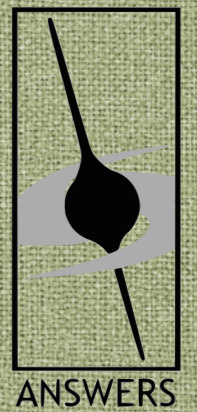




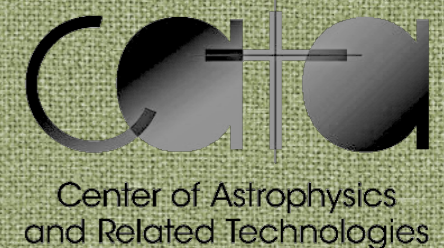
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Magnetic field evolution in neutron star cores

Francisco Castillo, A. Reisenegger, J. A. Valdivia

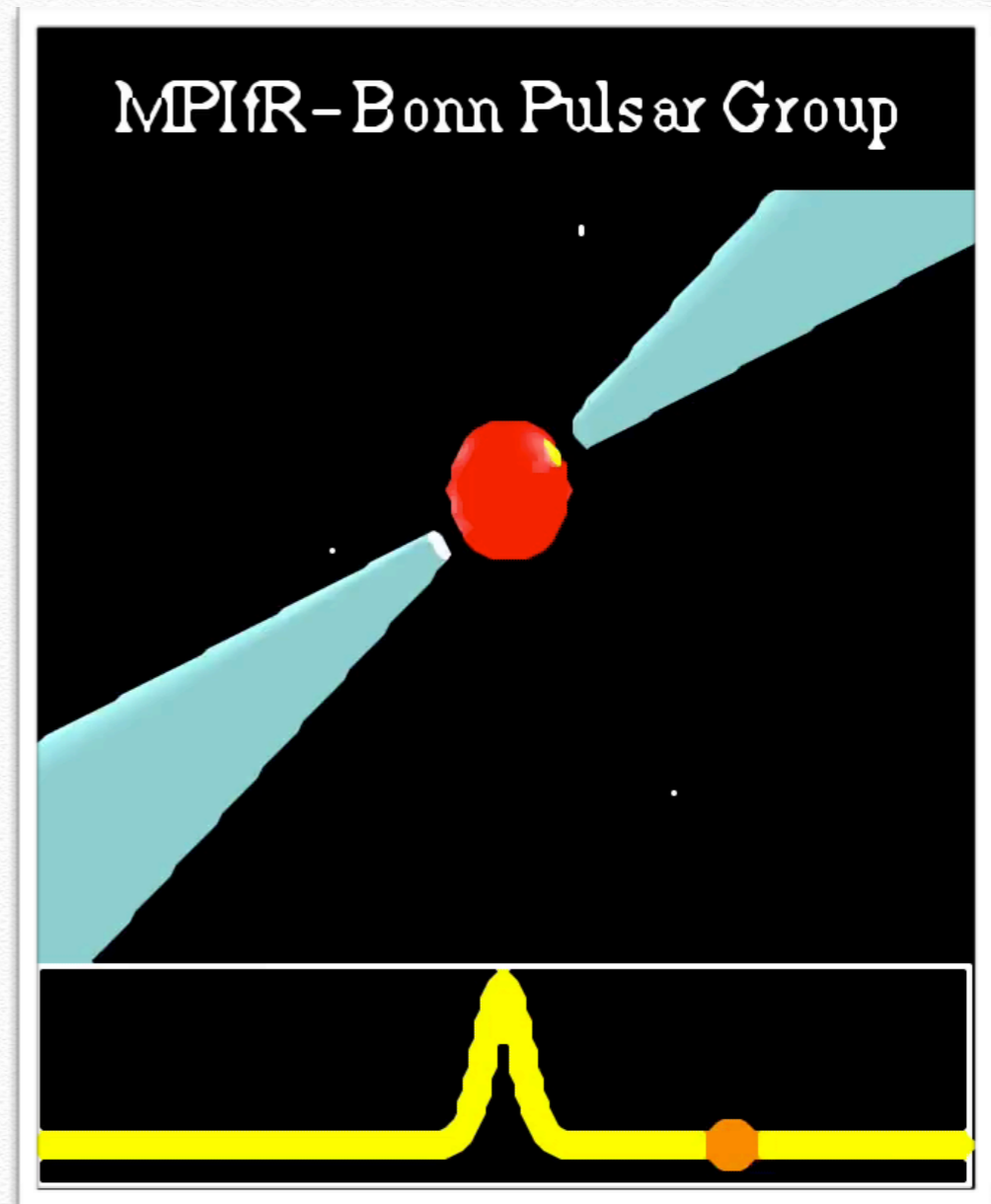
CoCoNuT Meeting 2018



Evidence for evolution?

$$B = 3.2 \times 10^{19} \sqrt{P\dot{P}/s} \text{ G.}$$

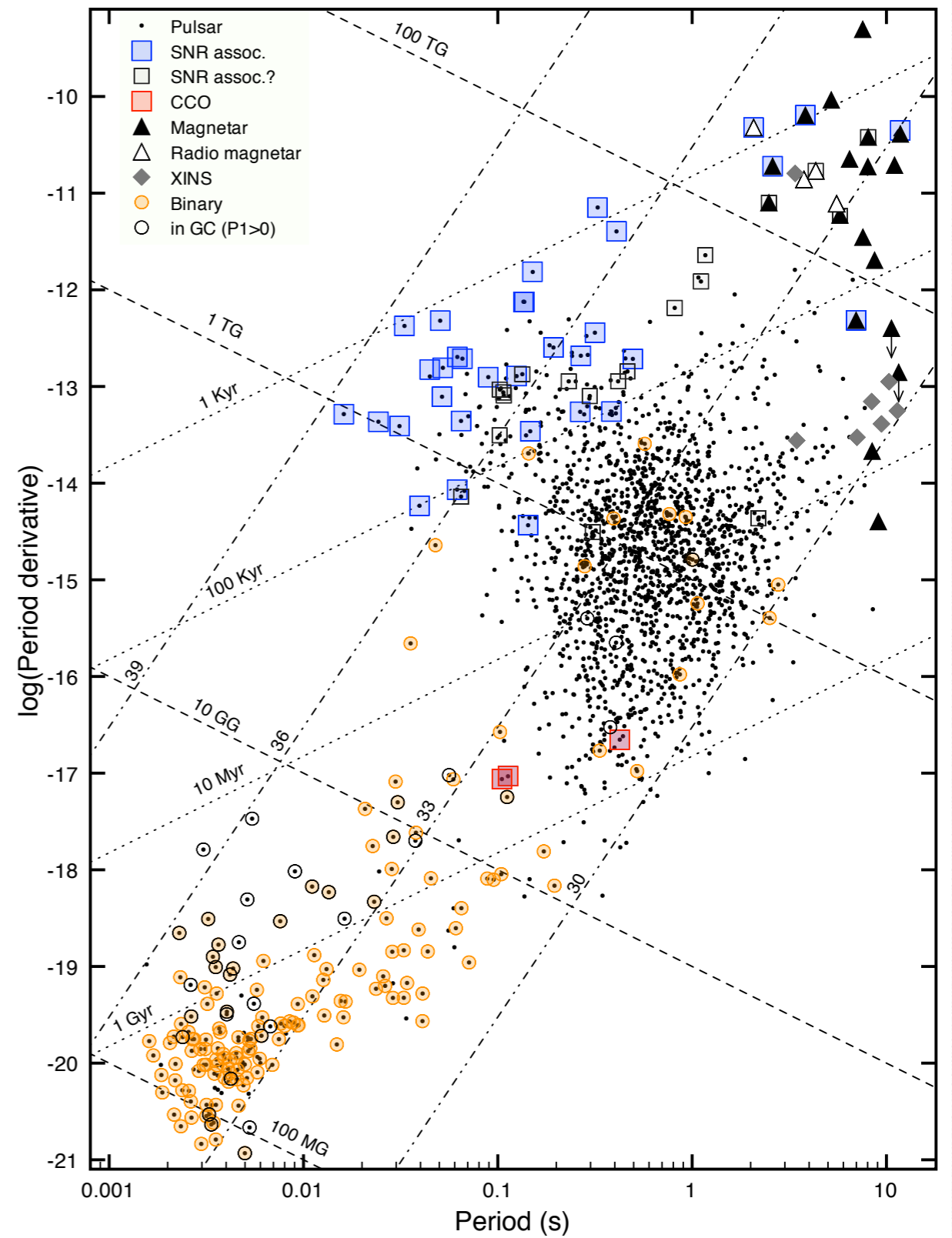
$$\tau = \frac{P}{2\dot{P}}.$$



Evidence for evolution?

$$B = 3.2 \times 10^{19} \sqrt{P\dot{P}/s} \text{ G.}$$

$$\tau = \frac{P}{2\dot{P}}.$$



How does B evolve?

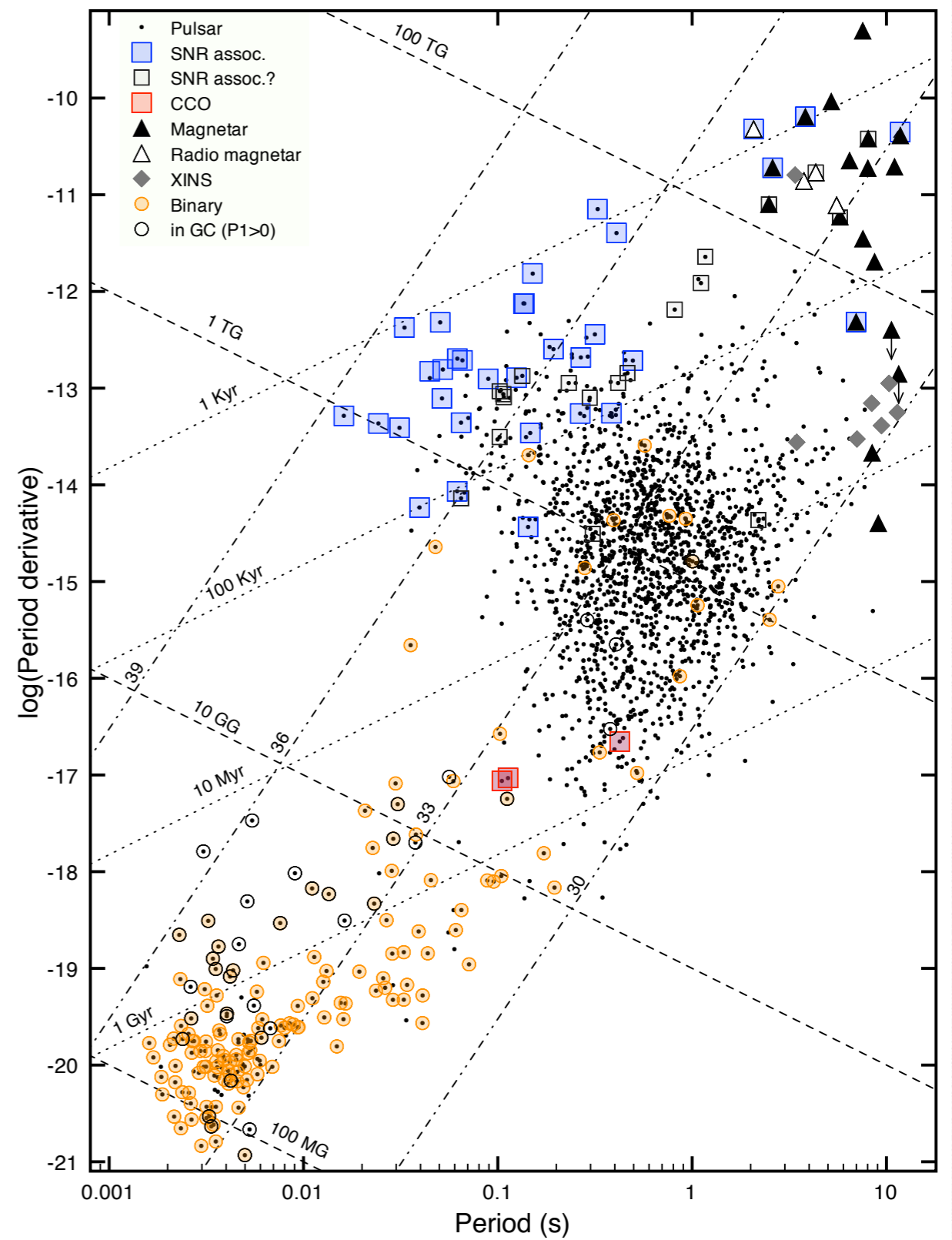


Figure: C. Espinoza.

Neutron star structure

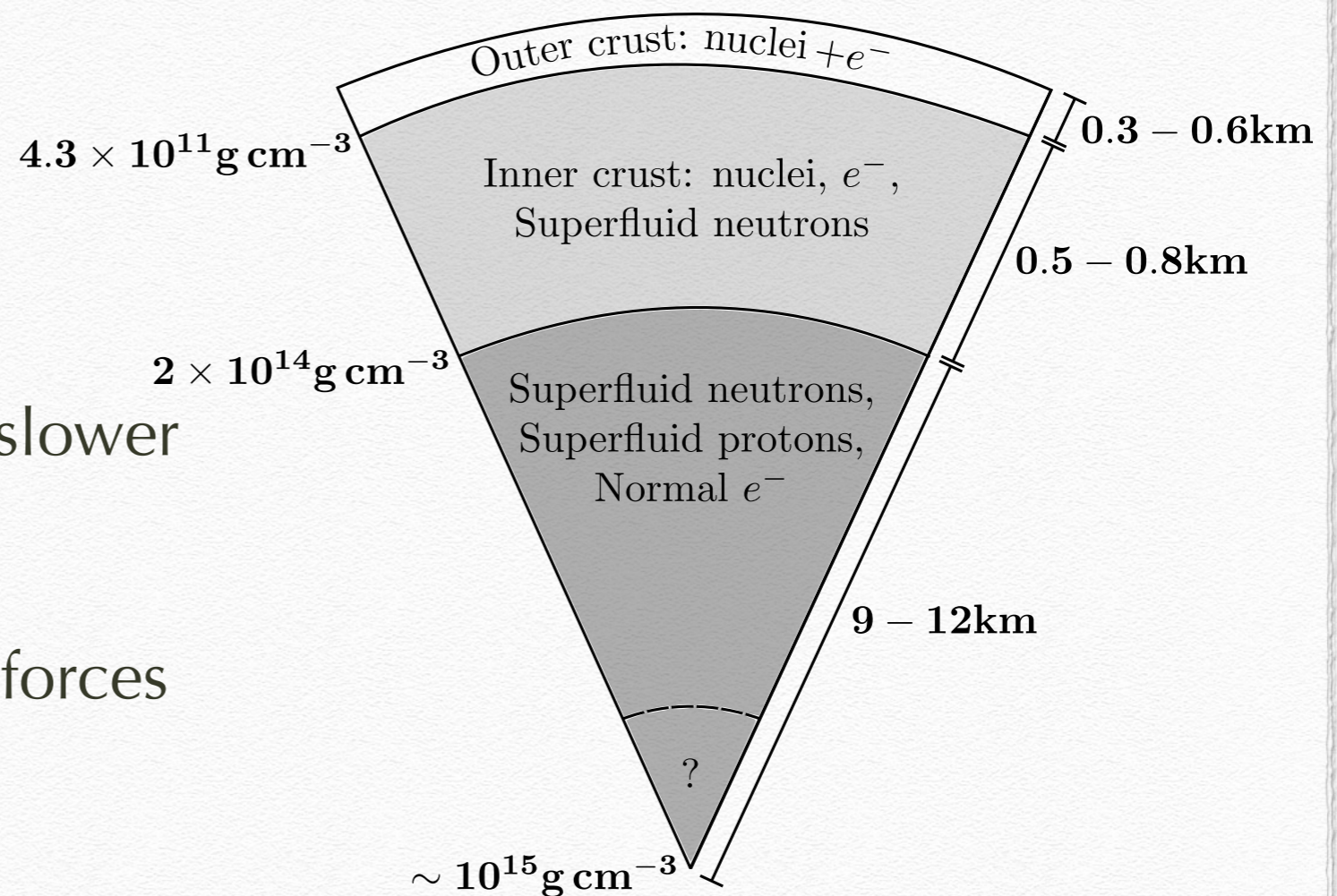
❖ Crust

❖ Core

❖ High conductivity:
Hall+Ohm: Much slower

❖ All particles
can move: friction forces

❖ Weak interactions



B evolution: Core vs. Crust

$$t_{\text{Ohm}} \sim \frac{5.7 \times 10^6}{Q} \text{ yr} \quad t_{\text{ambip}} \sim 3 \times 10^7 \frac{L_6^2 T_6^2}{B_{12}^2} \text{ yr}$$

**We treat the crust as a
vacuum**

(whose magnetic field at any time is fully determined by the field in the core).

What happens in the core?

Formalism: Goldreich & Reisenegger (1992); Hoyos et al. (2008)

❖ *Charge neutrality.*

❖ *Uniform background in hydrostatic and chemical equilibrium*

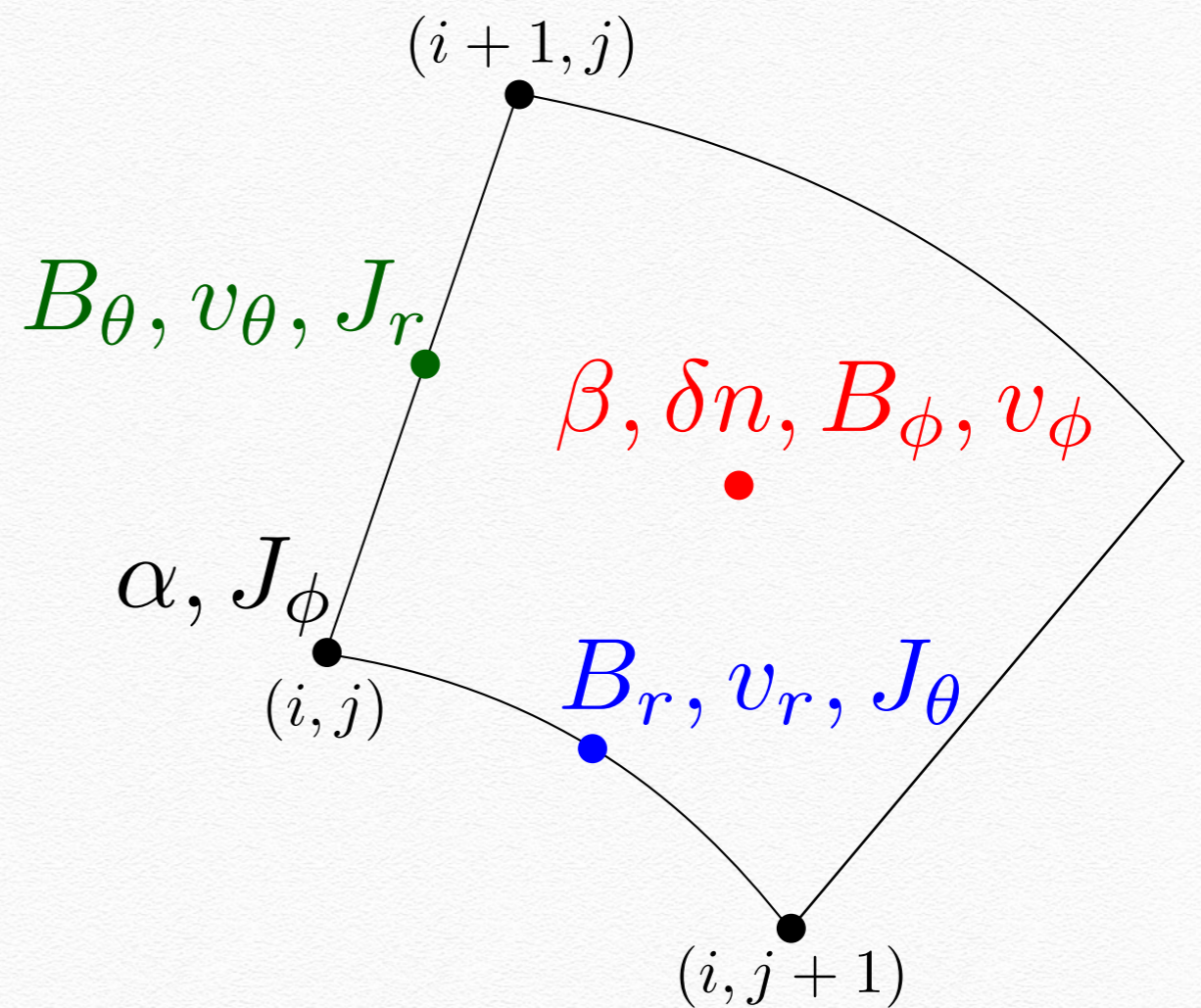
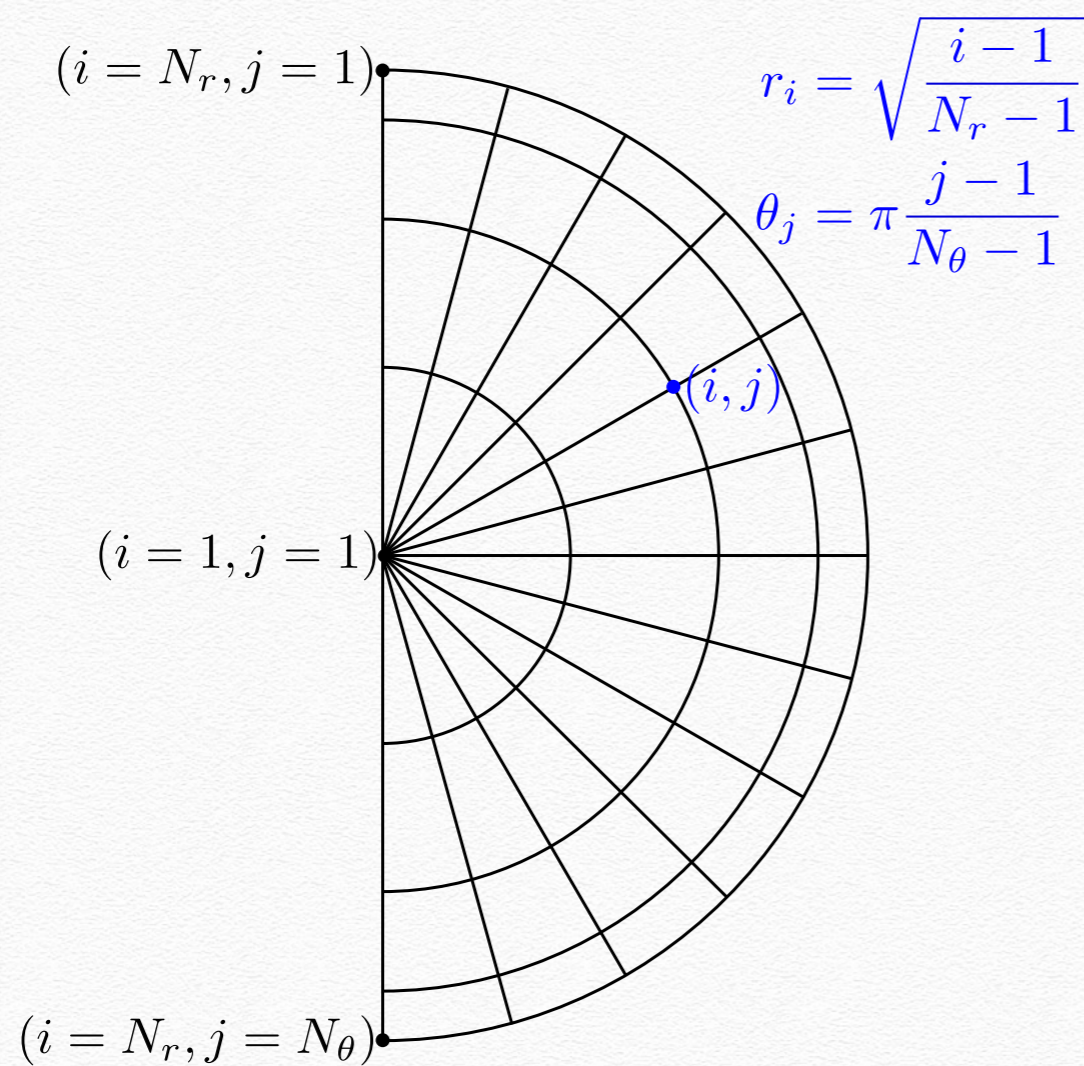
$$\nabla \mu + \frac{\mu}{c^2} \nabla \Psi = 0.$$

❖ Ratio between the magnetic pressure and degeneracy pressure: $B^2/8\pi P \lesssim 10^{-6}$

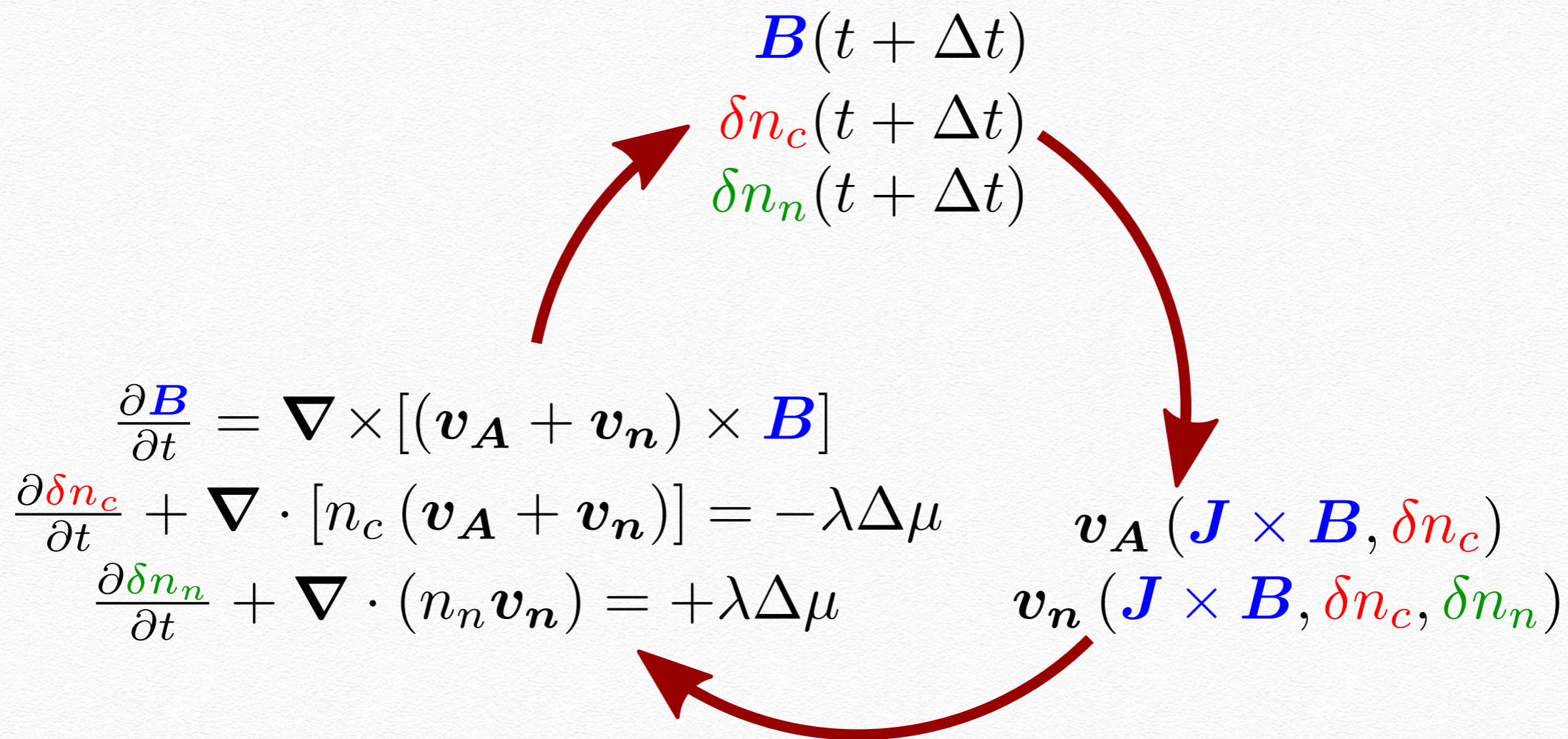
$$n_j(\vec{r}, t) \rightarrow n_j + \delta n_j(\vec{r}, t) \quad j = n, c$$

$$\mu_j(\vec{r}, t) \rightarrow \mu + \delta \mu_j(\vec{r}, t) \quad \mu = \mu_c = \mu_n$$

We need numerical simulations

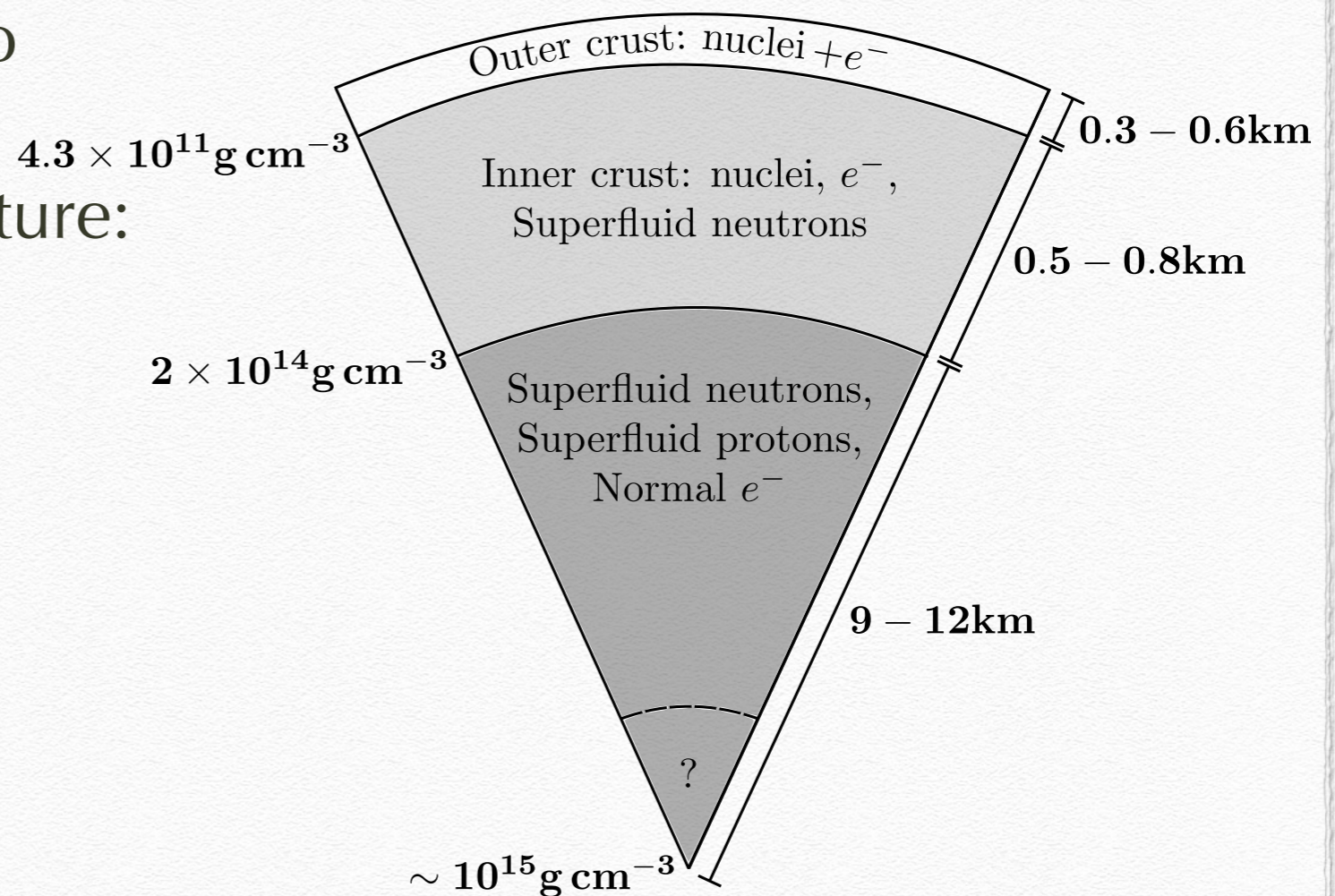


We need numerical simulations



What happens in the core?

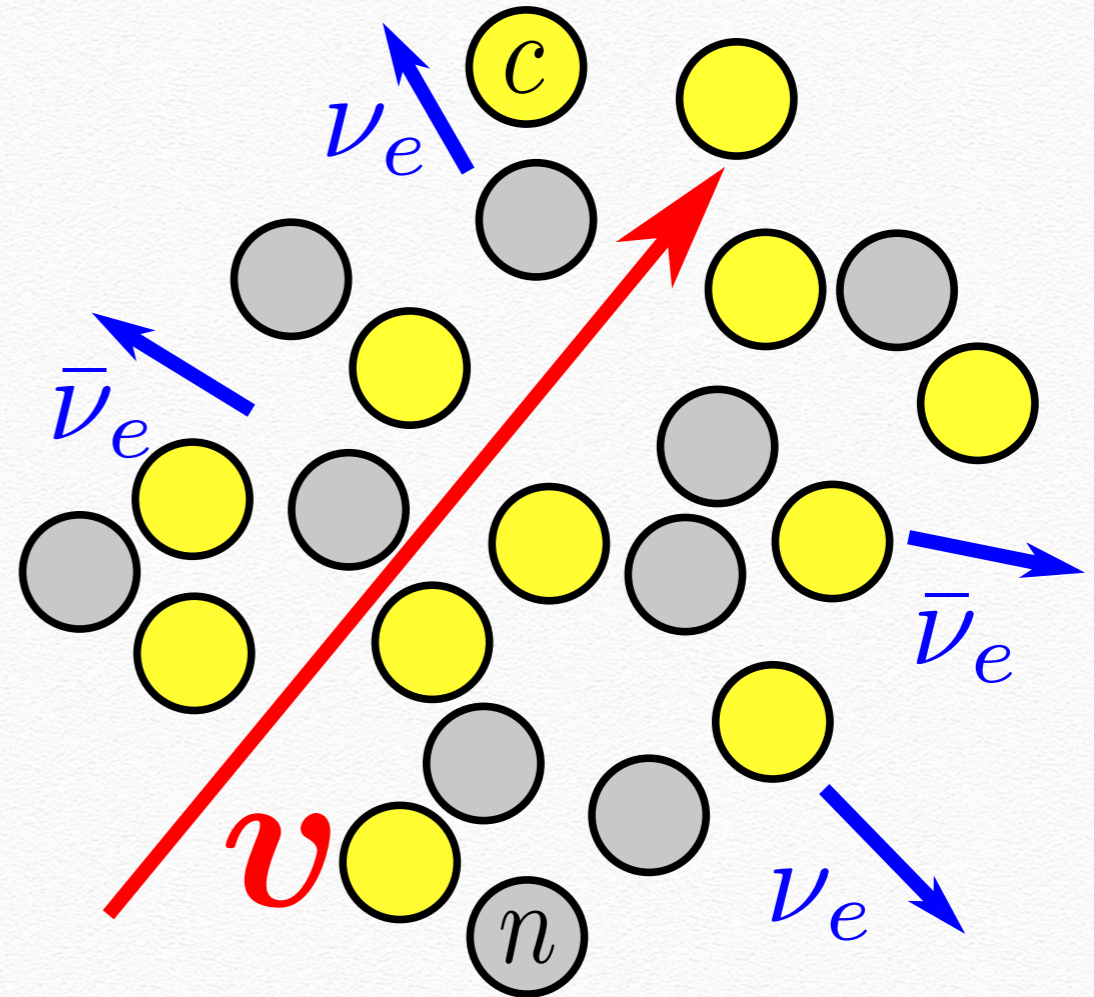
We can distinguish two
different regimes
depending on temperature:



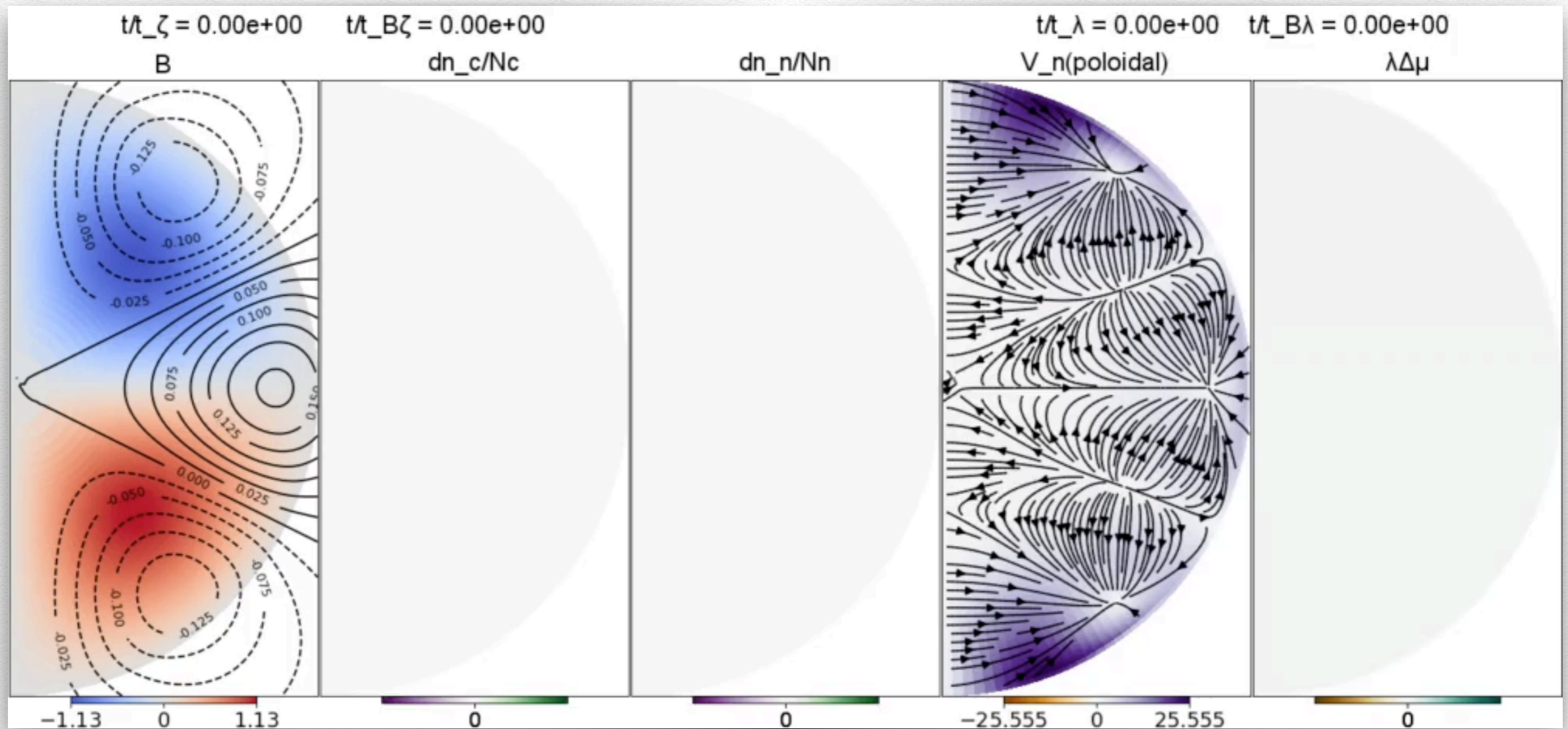
What happens in the core?

We can distinguish two different regimes depending on temperature:

- ❖ **High temperature:**
strong coupling regime
- ❖ Joint radial motions:
 - ❖ Opposed by strong buoyancy forces
 - ❖ Feasible only during a very short time after the neutron star birth

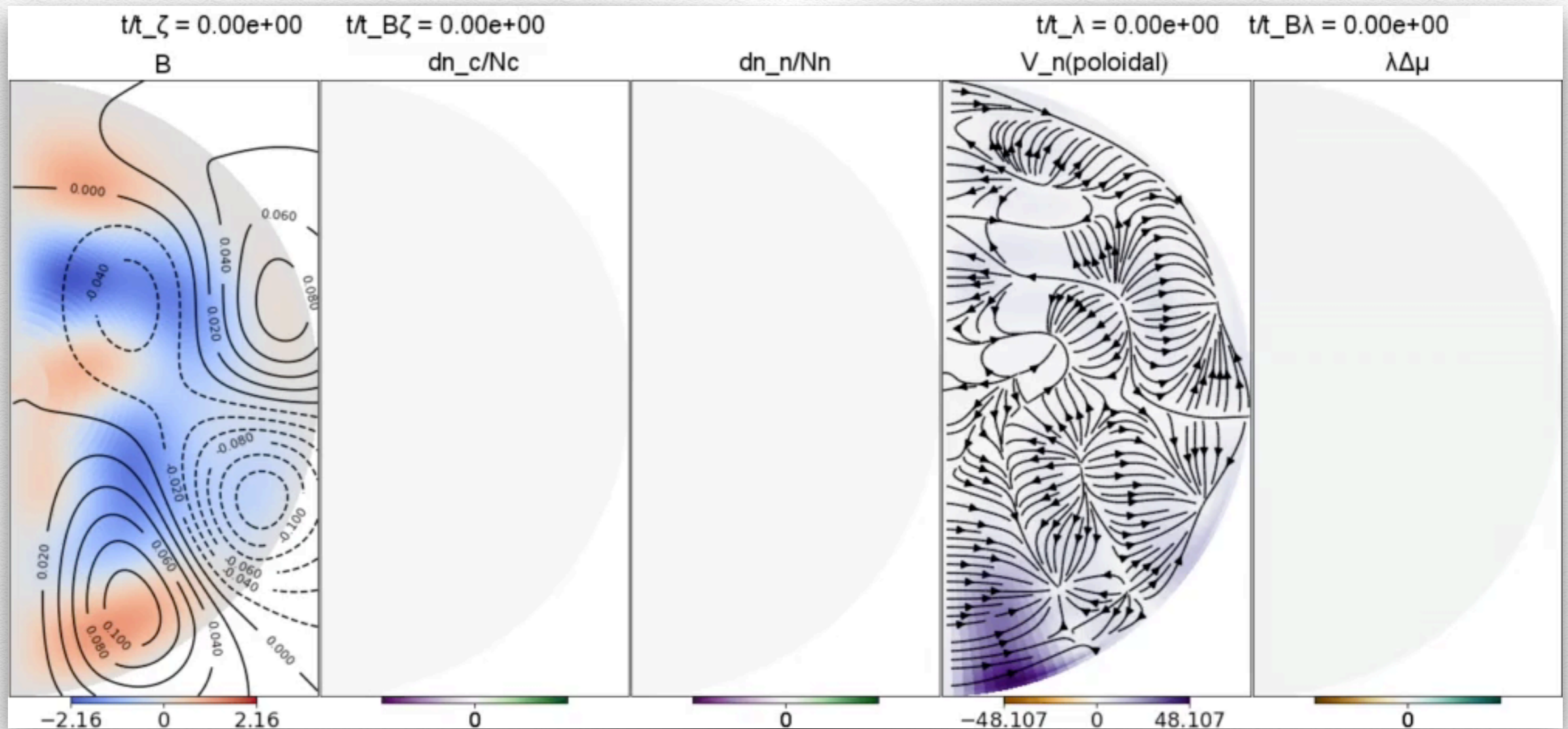


Strong coupling regime



Simulation performed taking $t_\zeta : t_{B\zeta} : t_\lambda : t_{B\lambda} = 1 : 100 : 640 : 64000$

Strong coupling regime

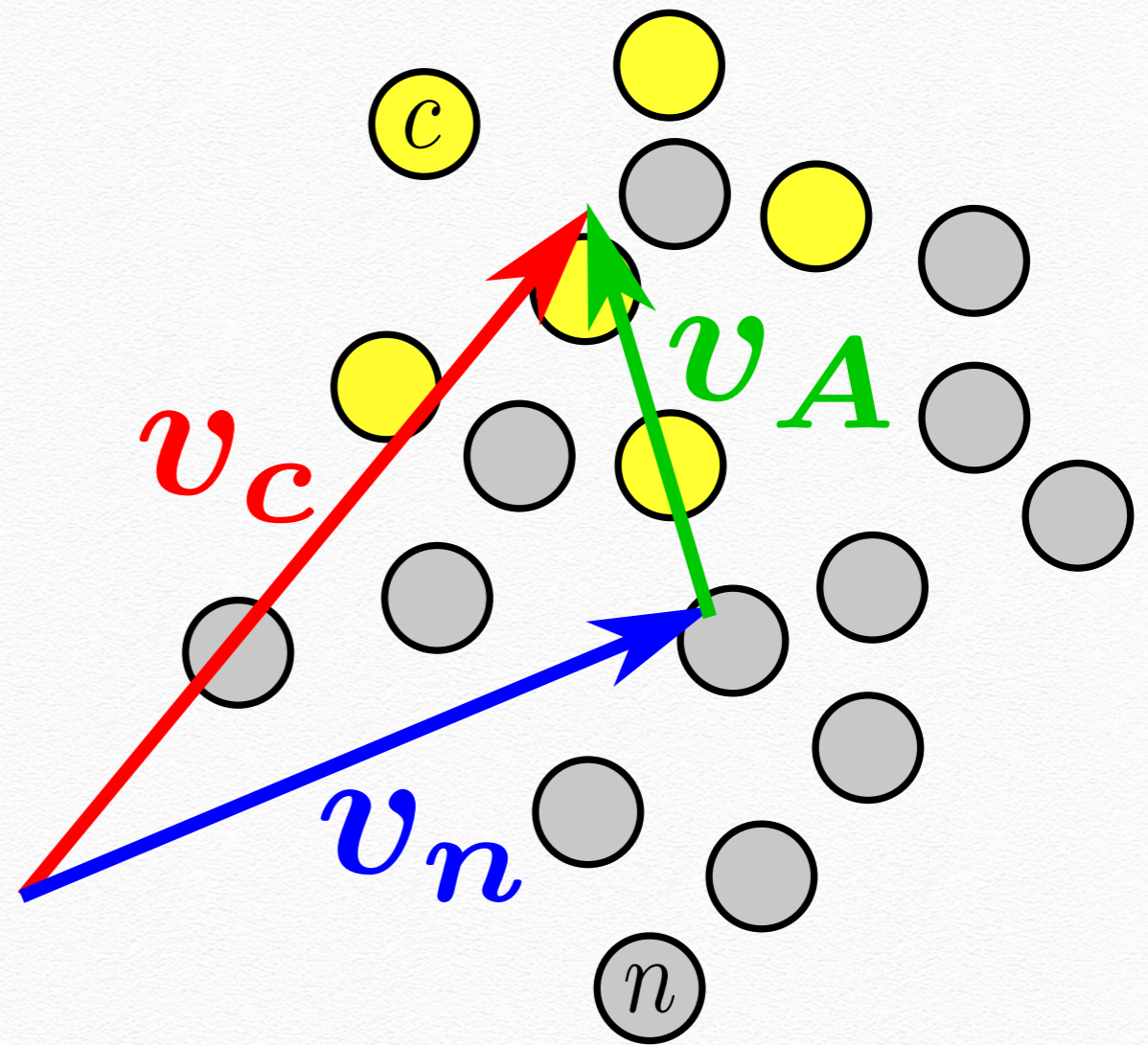


Simulation performed taking $t_\zeta : t_{B\zeta} : t_\lambda : t_{B\lambda} = 1 : 100 : 640 : 64000$

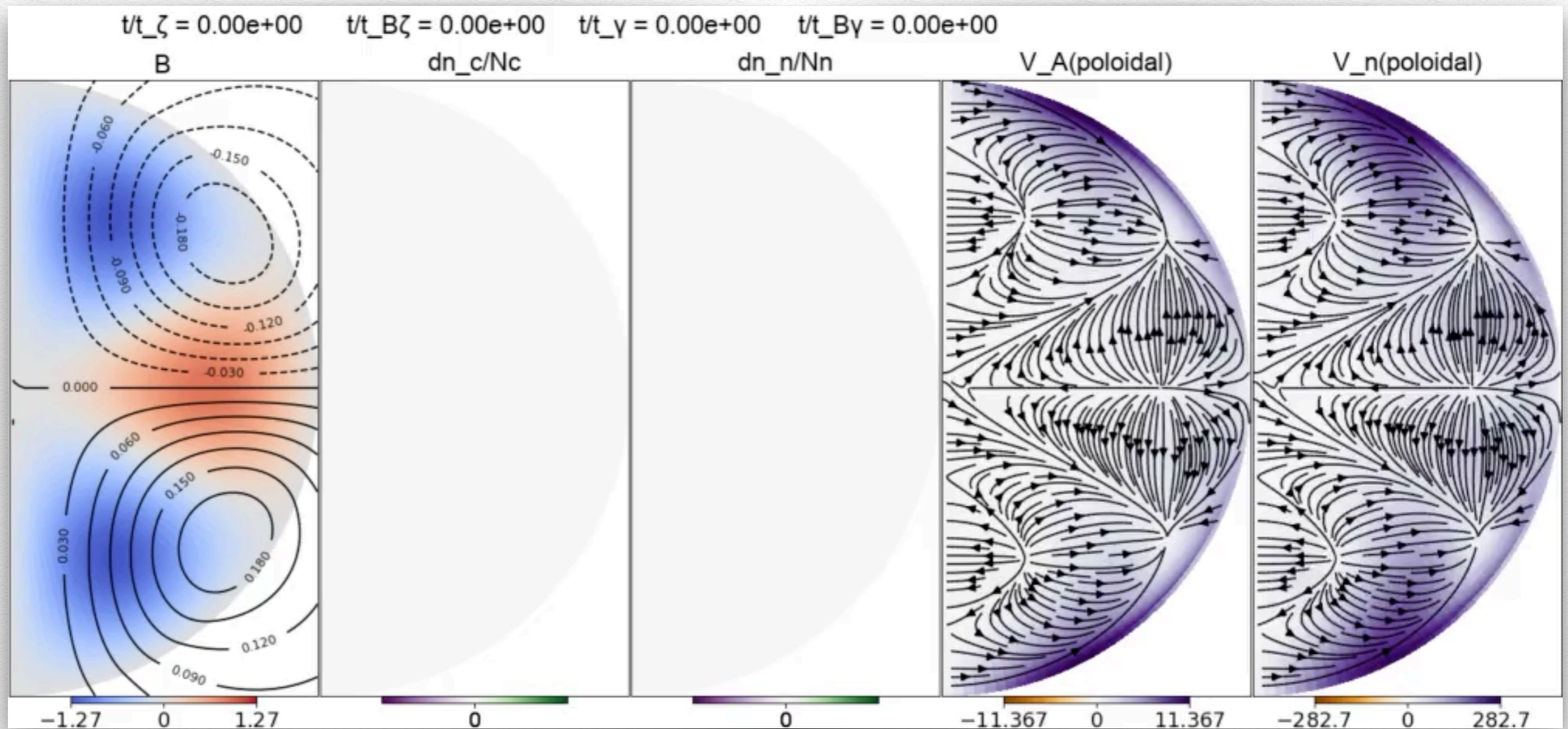
What happens in the core?

We can distinguish two different regimes depending on temperature:

- ❖ **Lower temperatures:**
weak coupling regime
 - ❖ Weak interactions are strongly suppressed
 - ❖ Collisional coupling is reduced: Relative motions
 - ❖ Ambipolar diffusion



Weak coupling regime



Simulation performed taking $t_\zeta : t_{B\zeta} : t_\gamma : t_{B\gamma} = 1 : 500 : 2000 : 57290$

Summary and Conclusions

- ❖ Simulations that evolve simultaneously the magnetic field and the density perturbations it induces for neutrons and charged particles.
- ❖ Two regimes of magnetic field evolution:
 - ❖ Strong coupling
 - ❖ Urca reactions can adjust the composition “in real time”.
 - ❖ Matter behaves as a single fluid (with time-varying composition) which moves together with the magnetic field.
 - ❖ Weak coupling
 - ❖ Relative motion of neutrons and charged particles: Ambipolar diffusion
- ❖ The star evolves towards hydromagnetic equilibria

Summary and Conclusions

- ❖ Caveats:
 - ❖ Neglected the currents in the crust (assumed to have a very low conductivity).
 - ❖ Superfluidity/superconductivity.

Summary and Conclusions

- ❖ Previous work indicates that there are no stable hydromagnetic equilibria in barotropic stars:
 - ❖ Resulting equilibria are likely to become unstable if considered in full 3D.
 - ❖ **Currently:** I aim to develop a 3D code:
 - ❖ Challenge: Integration time
 - ❖ **Any input appreciated**