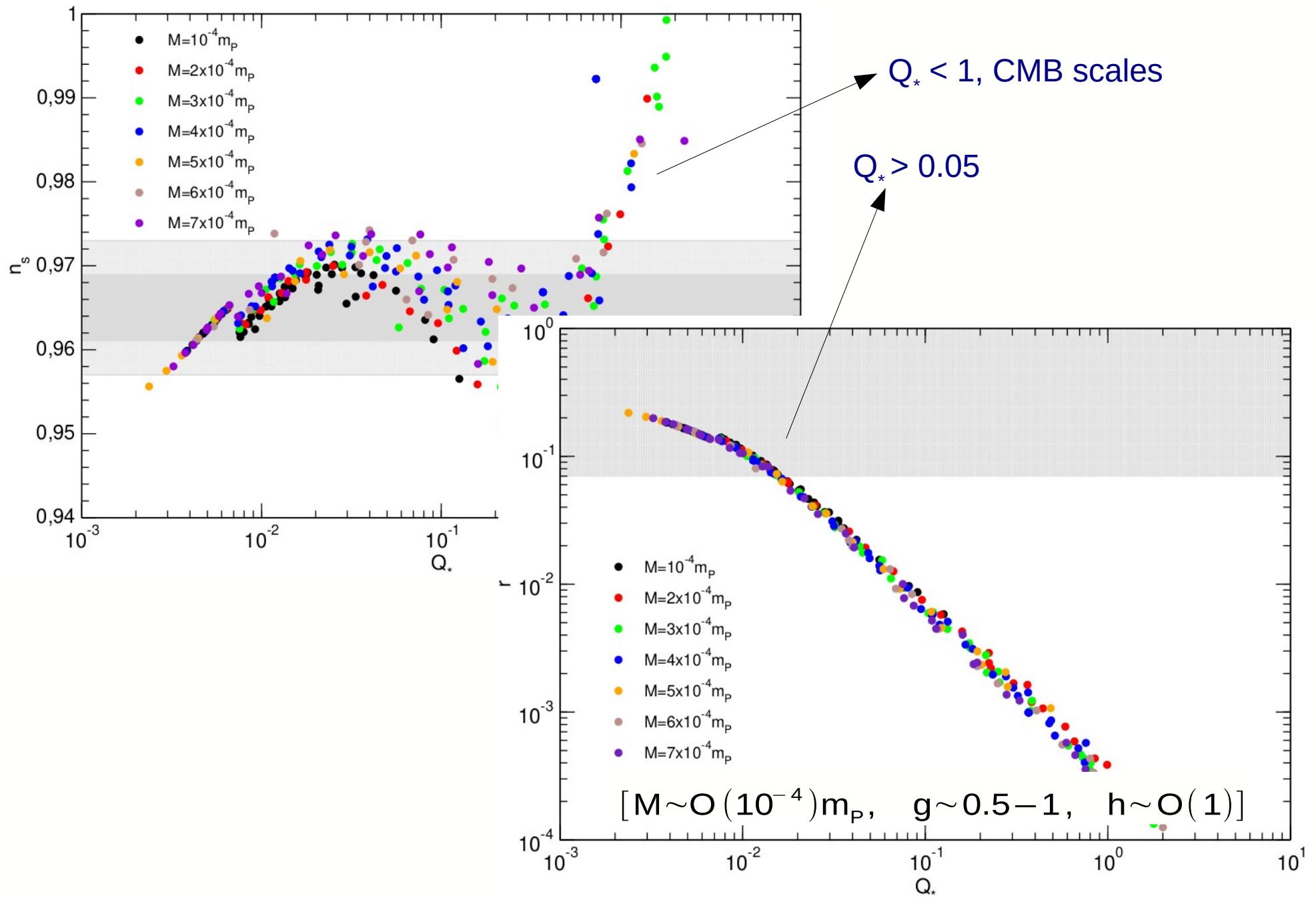
Smooth transition from inflation to radiation

# Spectral index & tensor-to-scalar ratio

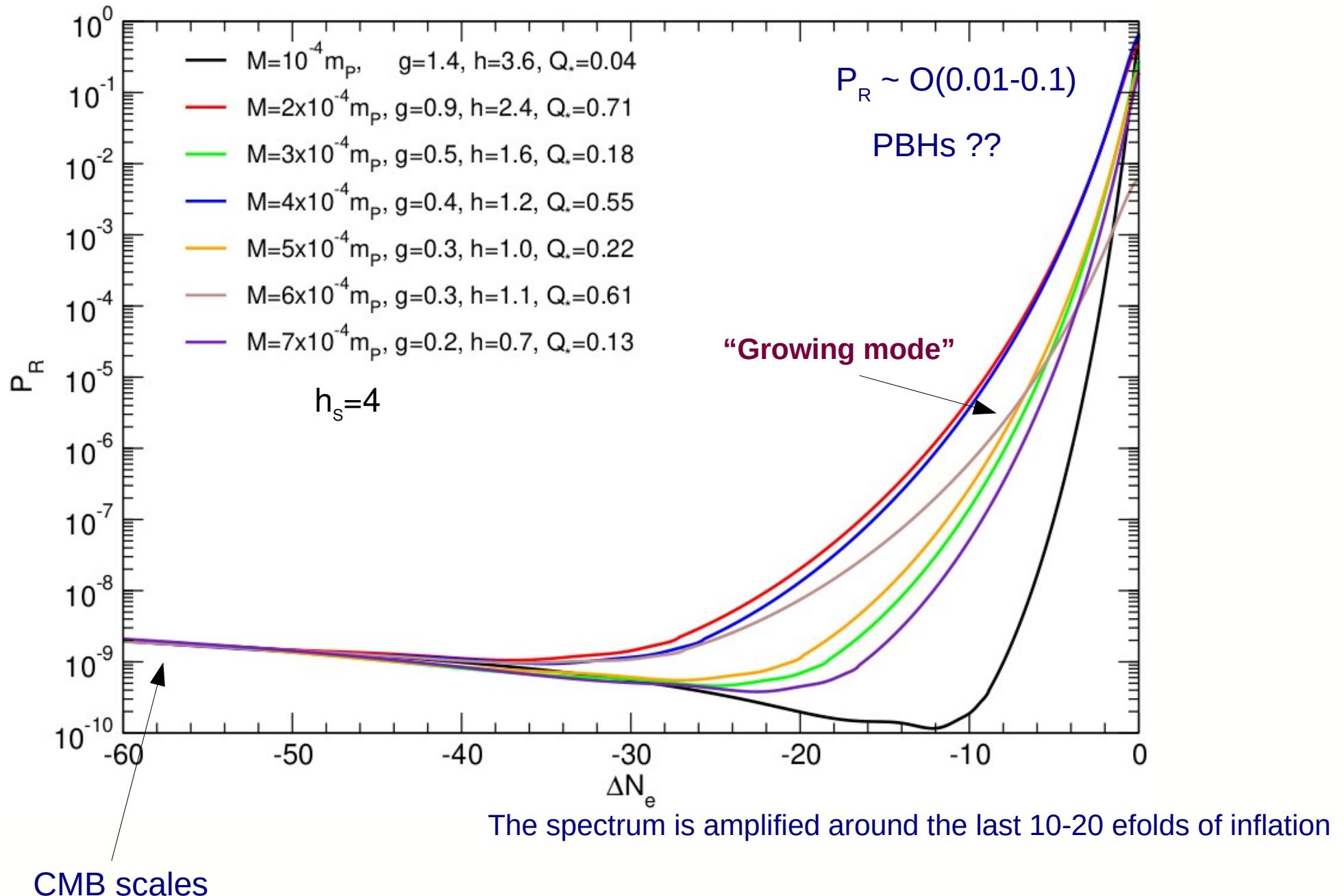


[  $\lambda$  fixed using Planck normalization of the primordial spectrum ]

# Spectral index & tensor-to-scalar ratio



# Primordial scalar spectrum



# Primordial Black Holes

$P_R \sim O(0.01-0.1)$  will lead to PBH formation on re-entry

$$M_{\text{PBH}}(k) \sim \frac{4\pi m_P^2}{H_M}$$

H on re-entry, close to the end of inflation

$M_{\text{PBH}} \sim [5 \times 10^4 \text{ g}, 10^6 \text{ g}]$  Light, evaporating black hole

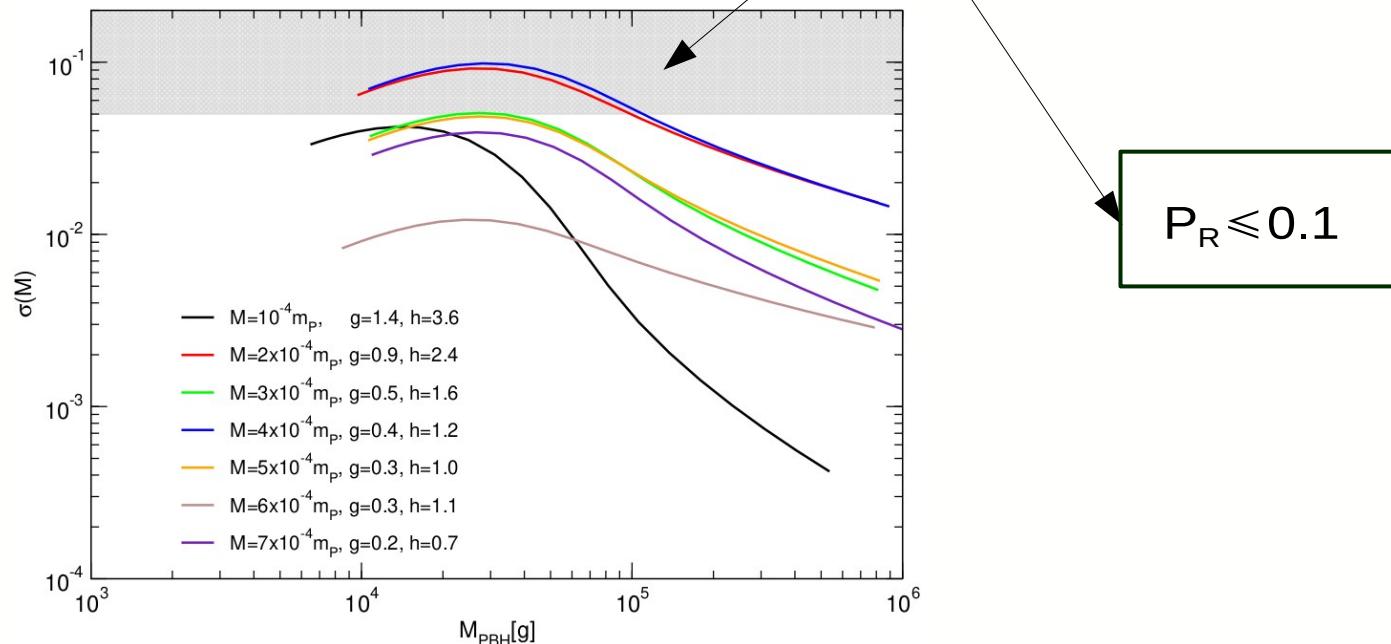
Planck relics may be produced during their evaporation

To avoid the overclosure of the universe:

$\sigma(M) < 0.05$

Variance of the density fluctuations

[Carr et al. PRD81 '10]



# Induced 2<sup>nd</sup> order GW

Although at linear order scalar, vector and tensor perturbations decouple,  
large scalar fluctuations source tensors at second order

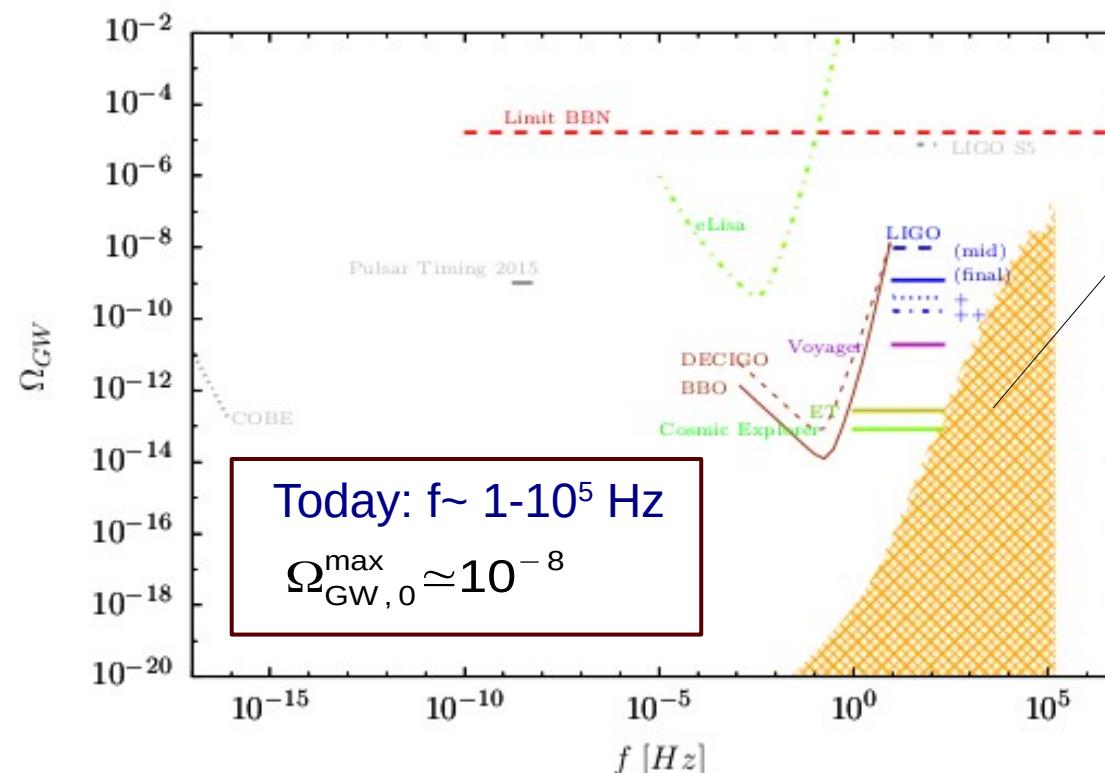
$$\ddot{h}_k + 3H\dot{h}_k + \frac{k^2}{a^2}h_k = S_k[\Phi_k]$$

Primordial spectrum

Gravitational potential:  $\Phi_k = T[k\tau]\phi_k, \quad \langle\phi_k\phi_q\rangle = \delta(\vec{k}+\vec{q}) \frac{2\pi^2}{k^3} \left(\frac{3+3w}{5+3w}\right)^2 P_\zeta(k)$  [w=1/3 radiation]

GW spectral density:  $\Omega_{GW,0}(k) \simeq 0.4 \Omega_{r,0} \times \frac{1}{24} \left(\frac{k}{aH}\right)^2 P_h(k, \tau_c)$

[Kohri & Terada 1804.08577]



# Summary

- Dissipative effects due to decaying fields can be relevant during inflation, and modify the inflationary predictions
- For a  $T$  dependent dissipative coefficient, the field and radiation perturbation EOM form a coupled system: field fluctuations are amplified before freeze-out
- “High  $T$ ” regime at CMB scales, compatible with observations:  $Y = C_T / T$

Inflaton a PNGB of a broken U(1) symmetry + pair of scalars + exchange sym.

Light scalars:  $gM < T$  + thermal corrections under control + minimal matter content

$\lambda\phi^4$  compatible with data,  $Q_* \sim 0.01-1$ ,  $r \sim 0.1-10^{-4}$

- “Low  $T$ ” regime before the end of inflation:  $Y = C_T T^c \rightarrow$  amplification of the primordial spectrum  $P_R \sim O(0.01-0.1)$  (last 10-20 efolds)

- On re-entry:

- Light, evaporating PBHs:  $M_{PBH} \sim O(10^4-10^6) g$
- Primordial GW:  $f \sim O(1-10^5) Hz$ ,  $\Omega_{GW,0}^{\max} \simeq 10^{-8}$

