

problems of horizon in effective gravity

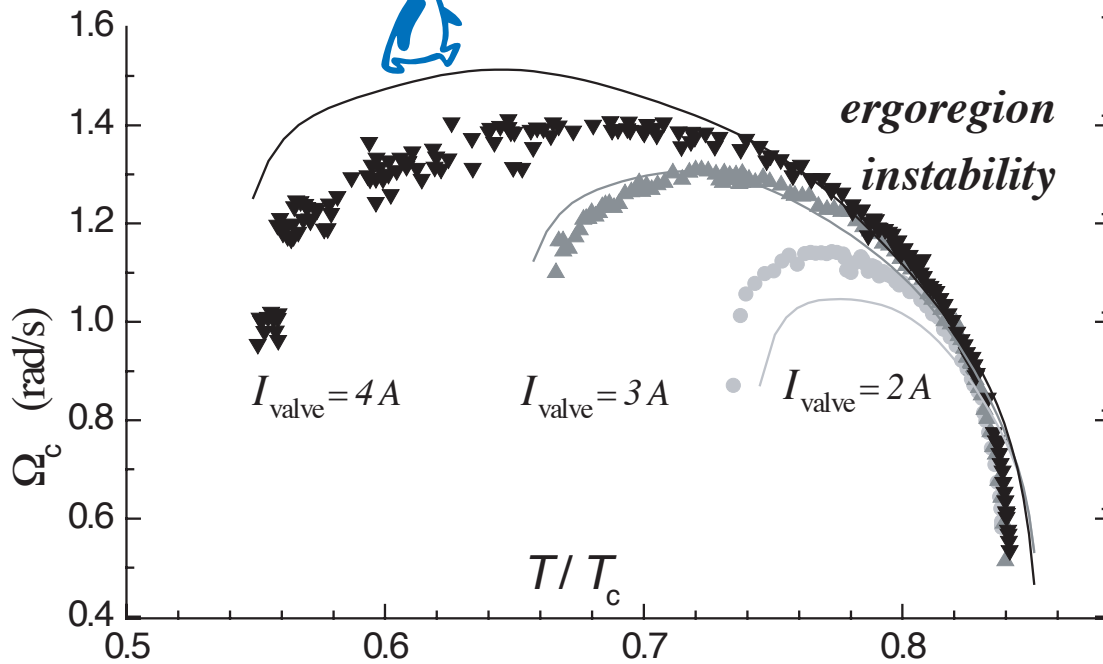
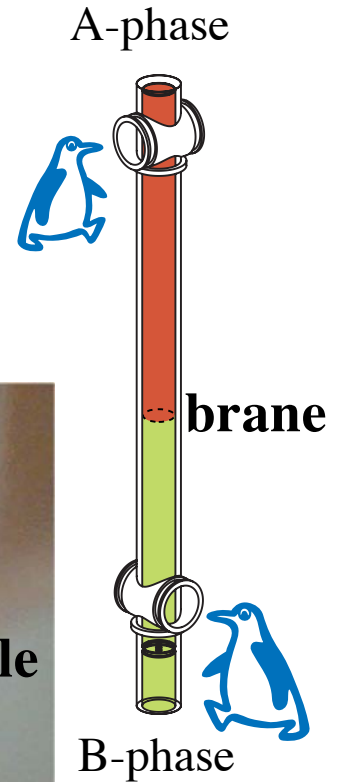
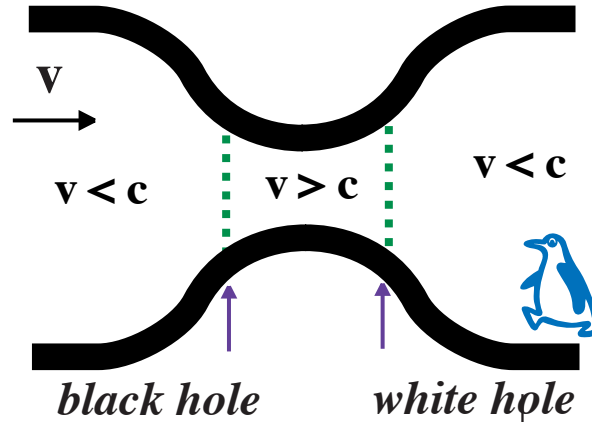
Helsinki U. Technology

G. Volovik

Landau Institute

Valencia Feb. 4, 2009

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problems of horizon in effective gravity

RUSSIAN ACADEMY OF SCIENCES

L. D. Landau
INSTITUTE FOR
THEORETICAL
PHYSICS



* sources of effective gravity in condensed matter



* black and white hole horizons

* Hawking radiation as quantum tunneling

* analog of Zeldovich-Starobinsky radiation from rotating BH

* vacuum instability in the presence of horizons & ergoregions

* from condensed matter to quantum gravity:
quantum vacuum as self-sustained Lorentz invariant medium

* possible instability of astronomical black holes



sources of effective gravity in condensed matter

*** Fermi point gravity**

*** acoustic gravity**



*** ripplon 2+1 gravity**

*** magnon BEC gravity** (*application of Cornell idea to magnon BEC*)

...

*** optical space-times**

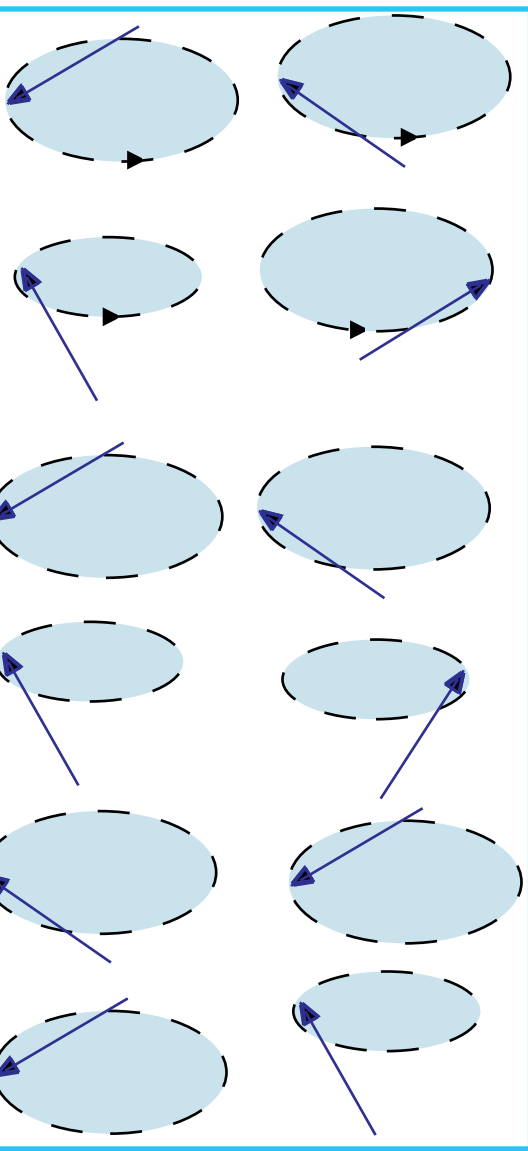
moving dielectric, optical solitons, slow light ...

*** elasticity theory of dislocations and disclinations**

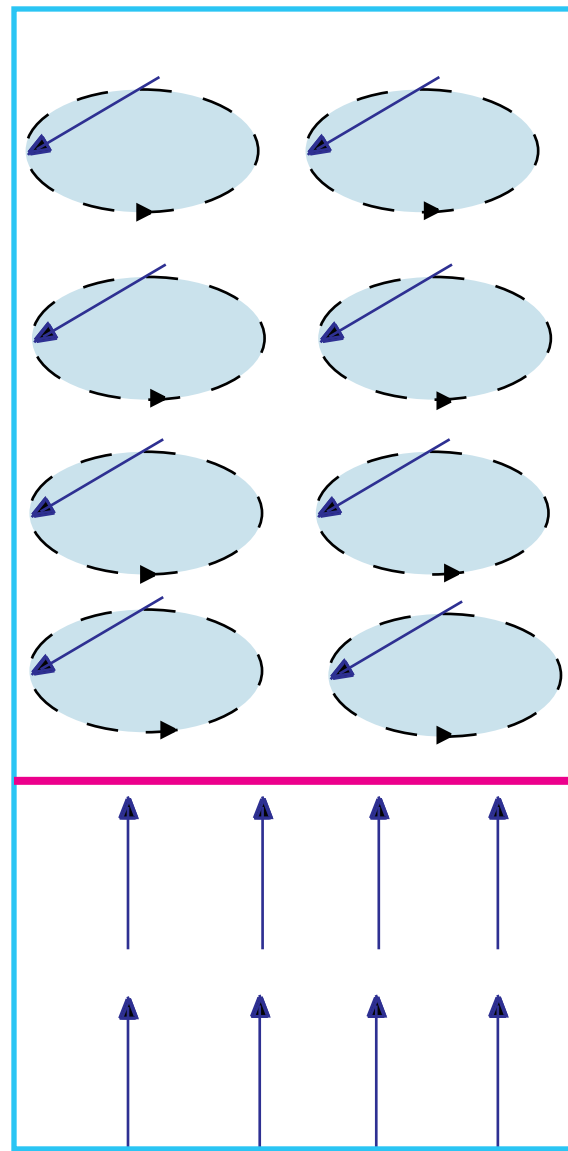
4D world crystal

Spontaneous phase-coherent precession

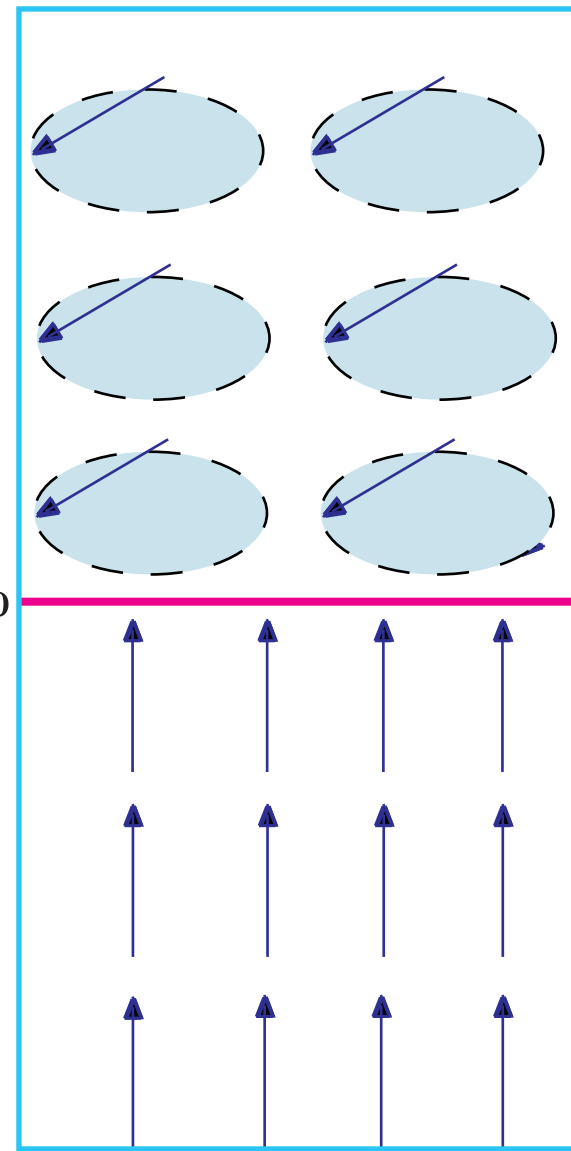
disordered precession
after pumping $S-S_z$



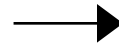
spontaneously organized
two-domain precession
with the same total spin S_z



two-domain precession decays
since S_z is not conserved,
but coherence is preserved



$\gamma H < \omega$



$\gamma H = \omega$

$\gamma H = \omega$

$\gamma H > \omega$

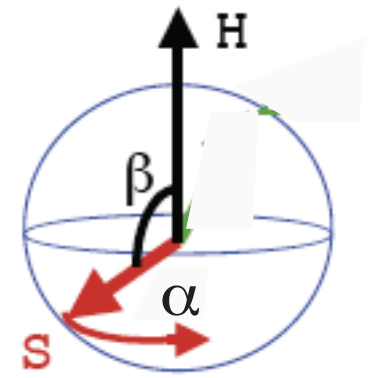
Ginzburg-Landau energy for magnon BEC

$$\Psi = |\Psi| e^{i\alpha(t)}$$

phase of precession
≡ condensate phase

$$n = |\Psi|^2 = \frac{S - S_z}{\hbar}$$

deficit of
spin projection
along field
≡ magnon number



$$\alpha(t) = \omega t + \alpha_0$$

$$F_{GL} = \frac{|\nabla\Psi|^2}{2m} + (\omega_L(\mathbf{r}) - \omega) |\Psi|^2 + F_D(|\Psi|^2)$$

magnon
mass

local
Larmor
frequency
≡
external
potential

frequency of
coherent
precession
≡
chemical
potential

spin-orbit energy
≡
interaction
between magnons

$$\dot{\alpha} = \omega$$

$$\dot{\alpha} = \mu$$

spontaneous coherent precession
=
magnon BEC

$$F_{GL} = \frac{|\nabla\Psi|^2}{2m} + (U(\mathbf{r}) - \mu) |\Psi|^2 + F(|\Psi|^2)$$



Sonic metric: effective metric in Landau two-fluid hydrodynamics

Doppler shifted phonon spectrum in moving superfluid and BEC

review:

Barcelo, Liberati & Visser,
Analogue Gravity
Living Rev. Rel. **8** (2005) 12

$$E = cp + \mathbf{p} \cdot \mathbf{v}_s$$

c speed of sound
 \mathbf{v}_s superfluid velocity

move $\mathbf{p} \cdot \mathbf{v}_s$ to the left

$$E - \mathbf{p} \cdot \mathbf{v}_s = cp$$

take square

$$(E - \mathbf{p} \cdot \mathbf{v}_s)^2 - c^2 p^2 = 0$$

$$g^{\mu\nu} p_\mu p_\nu = 0$$

$$p_\nu = (-E, \mathbf{p})$$

$$g^{00} = -1 \quad g^{0i} = -v_s^i \quad g^{ij} = c^2 \delta^{ij} - v_s^i v_s^j$$

effective metric

inverse metric $g_{\mu\nu}$ determines effective spacetime in which phonons move along geodesic curves

$$ds^2 = -c^2 dt^2 + (d\mathbf{r} - \mathbf{v}_s dt)^2$$

reference frame for phonon is dragged by moving liquid

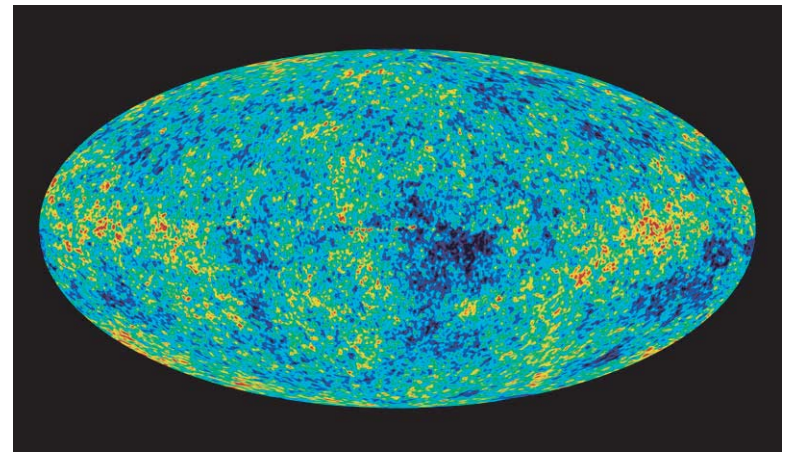
$$ds^2 = g_{\mu\nu} dx^\mu dx^\nu$$

geometry is emergent

Superfluid 4He and BEC



Universe



acoustic gravity

metric theories of gravity

general relativity



geometry of effective space time
for quasiparticles (phonons)

$$g_{\mu\nu}$$

geometry of space time
for matter



geodesics for phonons

$$ds^2 = g_{\mu\nu} dx^\mu dx^\nu = 0$$

geodesics for photons

Landau two-fluid equations

$$\dot{\rho} + \nabla \cdot (\rho \mathbf{v}_s + \mathbf{P}^{\text{Matter}}) = 0$$

dynamic equations
for metric field $g_{\mu\nu}$

$$\dot{\mathbf{v}}_s + \nabla(\mu + \mathbf{v}_s^2/2) = 0$$

equations
for superfluid
component

equation
for normal
component

$$T_{;\nu}^{\mu\nu} \text{ Matter} = 0$$

equation
for matter

Einstein equations of GR

$$\frac{1}{8\pi G} (R_{\mu\nu} - g_{\mu\nu} R/2) = T_{\mu\nu}^{\text{Matter}}$$

1/2 of GR



message from: **cond-mat**

to: **quantum gravity**



is gravity fundamental ?

it may emerge as classical output of underlying quantum system



as hydrodynamics ?

underlying microscopic quantum system at high energy

superfluid helium

quantum vacuum

classical 2-fluid hydrodynamics

emergent low-energy effective theory

classical general relativity

$$\dot{\rho} + \nabla \cdot (\rho \mathbf{v}_s + \mathbf{P}^{\text{Matter}}) = 0$$

$$\dot{\mathbf{v}}_s + \nabla (\mu + \mathbf{v}_s^2/2) = 0$$

$$\nabla_\nu T_{\text{Matter}}^{\mu\nu} = 0$$

$$\frac{1}{8\pi G} (R_{\mu\nu} - g_{\mu\nu} R/2) - \Lambda g_{\mu\nu} = T_{\mu\nu}^{\text{Matter}}$$

Schwarzschild-Painleve-Gulstrand acoustics

Superfluid 4He & BEC

Gravity

speed of sound

c

speed of light

phonon spectrum

$$g^{\mu\nu} p_\mu p_\nu = 0$$

$g^{\mu\nu}$

geometry

$$ds^2 = g_{\mu\nu} dx^\mu dx^\nu$$



$$ds^2 = - dt^2 c^2 + (dr - v dt)^2$$

Acoustic metric for phonons
propagating in radial flow $v(r)$

Painleve-Gulstrand metric

outside gravitating body of mass M

$$ds^2 = - dt^2 (c^2 - v^2) + 2 v dr dt + dr^2 + r^2 d\Omega^2$$

g_{00}

g_{0r}

after transformation

$$ds^2 = - dt^2 (c^2 - v^2) + 2 v dr dt + dr^2 + r^2 d\Omega^2$$

g_{00}

g_{0r}

$$dt = dt - v dr / (c^2 - v^2)$$

Schwarzschild metric is obtained:

$$ds^2 = - dt^2 (c^2 - v^2) + dr^2 / (c^2 - v^2) + r^2 d\Omega^2$$

g_{00}

g_{rr}

$$v^2(r) = \frac{2GM}{r}$$

Kinetic energy of superflow = potential of gravitational field

Sonic Black Hole



Liquids, BEC & superfluids

acoustic horizon

(Unruh, 1981)

Gravity



black hole horizon

Painleve-Gulstrand metric

$$ds^2 = - dt^2 (c^2 - v^2) + 2 v dr dt + dr^2 + r^2 d\Omega^2$$

ξ_{00} \uparrow \uparrow ξ_{0r}

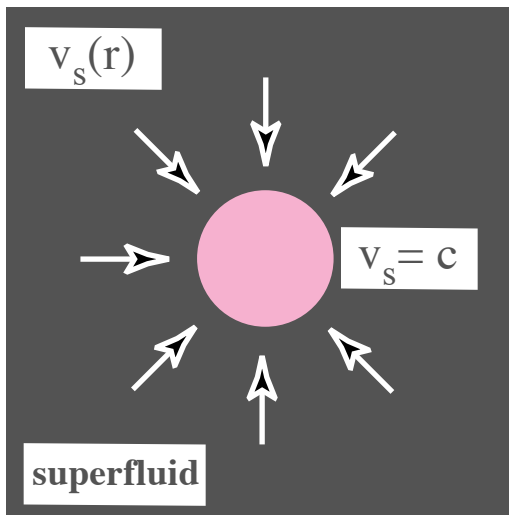
Schwarzschild metric

$$ds^2 = - dt^2 (c^2 - v^2) + dr^2 / (c^2 - v^2) + r^2 d\Omega^2$$

ξ_{00} \uparrow \uparrow ξ_{rr}

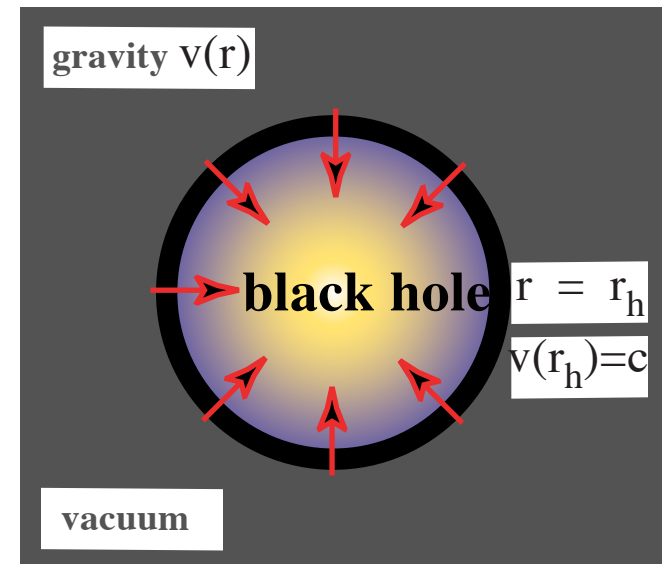
$$v^2(r) = c^2 \frac{r_h}{r}$$

$$v^2(r) = \frac{2GM}{r} = c^2 \frac{r_h}{r}$$




Information from interior region cannot be transferred by light or sound

horizon at $\xi_{00} = 0$
(or $v(r_h) = c$)



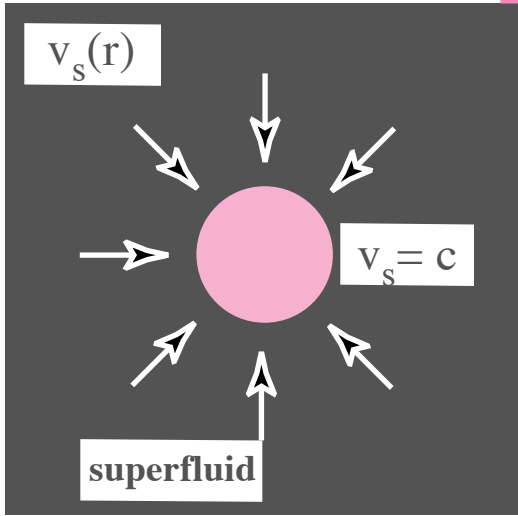
Landau critical velocity = black hole horizon

BEC, superfluid ⁴He



acoustic horizon

$$v^2(r) = c^2 \frac{r_h}{r}$$



(Unruh, 1981)

Painleve-Gulstrand metric

$$ds^2 = - dt^2 (c^2 - v^2) + 2 v dr dt + dr^2 + r^2 d\Omega^2$$

g_{00} ↑ ↑ g_{0r}

Schwarzschild metric

$$ds^2 = - dt^2 (c^2 - v^2) + dr^2 / (c^2 - v^2) + r^2 d\Omega^2$$

g_{00} ↑ ↑ g_{rr}

horizon at $g_{00} = 0$
 where flow velocity
 (velocity of local frame in GR)
 reaches Landau critical velocity

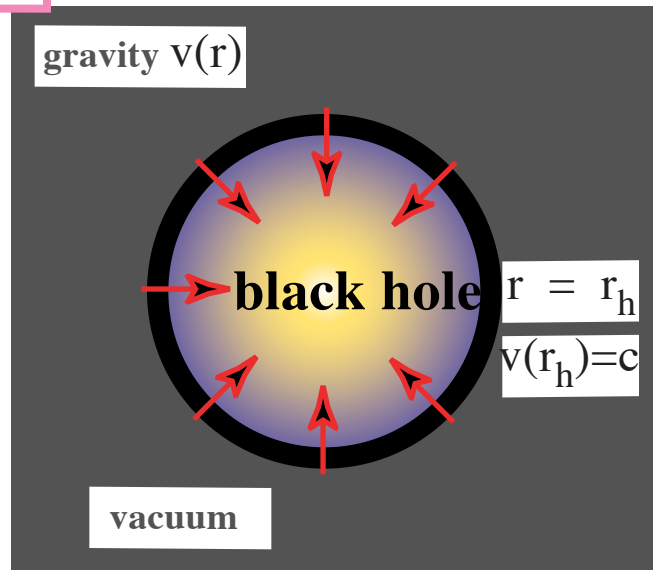
$$v(r_h) = v_{\text{Landau}} = c$$

Gravity



black hole horizon

$$v^2(r) = \frac{2GM}{r} = c^2 \frac{r_h}{r}$$



Hawking radiation is phonon/photon creation
 above Landau critical velocity



gravity emerging near Fermi point

Atiyah-Bott-Shapiro construction:

linear expansion of Hamiltonian near the nodes in terms of Dirac γ -matrices

emergent relativity

linear expansion near Fermi point

$$H = e_i^k \Gamma^i \cdot (p_k - p_k^0)$$

emergent fierbein gravity emergent γ - matrices

$$g^{\mu\nu} (p_\mu - eA_\mu - e\tau \cdot \mathbf{W}_\mu) (p_\nu - eA_\nu - e\tau \cdot \mathbf{W}_\nu) = 0$$

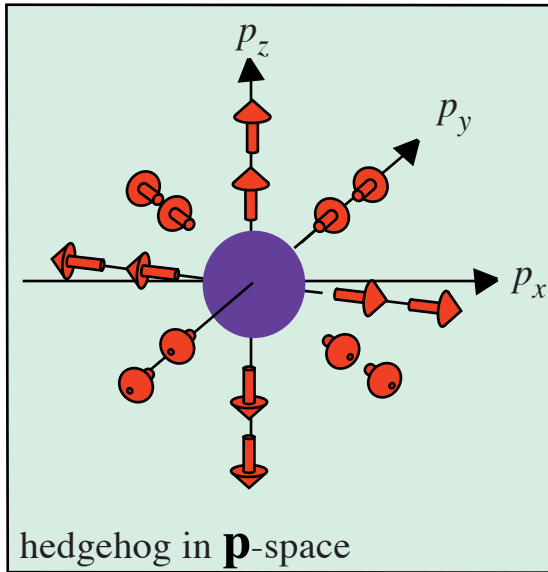
effective metric:
emergent gravity

effective $SU(2)$ gauge field

effective isotopic spin

effective electromagnetic field

effective electric charge
 $e = +1$ or -1



**all ingredients of Standard Model :
chiral fermions & gauge fields
emerge in low-energy corner**

**together with spin, Dirac Γ -matrices, gravity & physical laws:
Lorentz & gauge invariance, equivalence principle, etc**

*gravity & gauge fields
are collective modes
of vacua with Fermi point*



crossover from hydrodynamics to Einstein general relativity

they represent two different limits of hydrodynamic type equations

equations for $g^{\mu\nu}$ depend on hierarchy of ultraviolet cut-off's:
Planck energy scale E_{Planck} vs Lorentz violating scale E_{Lorentz}



$E_{\text{Planck}} \gg E_{\text{Lorentz}}$
emergent hydrodynamics

$E_{\text{Planck}} \ll E_{\text{Lorentz}}$
emergent general relativity



${}^3\text{He-A}$ with Fermi point

$E_{\text{Lorentz}} \ll E_{\text{Planck}}$
 $E_{\text{Lorentz}} \sim 10^{-3} E_{\text{Planck}}$

Universe

$E_{\text{Lorentz}} \gg E_{\text{Planck}}$
 $E_{\text{Lorentz}} > 10^9 E_{\text{Planck}}$

Fermionic horizons:

horizon in flowing 3He film & in moving soliton

Quasiparticle energy spectrum in moving texture (Doppler shift)

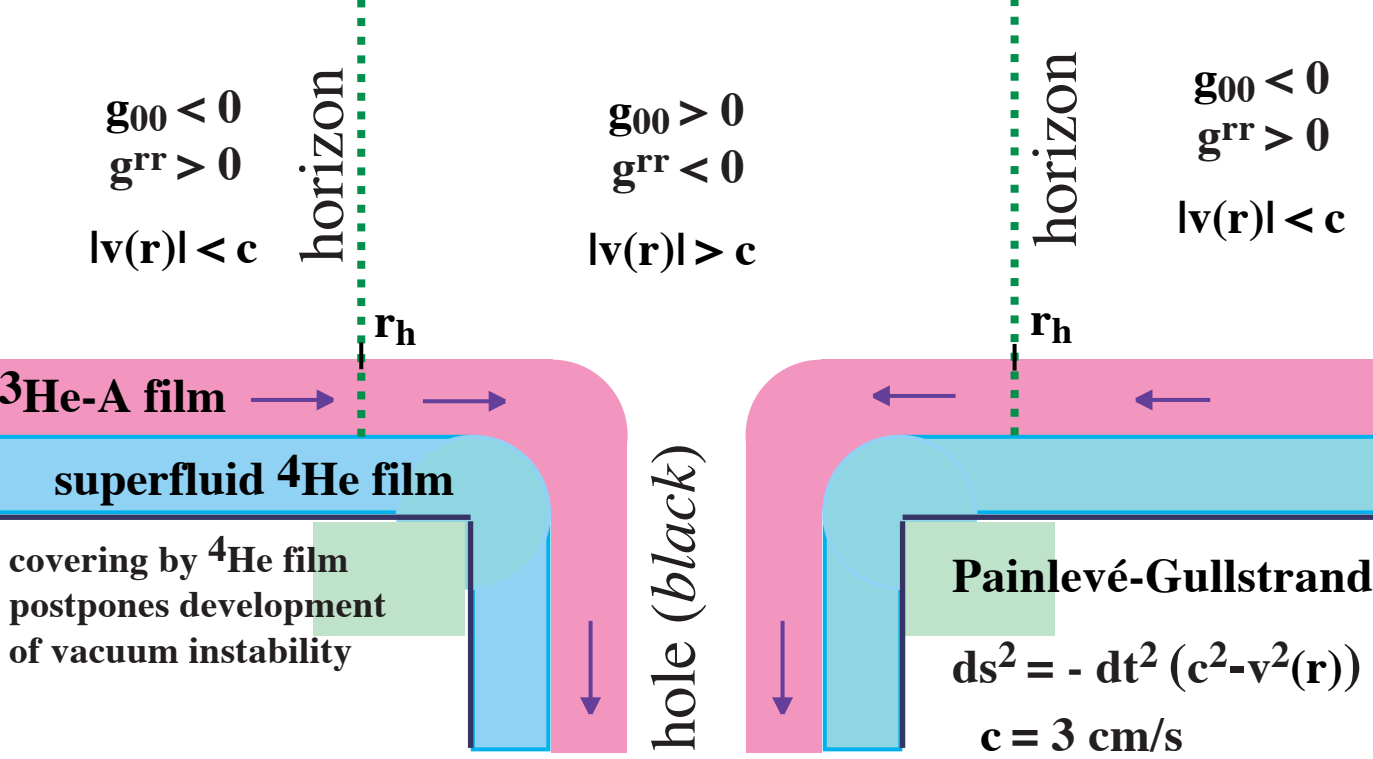
$$(E - vp_z)^2 = c_x^2 (p_x - p_F)^2 + c_z^2(z) p_z^2 + c_y^2 p_y^2$$

$$ds^2 = - dt^2 (1 - v^2/c_z^2(z)) - 2(v/c_z^2(z)) dz dt + c_x^{-2} dx^2 + c_y^{-2} dy^2 + c_z^{-2}(z) dz^2$$

Soliton: speed of light $c_z(z)$ changes sign across the soliton

Hawking temperature $T_H = (\hbar/2\pi) (dc_z/dz)_{\text{hor}}$

black hole in ${}^3\text{He-A}$ film

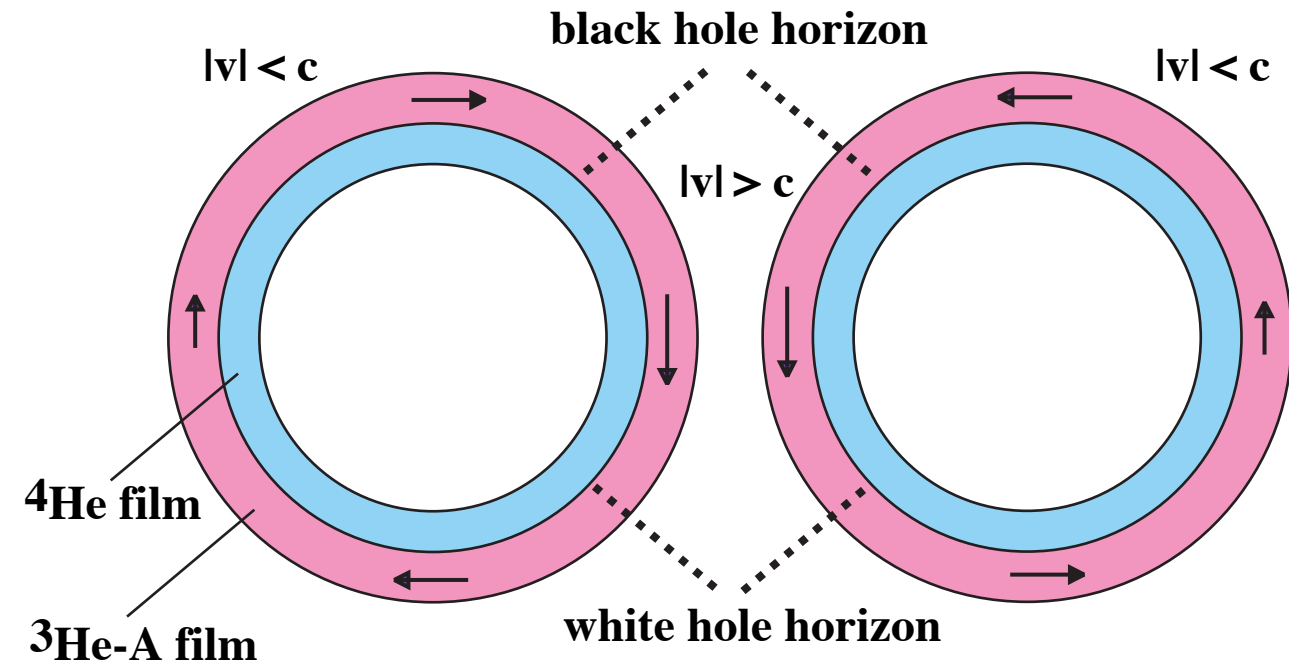


Painlevé-Gullstrand form of 2D black hole:

$$ds^2 = - dt^2 (c^2 - v^2(r)) + 2v(r) dr dt + dr^2 + r^2 d\phi^2$$

$c = 3 \text{ cm/s}$

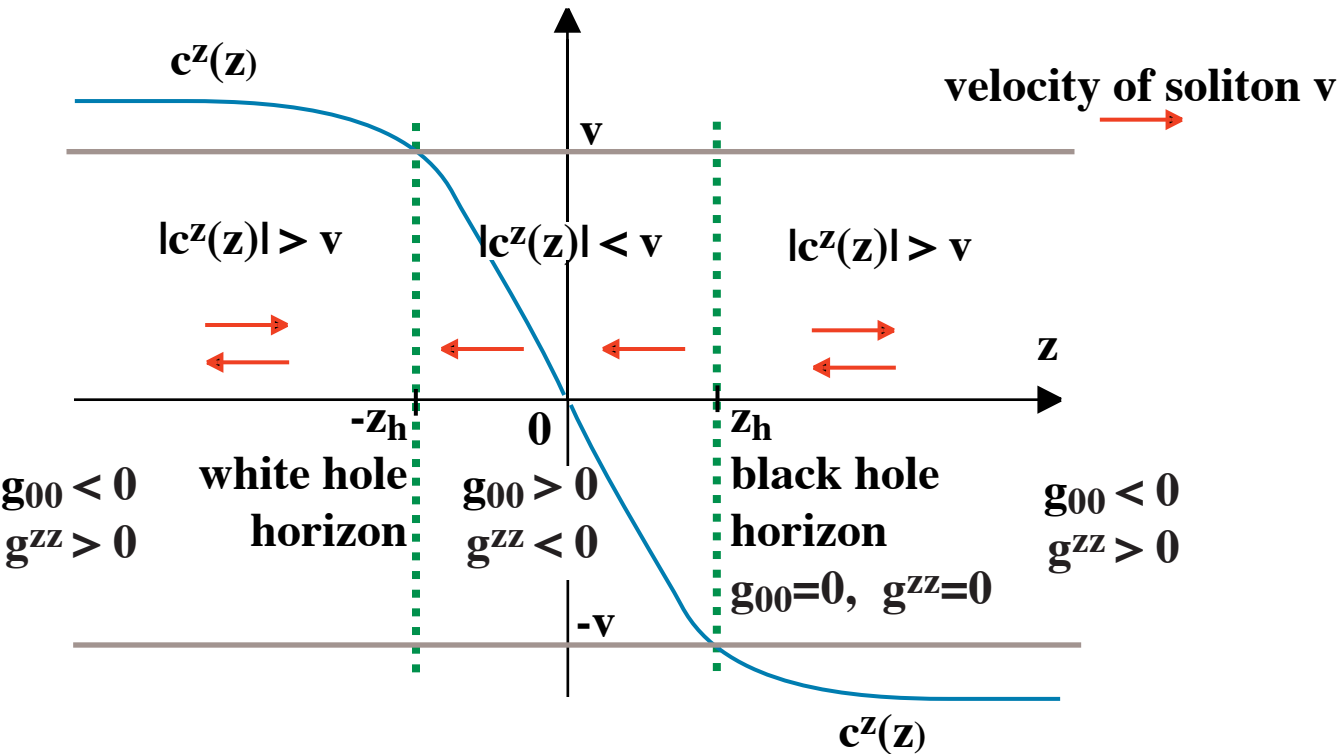
If ${}^3\text{He-A}$ film is moving
to the hole $v(r) = b / r$
horizon is at $r_h = b / c$



Horizon in moving soliton

Jacobson-GV, PRD 58, 064021 (1998)

speed of light $c_z(z)$ changes sign across the soliton



Hawking radiation leads to deceleration of soliton until horizons shrink leaving bare singularity

physical consequence of Hawking radiation is the decay of horizon

this is what to be measured

effective metric

$$ds^2 = - dt^2 (1 - v^2/c_z^2(z)) - 2(v/c_z^2(z))dzdt + c_x^{-2}dx^2 + c_y^{-2}dy^2 + c_z^{-2}(z)dz^2$$

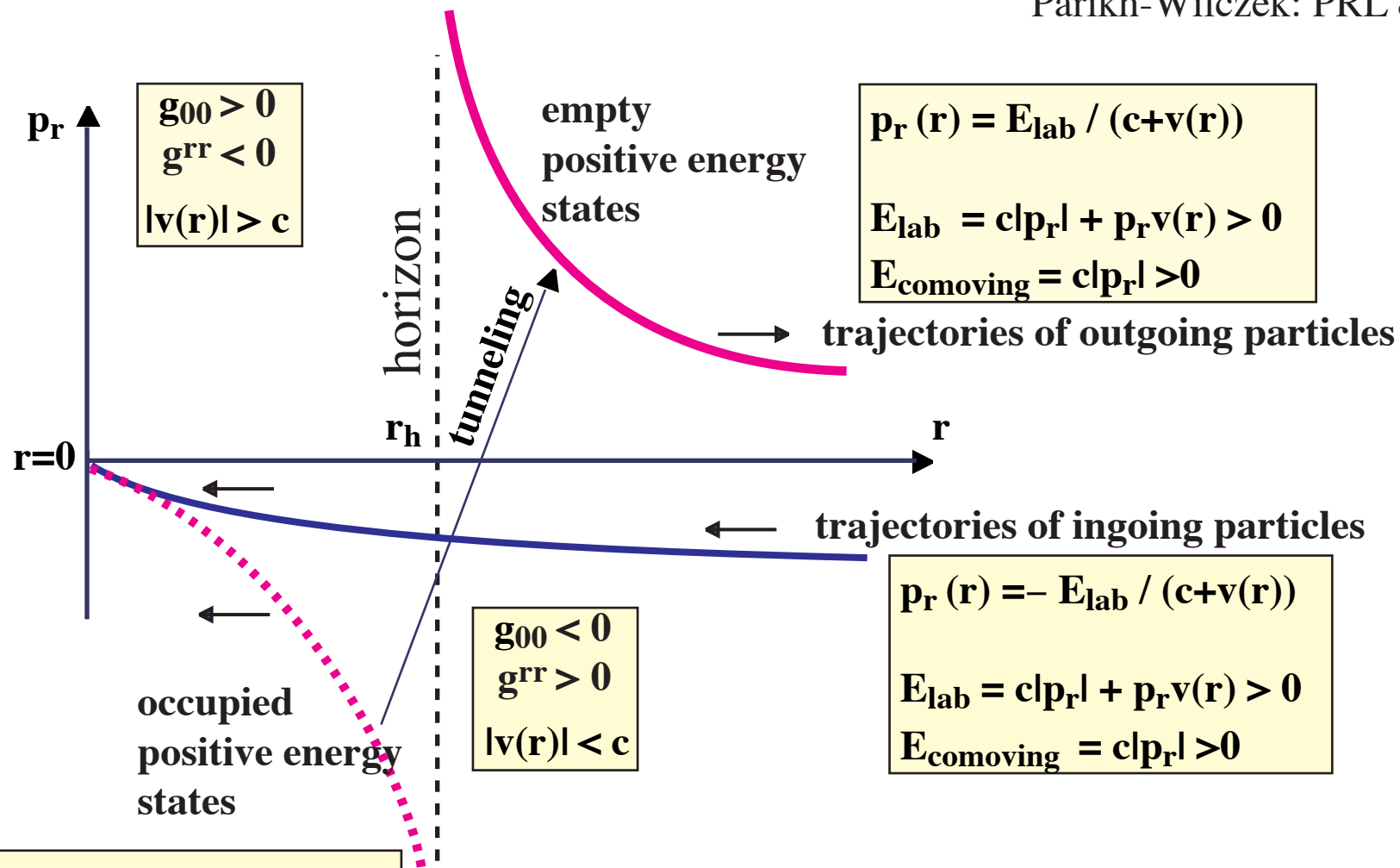
horizons at $c_z(z_h) = \pm v$ $g_{00}=0$ $g^{zz}=0$ between horizons particles move in one direction

$z=0$ - curvature singularity : $c_z(z)=0$

Hawking radiation in semiclassical description: quantum tunneling between classical trajectories

GV: JETP Lett. **69**,705 (1999)

Parikh-Wilczek: PRL **85**,5042 (2000)



$p_r(r) = E_{lab} / (c+v(r))$
 $E_{lab} = -c|p_r| + p_r v(r) > 0$
 $E_{comoving} = -c|p_r| < 0$

tunneling exponent and Hawking temperature

$$S(E) = 2 \operatorname{Im} \int dr p_r(r) = 2E \operatorname{Im} \int dr / (c+v(r)) = E / T_H$$

$$W = w e^{-S(E)} = w e^{-E/T_H} \quad T_H = v'/2\pi$$

problem of Hawking radiation in de Sitter universe

GV: 0803.3367

$$ds^2 = - dt^2 c^2 + (dr - v dt)^2 \quad ds^2 = - dt^2 (c^2 - H^2 r^2) + 2 Hr dr dt + dr^2 + r^2 d\Omega^2$$

$$v(\mathbf{r}) = Hr \quad \text{de Sitter expansion}$$

* tunneling approach

outside de Sitter horizon $r > c/H$

$$p_r(\mathbf{r}) = E / (-c + v(\mathbf{r})) > 0$$

$$E = -c|p_r| + p_r v(\mathbf{r}) > 0$$

$$E_{\text{comoving}} = -c|p_r| < 0$$

inside de Sitter horizon $r < c/H$

$$p_r(\mathbf{r}) = E / (-c + v(\mathbf{r})) < 0$$

$$E = c|p_r| + p_r v(\mathbf{r}) > 0$$

$$E_{\text{comoving}} = c|p_r| > 0$$

tunneling from occupied
positive energy states

$$S(E) = 2 \operatorname{Im} \int dr p_r(r) = 2E \operatorname{Im} \int dr / (-c + Hr) = 2\pi E / H$$

$$W = w e^{-S(E)} = w e^{-E/T_H} \quad T_H = H / 2\pi$$

* Does Hawking radiation really occur in de Sitter universe ?

BH has preferred reference frame, de Sitter Universe does not

* cond-mat simulation $ds^2 = - dt^2 + d\mathbf{r}^2 / c^2 \quad c(t) = e^{-Ht}$

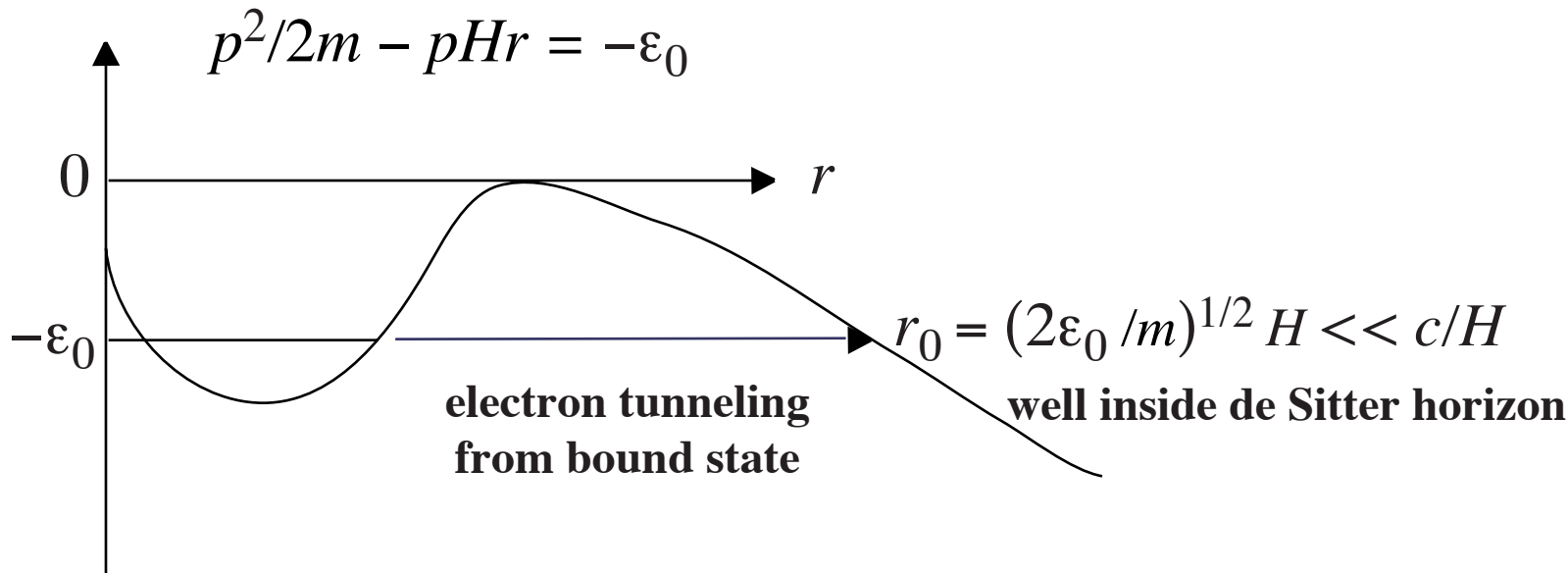
Schutzhold: PRL **95** (2005) 135703; GV: J.Low.Temp.Phys. **113** (1998) 667

activation in de Sitter space-time

$$ds^2 = - dt^2 c^2 + (dr - v dt)^2 \quad ds^2 = - dt^2 (c^2 - H^2 r^2) + 2 Hr dr dt + dr^2 + r^2 d\Omega^2$$

$$v(\mathbf{r}) = Hr \quad \text{de Sitter expansion}$$

* ionization rate of an atom from electron at atomic level $-\varepsilon_0$



$$S(-\varepsilon_0) = 2 \operatorname{Im} \int_0^{r_0} dr p_r(r) = \pi \varepsilon_0 / H$$

$$W = w e^{-S(-\varepsilon_0)} = w e^{-\varepsilon_0/T}$$

effective activation temperature

$$T = H / \pi$$

is twice the Hawking temperature

$$T_H = H / 2\pi$$

suggestion for Hawking type radiation in de Sitter space-time GV: 0803.3367

$$ds^2 = - dt^2 c^2 + (dr - v dt)^2 \quad ds^2 = - dt^2 (c^2 - H^2 r^2) + 2 Hr dr dt + dr^2 + r^2 d\Omega^2$$

$v(\mathbf{r}) = Hr$ de Sitter expansion

pure de Sitter vacuum does not radiate

but observer views thermal bath with twice the Hawking temperature because he/she violates de Sitter symmetry

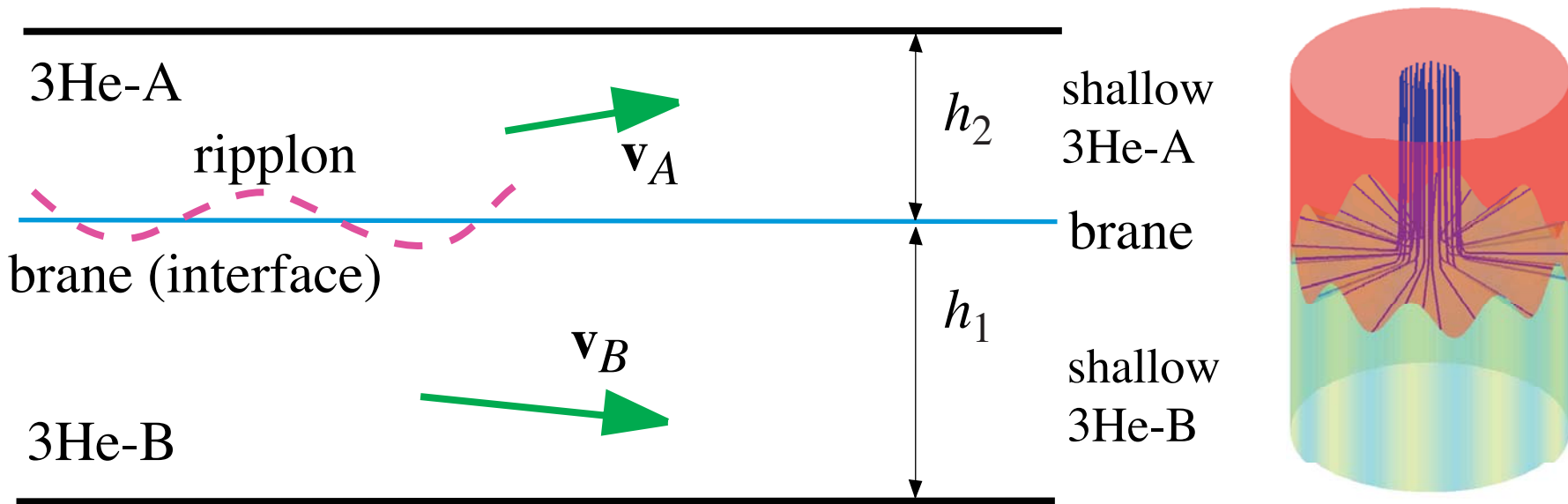
effective activation temperature

$$T = H / \pi$$

is twice the Hawking temperature

$$T_H = H / 2\pi$$

Relativistic riplons living on brane between two shallow superfluids



Spectrum of riplons $g^{\mu\nu}p_\mu p_\nu = 0$

$$\alpha_1(\omega - \mathbf{k} \cdot \mathbf{v}_A)^2 + \alpha_2(\omega - \mathbf{k} \cdot \mathbf{v}_B)^2 = c^2 k^2$$

Effective metric for riplons

$$ds^2 = g_{\mu\nu} dx^\mu dx^\nu$$

Speed of light

$$c^2 = (F/\rho_s) h_1 h_2 / (h_1 + h_2)$$

$$ds^2 = -dt^2 \frac{c^2 - W^2 - U^2}{c^2 - U^2} + dr^2 \frac{1}{c^2 - W^2 - U^2} + r^2 d\phi^2$$

F is restoring force:

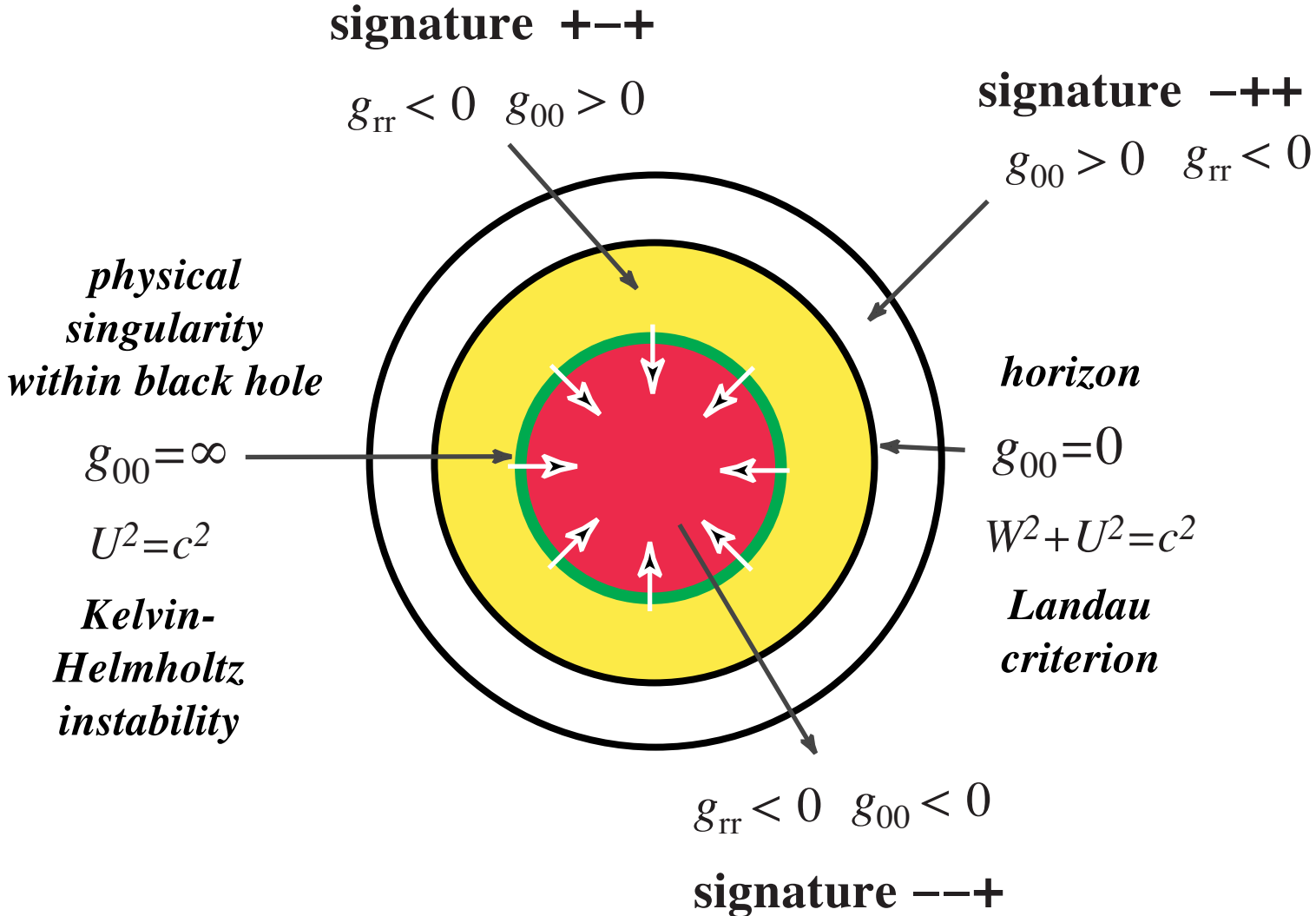
either gravity or gradient of magnetic field

$$U^2 = \alpha_1 \alpha_2 (\mathbf{v}_A - \mathbf{v}_B)^2$$

$$\mathbf{W} = \alpha_1 \mathbf{v}_A + \alpha_2 \mathbf{v}_B$$

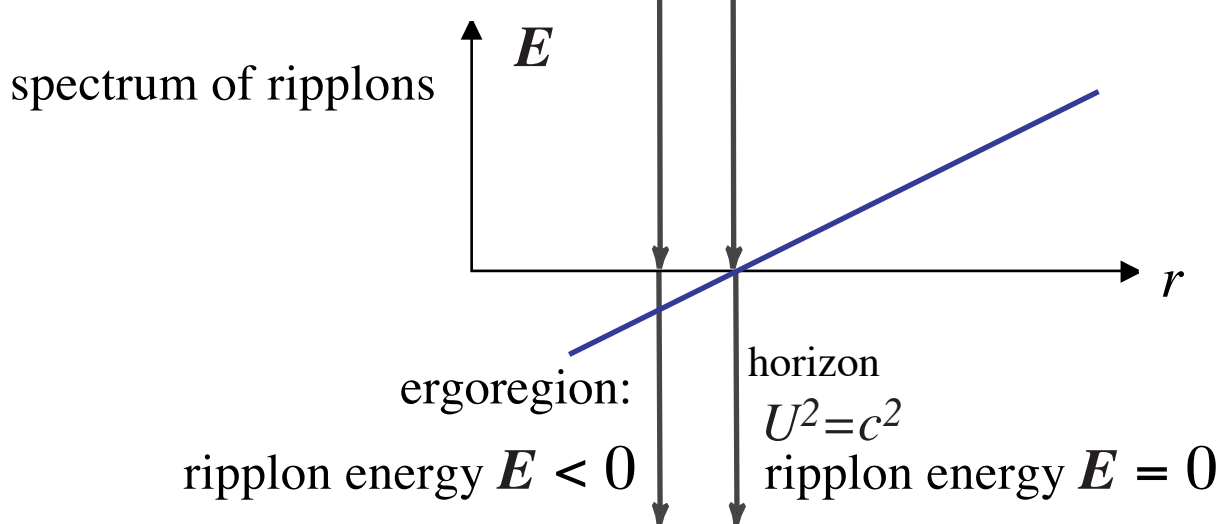
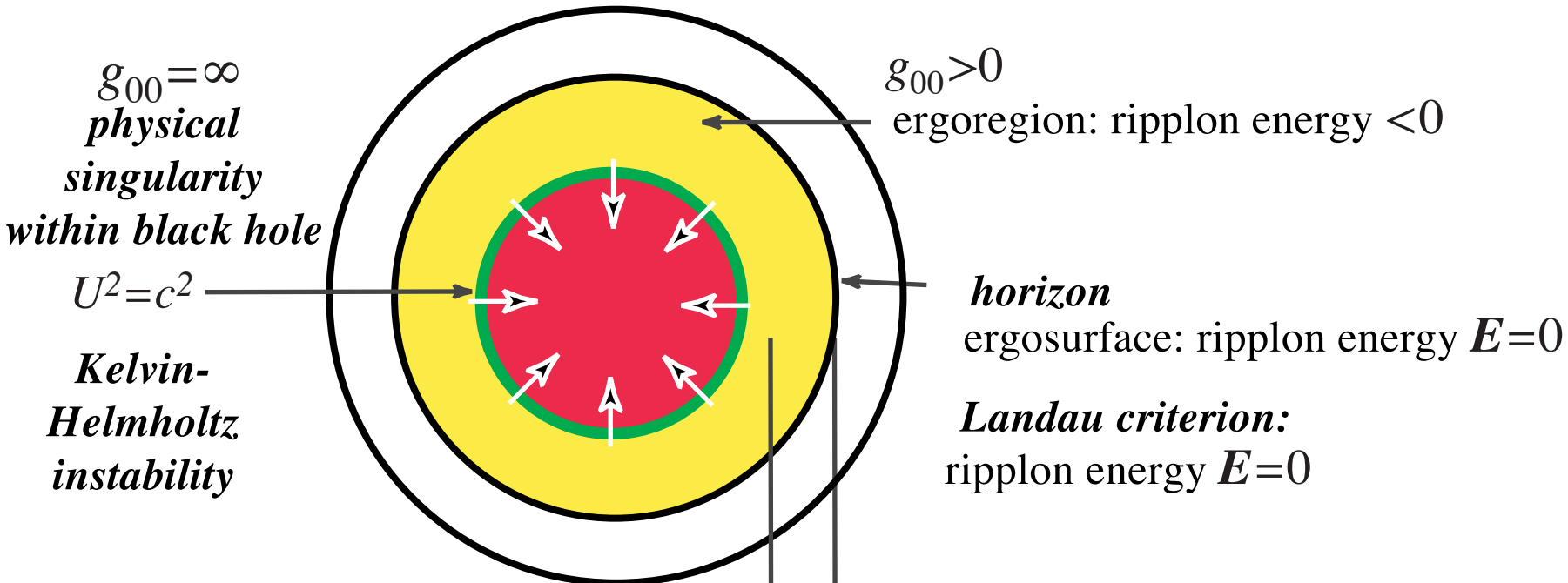
black hole for ripplons at AB-brane

$$ds^2 = -dt^2 \frac{c^2 - W^2 - U^2}{c^2 - U^2} + dr^2 \frac{1}{c^2 - W^2 - U^2} + r^2 d\phi^2$$



horizon as ergosurface

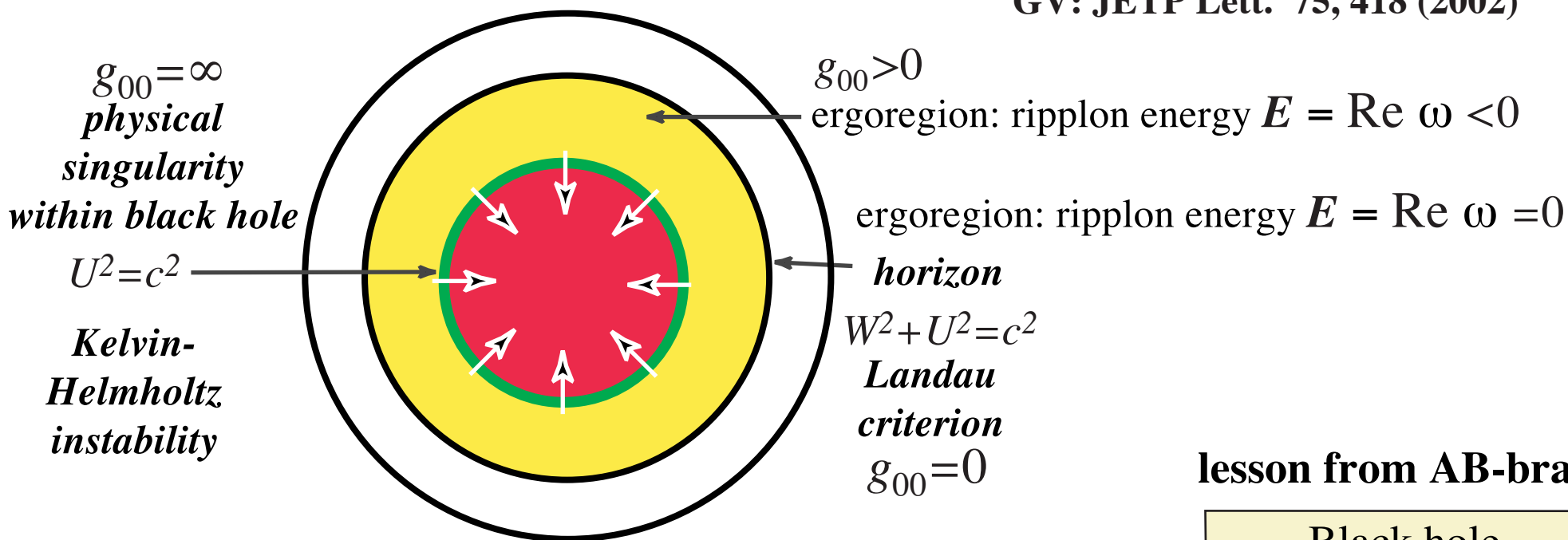
$$ds^2 = -dt^2 \frac{c^2 - W^2 - U^2}{c^2 - U^2} + dr^2 \frac{1}{c^2 - W^2 - U^2} + r^2 d\phi^2$$



instability in ergosurface

Schutzhold-Unruh: PRD 66, 044019 (2002)

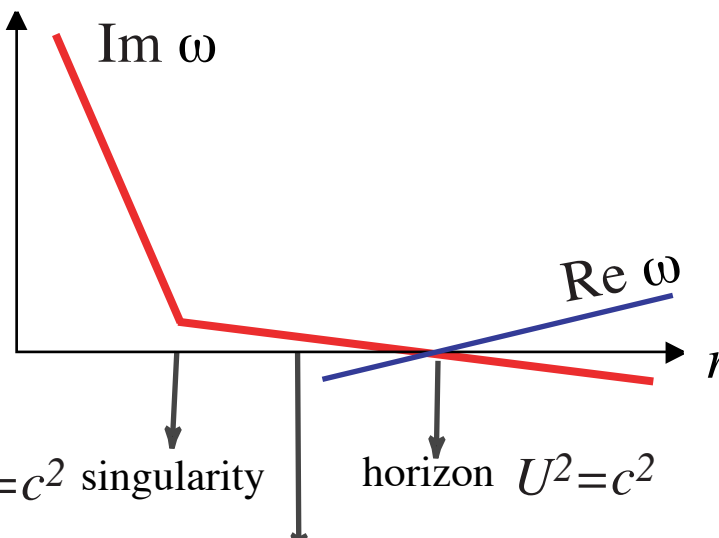
GV: JETP Lett. 75, 418 (2002)



lesson from AB-brane:

Black hole
may collapse due to
vacuum instability
in ergoregion

spectrum of ripples

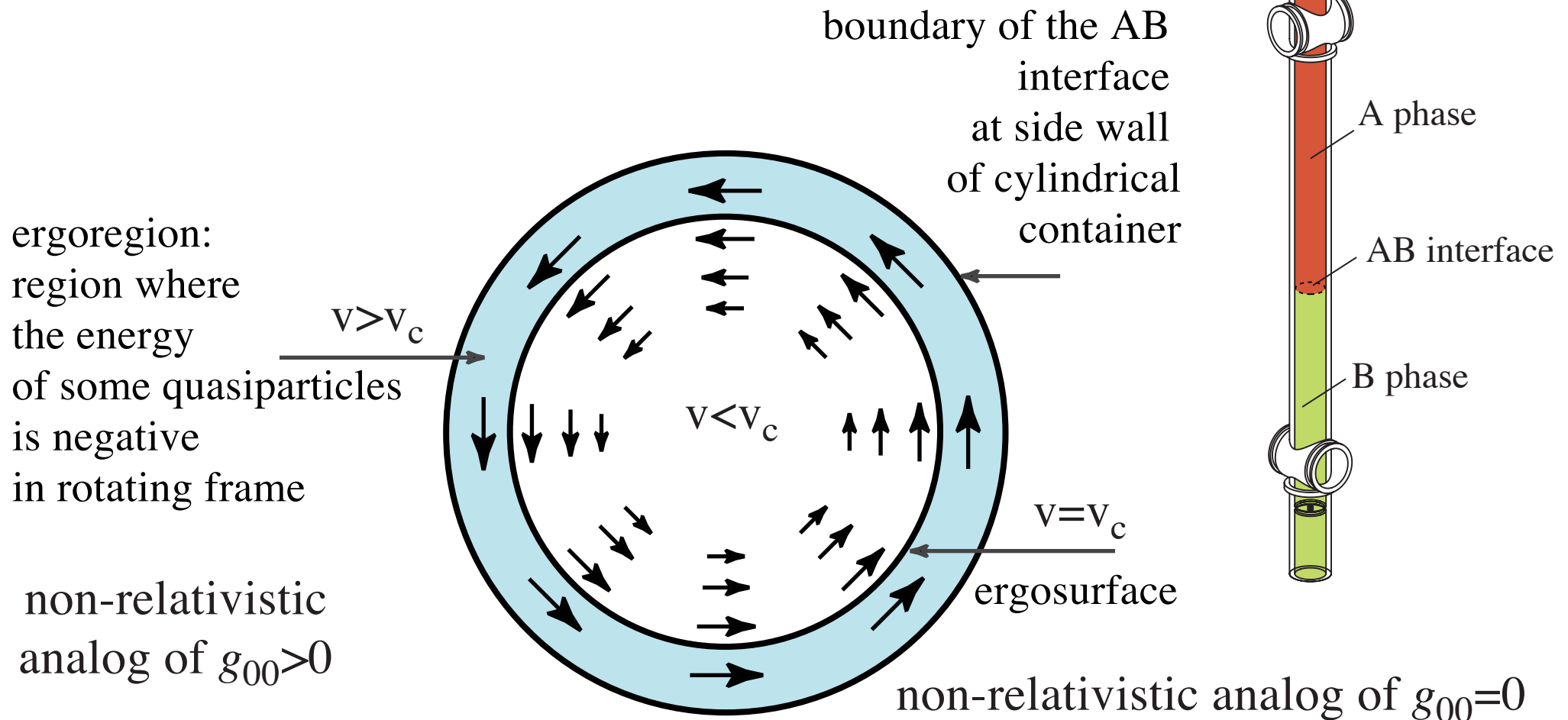


$\text{Im } \omega$ & $\text{Re } \omega$
cross 0 simultaneously:
behind horizon
attenuation transforms
to amplification

ergoregion: ripplon energy $E = \text{Re } \omega < 0$
also $\text{Im } \omega < 0$ which means instability

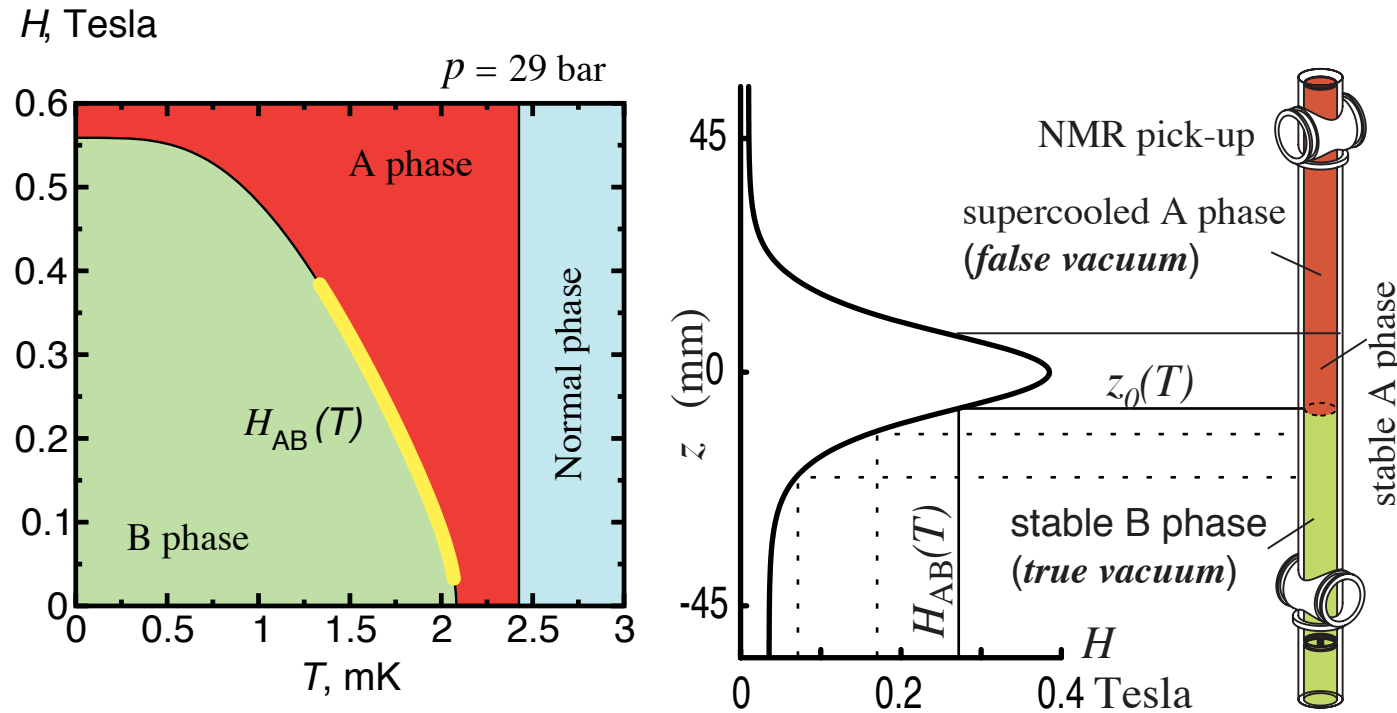
Ergoregion instability at the AB-brane in Helsinki experiments

interface between static B-phase and A phase circulating with solid-body velocity $v = \Omega r$
(velocity is shown by arrows)



experimental Landau criterion of ripplon ergregion instability

experimental set-up

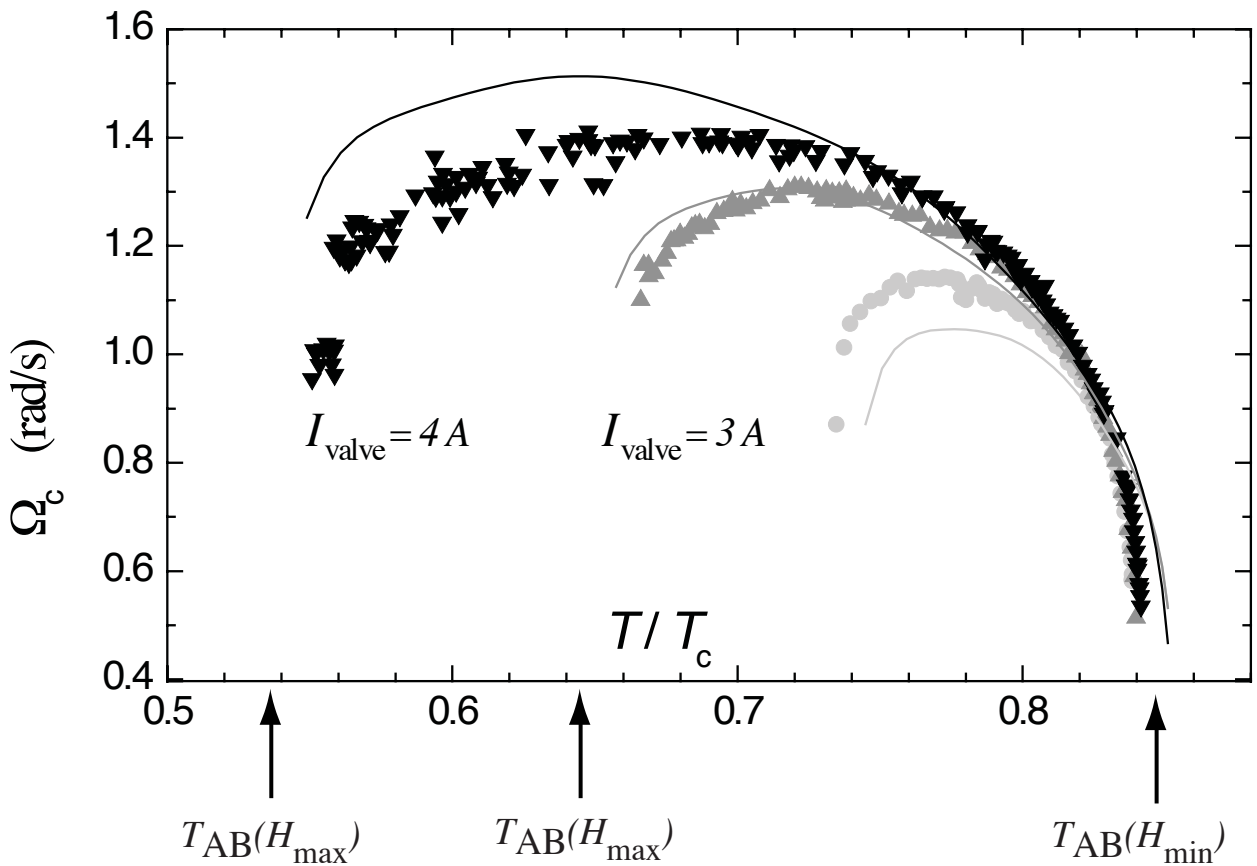


Landau criterion

$$\rho_{sB}(v_{sB} - v_n)^2 + \rho_{sA}(v_{sA} - v_n)^2 = 2\sqrt{F\sigma}$$

critical value $v_{sB}(T)$ depends on T via $\rho_{sB}(T)$, $F(T)$, $\sigma(T)$

experimental Landau criterion of ripplon ergregion instability

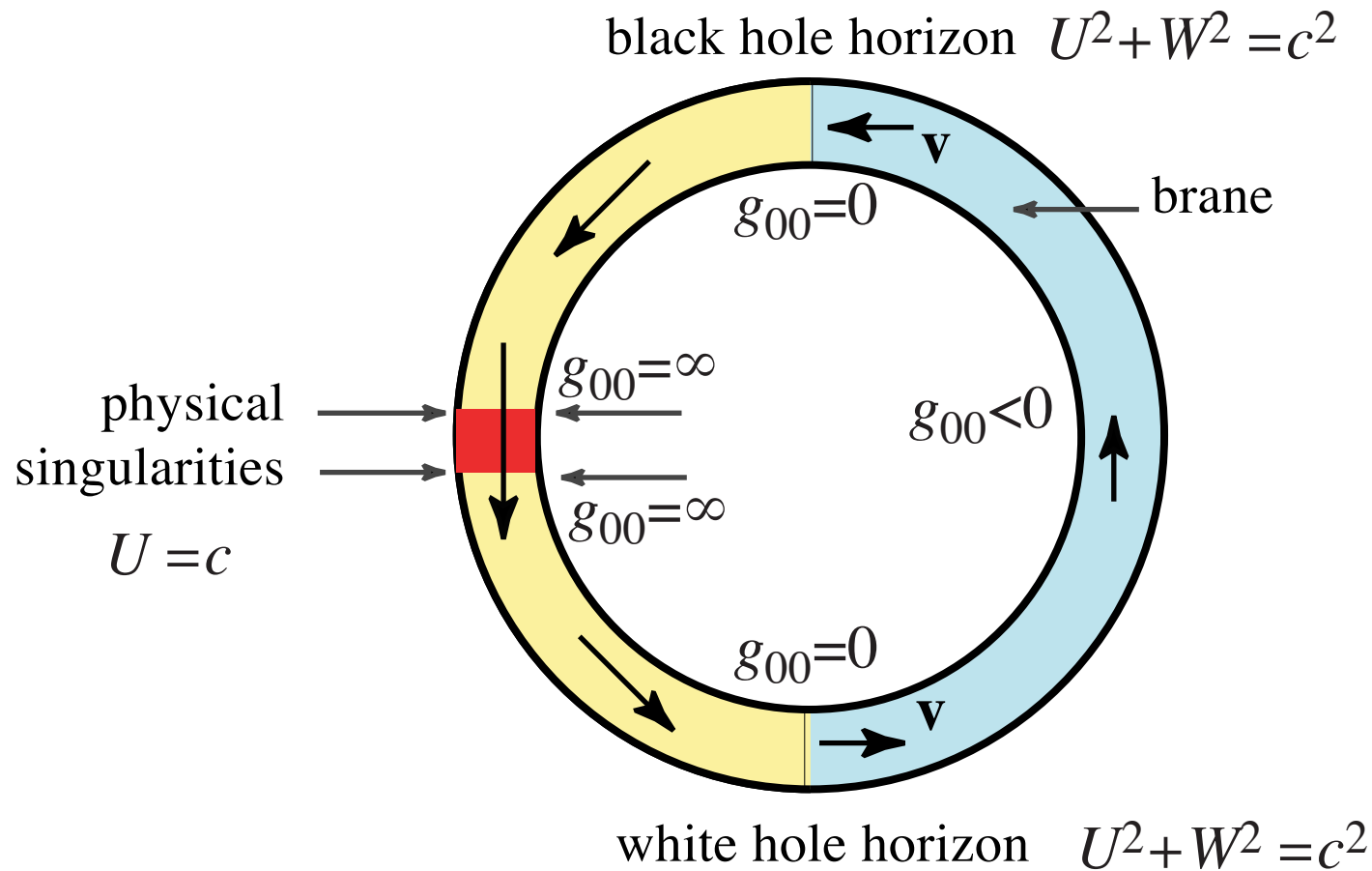


lines - theoretical values without fitting

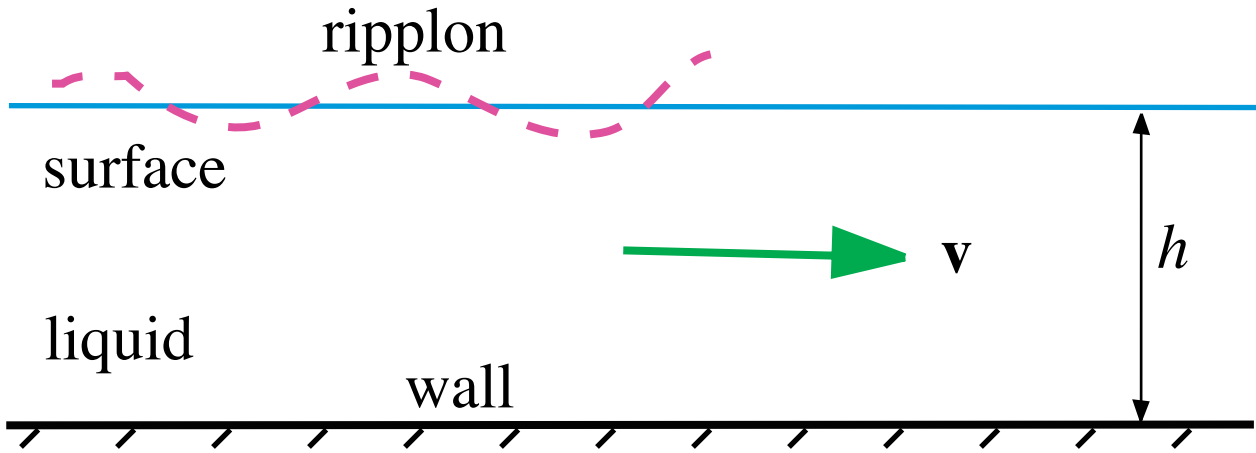
$$\rho_{sB}(v_{sB} - v_n)^2 + \rho_{sA}(v_{sA} - v_n)^2 = 2\sqrt{F\sigma}$$

Proposal for black hole for ripplons at AB-brane (azimuthal flow)

$$ds^2 = -dt^2 \frac{c^2 - W^2 - U^2}{c^2 - U^2} + \frac{r^2 d\phi^2}{c^2 - W^2 - U^2} + dr^2$$



Relativistic riplons in shallow water




Spectrum of riplons $g^{\mu\nu}p_\mu p_\nu = 0$

$$(\omega - \mathbf{k} \cdot \mathbf{v})^2 = c^2 k^2 + c^2 k^4 (\sigma / \rho g - h^2 / 3)$$

$$k \ll k_c = \sqrt{\rho g / \sigma}$$

$$kh \ll 1$$

Effective metric for riplons $ds^2 = g_{\mu\nu} dx^\mu dx^\nu$

$$ds^2 = -dt^2 (1 - v^2/c^2) + dr^2 \frac{1}{c^2 - v^2} + r^2 d\phi^2$$


Speed of "light"

$$c^2 = gh$$



ripplonic white hole

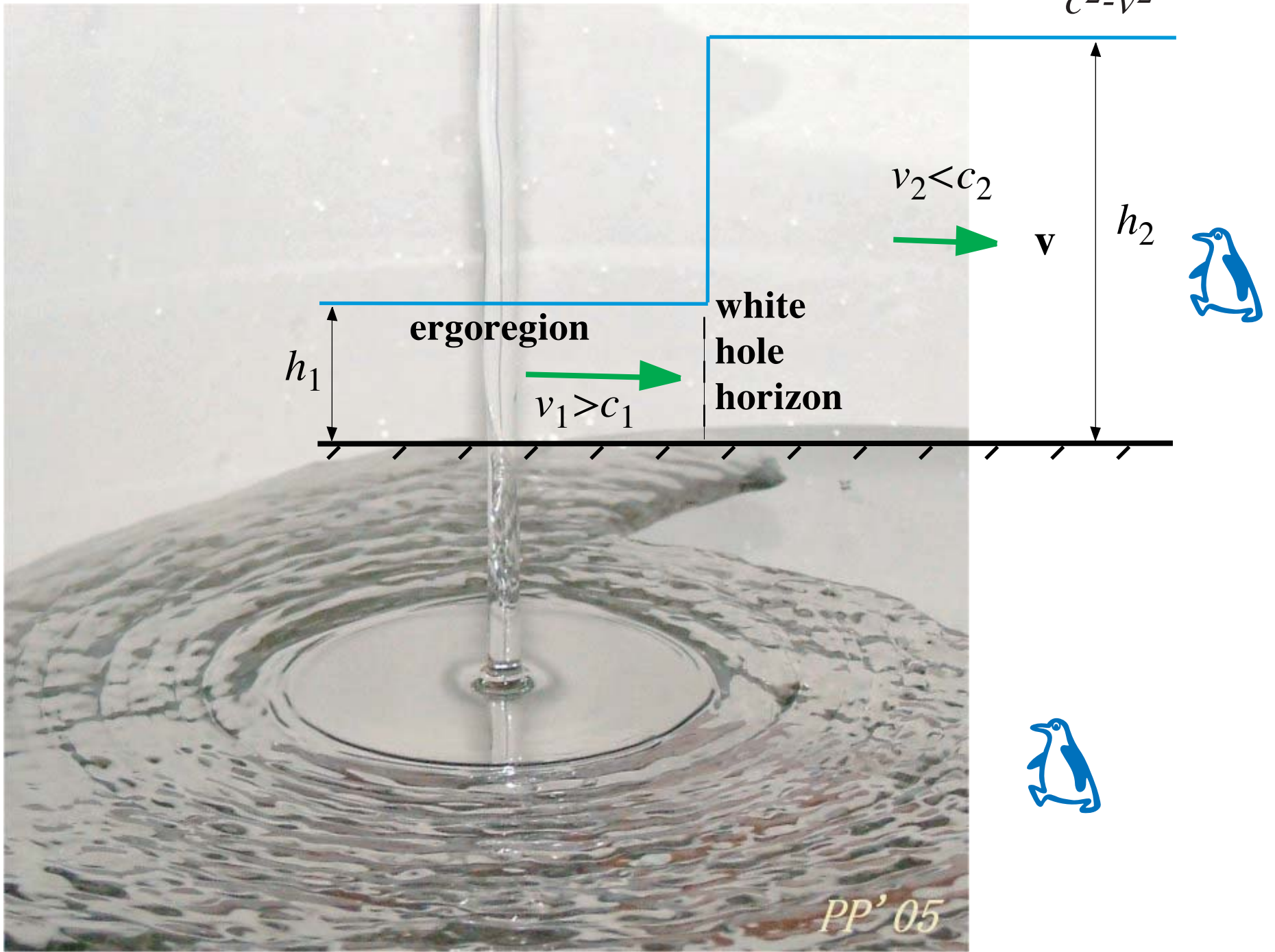
circular hydraulic jump

*Courtesy
Piotr
Pieranski*

PP' 05

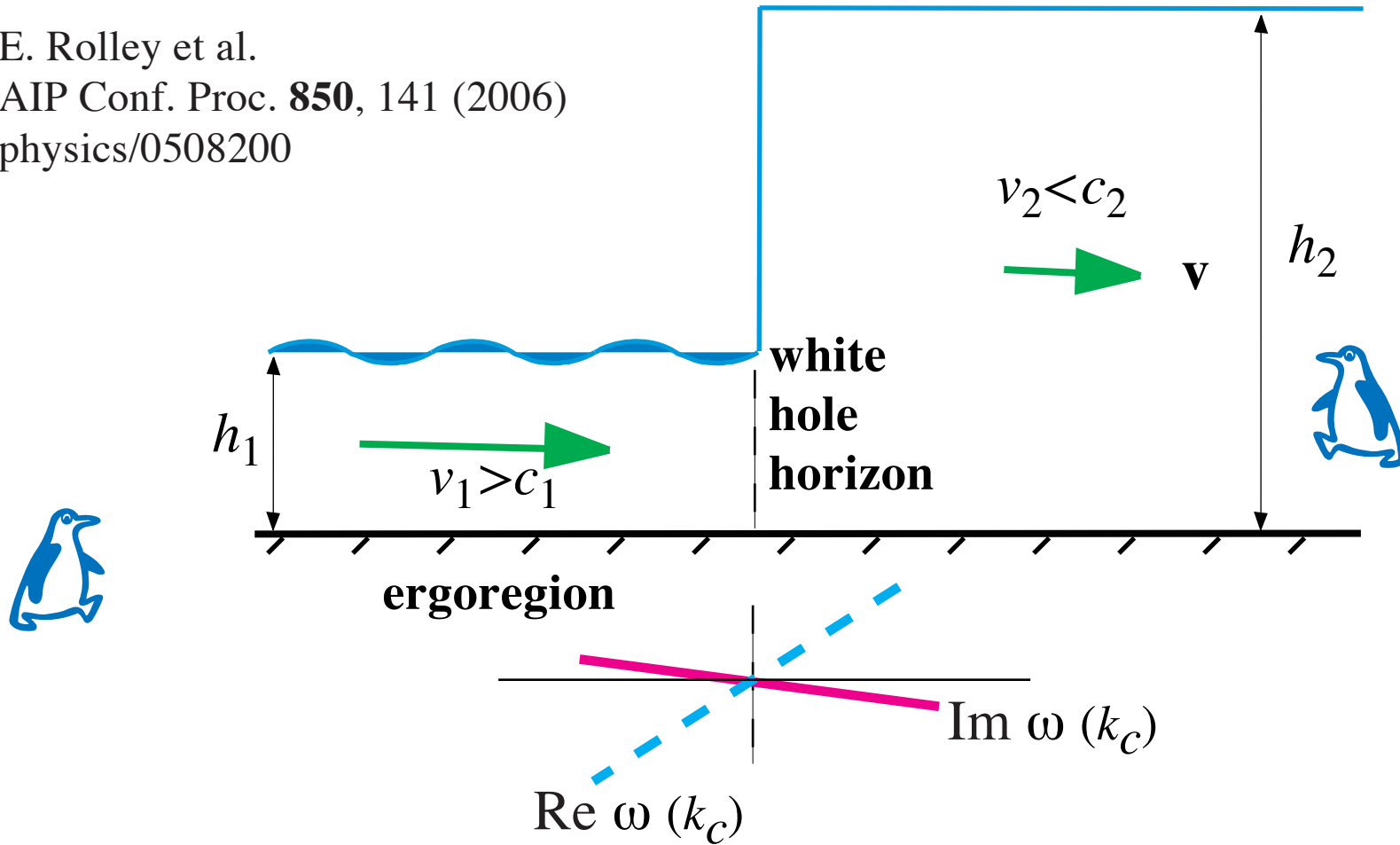
Hydraulic jump as white hole

$$ds^2 = -dt^2 (1-v^2/c^2) + dr^2 \frac{1}{c^2-v^2} + r^2 d\phi^2$$



Observation of instability in the ergoregion in superfluid ^4He

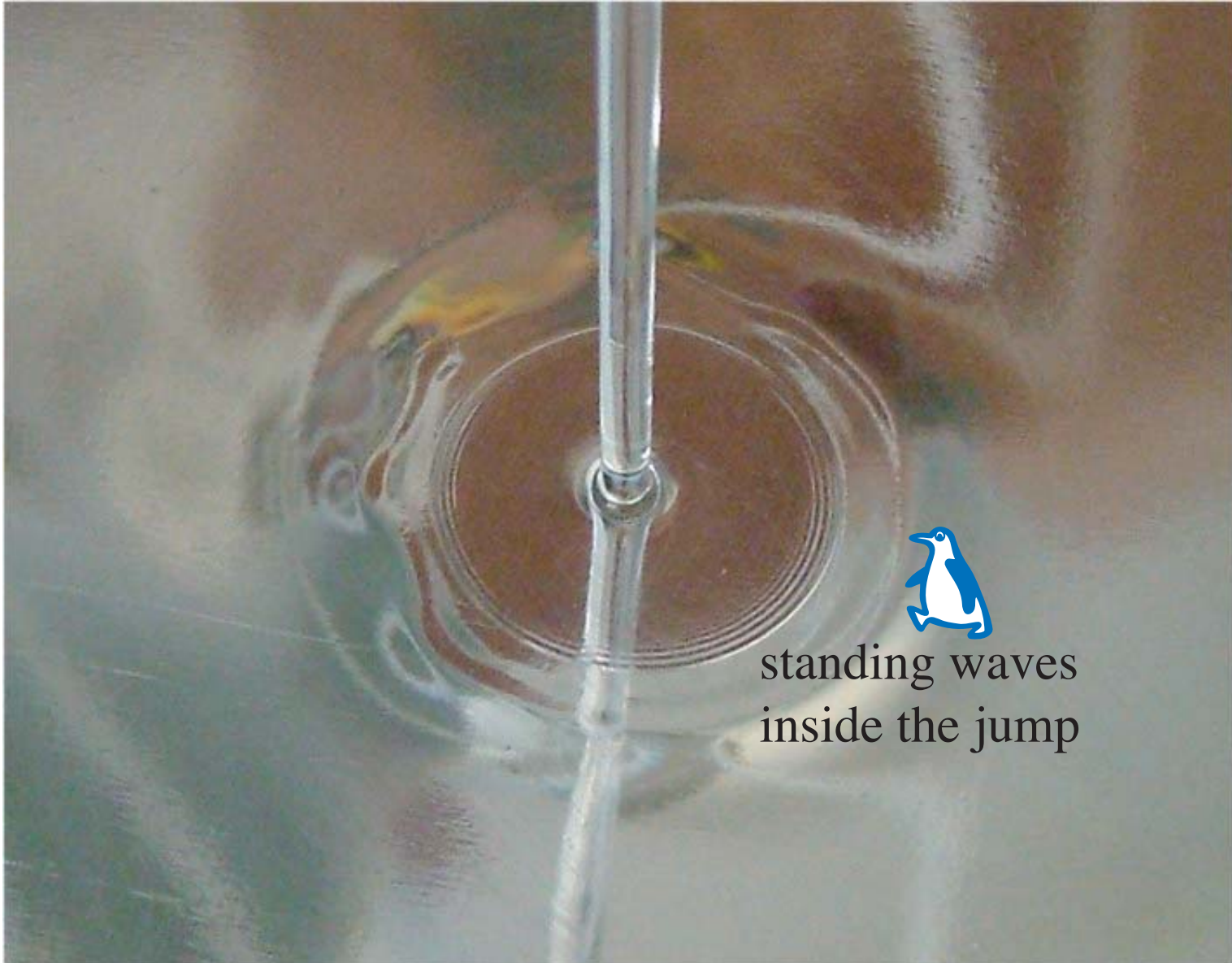
E. Rolley et al.
AIP Conf. Proc. **850**, 141 (2006)
physics/0508200



in the ergoregion
attenuation of ripplon
transforms to amplification

standing waves
are excited
in the ergoregion with
 $\text{Re } \omega(k_c) = 0$

ergoregion instability inside the "white hole" in shallow water



standing waves
inside the jump

Courtesy Piotr Pieranski

Zeldovich-Starobinsky effect, rotational Unruh effect & ergorgion instability

Radiation by object rotating in quantum vacuum & by circulating superflow

Analog to Kerr Black Hole

Hawking radiation
looks as thermal with
 $T_H \sim hg/c$
where g is gravity

Unruh effect:
radiation looks as thermal with
 $T_U \sim ha/c$
where a is acceleration of a body

metric in rotating frame

$$ds^2 = - dt^2 c^2 + (d\mathbf{r} - \mathbf{v}dt)^2$$

$$\mathbf{v}(\mathbf{r}) = \boldsymbol{\Omega} \times \mathbf{r}$$

H. Takeuchi, M. Tsubota and GV
*Zel'dovich-Starobinsky Effect
in Atomic Bose-Einstein Condensates:
Analogy to Kerr Black Hole*
J. Low Temp. Phys. **150**, 624 (2008)

$$ds^2 = - dt^2 (c^2 - \Omega^2 \rho^2) + 2\Omega \rho^2 d\phi dt + \rho^2 d\phi^2 + d\rho^2 + dz^2$$

tunneling approach

$$ds^2 = - dt^2 (c^2 - \Omega^2 \rho^2) + 2\Omega \rho^2 d\phi dt + \rho^2 d\phi^2 + d\rho^2 + dz^2$$

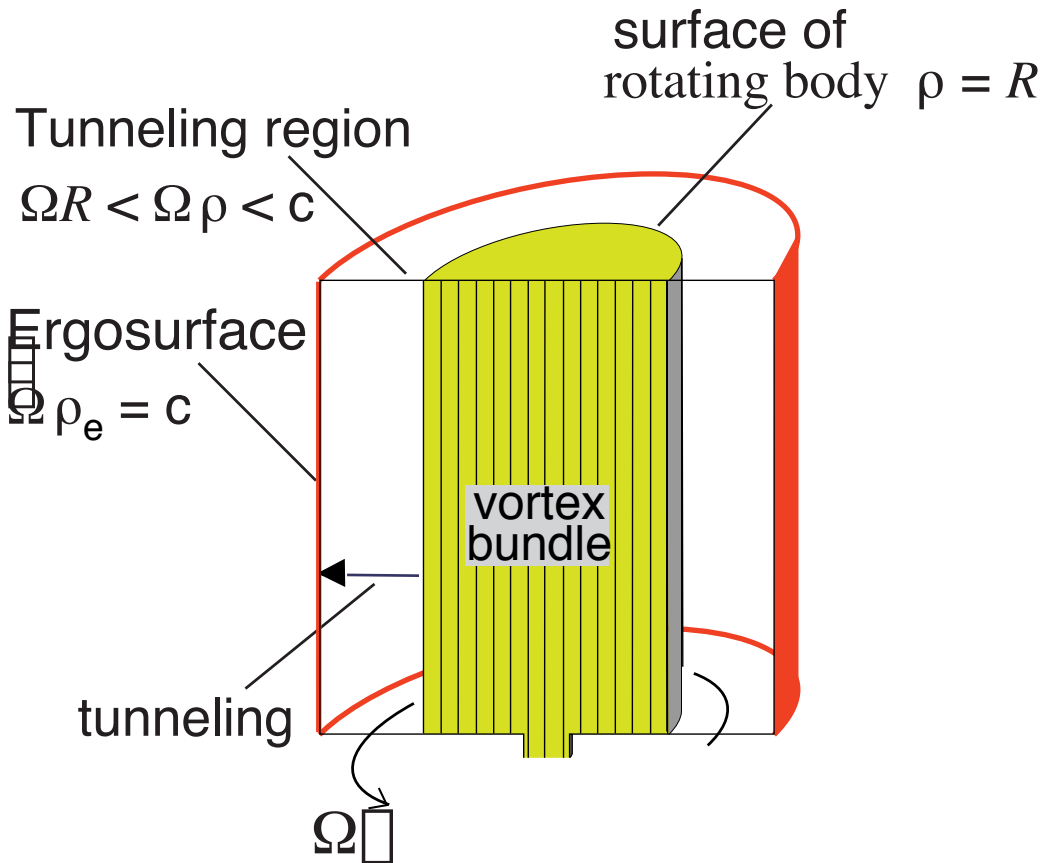
normal region $\rho < c/\Omega$

ergoregion $\rho > c/\Omega$

body surface $R < c/\Omega$ \longrightarrow ergosurface $\rho = c/\Omega$

tunneling from zero energy state at the body
to zero energy state at the ergosurface

Radiation by object rotating in quantum vacuum



optical potential made by
blue or red detuned laser
moving in static BEC
or vortex cluster rotating in static BEC

potential or a vortex moves with
linear velocity $\Omega R < c$,
but it radiates via quantum tunneling
to the ergoregion

Analog to Kerr Black Hole

tunneling exponent

$$ds^2 = - dt^2 (c^2 - \Omega^2 \rho^2) + 2\Omega \rho^2 d\phi dt + \rho^2 d\phi^2 + d\rho^2 + dz^2$$

body surface $R < c/\Omega$

—————▶
tunneling from zero energy state
at the body to the
zero energy state at the ergosurface

ergosurface $\rho = c/\Omega$

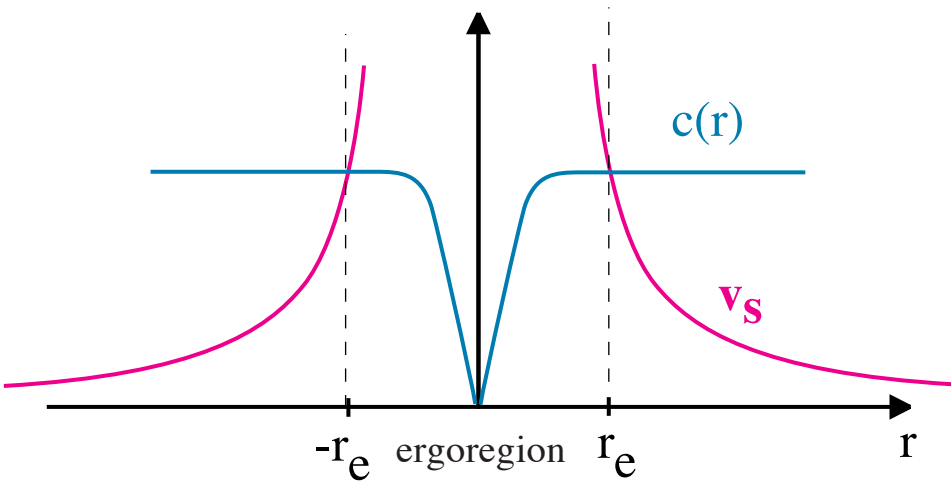
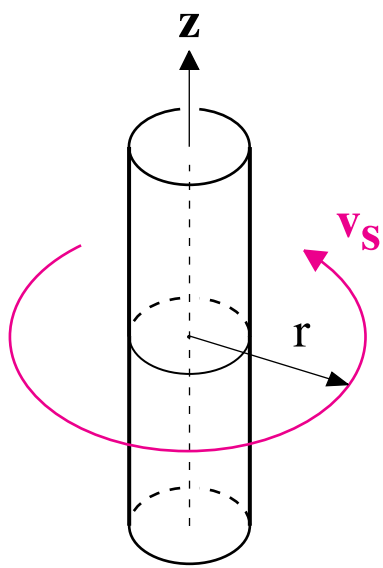
$$S(E=0) = 2 \operatorname{Im} \int d\rho p_\rho(\rho) = 2 \operatorname{Im} \int_R^{c/\Omega} d\rho (L^2/\rho^2 - L^2\Omega^2/c^2)^{1/2} = 2L \ln(c/\Omega R)$$

$$W = w e^{-S(E=0)} = w (\Omega R/c)^{2L} = w (\omega R/cL)^{2L}$$

$$\omega = \Omega L \quad \text{Zeldovich resonance condition}$$

Vacuum instability in ergoregion

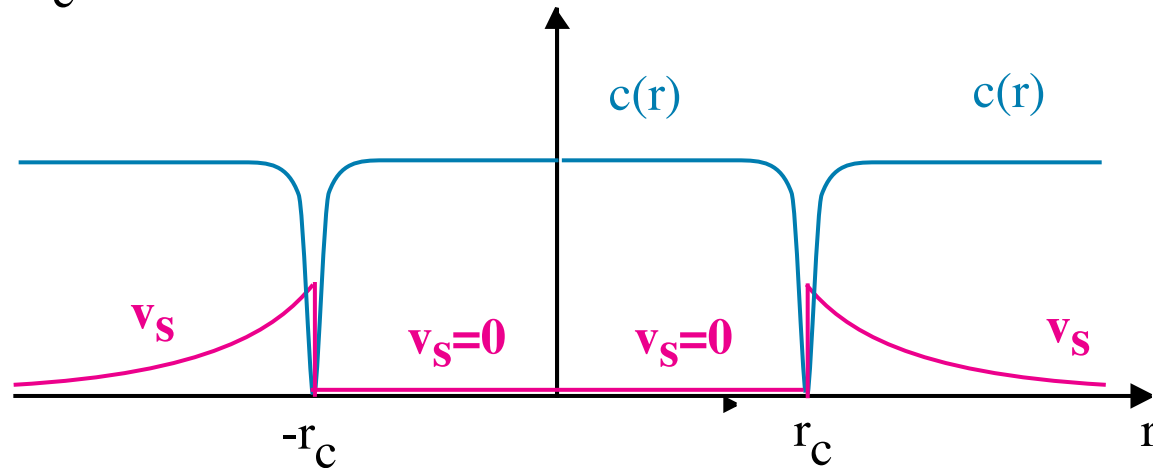
instability of ergoregion
in vortex core with $N \gg 1$:
no stable vacuum if $v_S > c$



$c < v_S$
 $g_{00} > 0$

reconstruction of vacuum
in the vortex core with $N \gg 1$:

ergoregion disappears, $v_S < c$ everywhere,
shell of Planck size and Planck energy
separates vacua with different $g_{\mu\nu}$

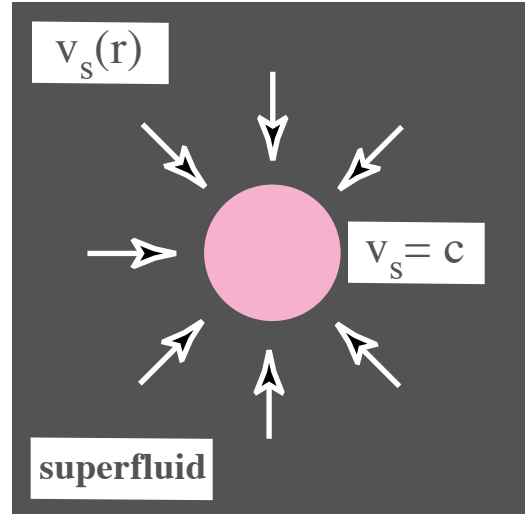
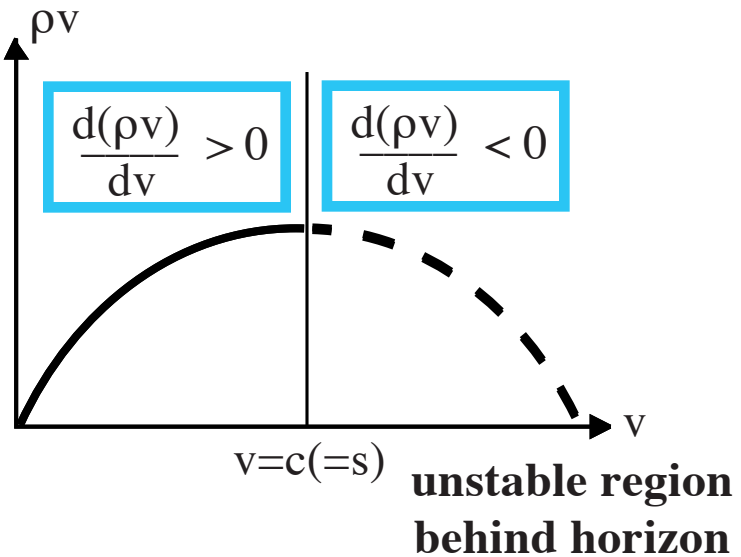


Vacuum resistance to formation of horizons

*spherical acoustic black or white holes
are not solutions of hydrodynamic equations*

equation along the stream line
of stationary flow:

$$\frac{d(\rho v)}{dv} = \rho(1 - v^2(r)/c^2)$$



horizon cannot be achieved
because continuity equation

$$\rho v = \text{Const} / r^2$$

requires

$$\frac{d(\rho v)}{dv} > 0$$

everywhere

Hydrodynamic instability is absent
if speed of "light" $c < s$ speed of sound

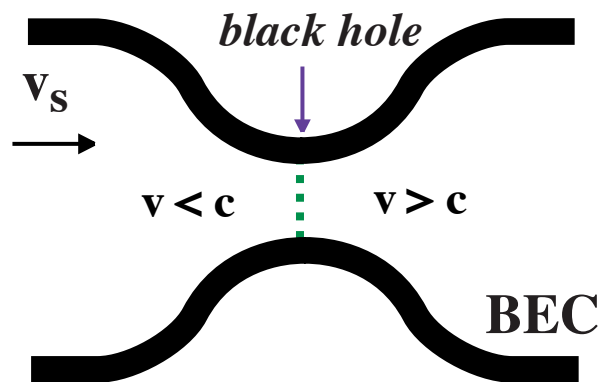
in Fermi superfluids $c \ll s$

possible horizons

Painleve-Gulstrand metric

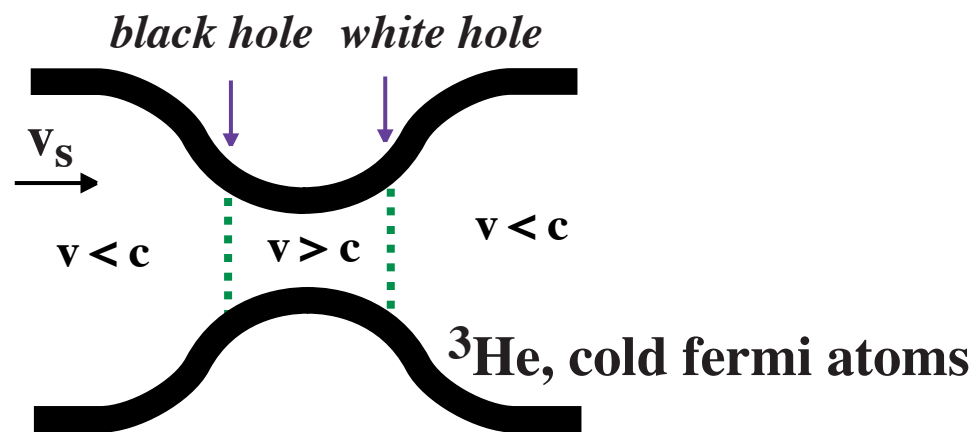
$$ds^2 = - dt^2 (c^2 - v^2) + 2 v dr dt + dr^2 + r^2 d\Omega^2$$

Acoustic horizon in Laval nozzle



horizon can be exactly in the middle

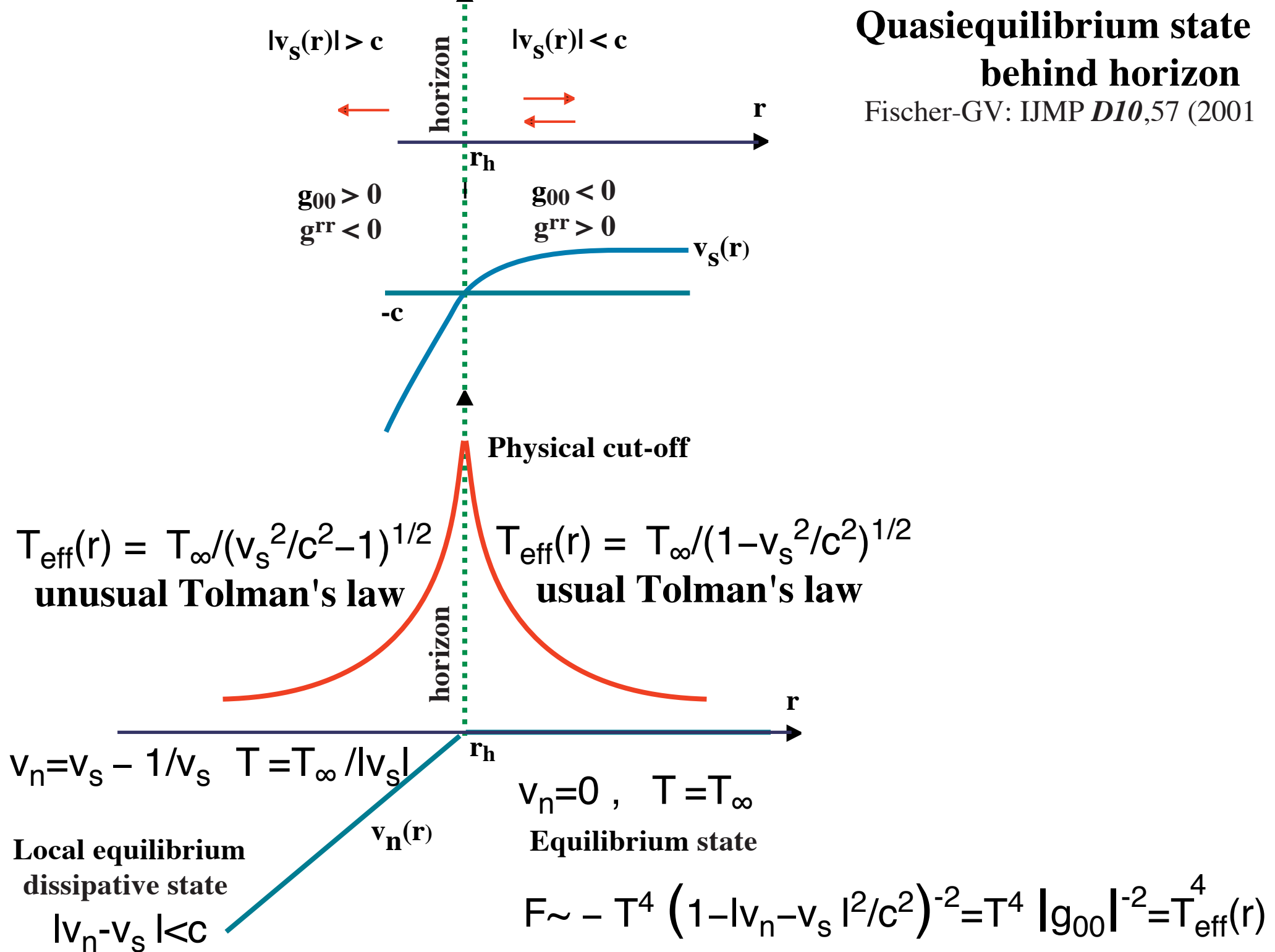
Horizons for fermionic quasiparticles



speed of "light" < speed of sound
no hydrodynamic instability

Quasiequilibrium state behind horizon

Fischer-GV: IJMP *D10*,57 (2001)



conclusion

* cond-mat sources:

sonic gravity, ripplon gravity, fermionic gravity

* messages from cond-mat to GR:

black hole is radiating, de Sitter is not GV: "On dS radiation via quantum tunneling" 0803.3367
but detector embedded in de Sitter sees radiation

vacuum resists to formation of horizon

vacuum may be unstable behind the horizon

* we need relativistic theory of quantum vacuum

black holes in dynamics of quantum vacuum

F.R. Klinkhamer and GV

"Self-tuning vacuum variable and cosmological constant", Phys. Rev. D **77**, 085015 (2008)

"Dynamic vacuum variable and equilibrium approach in cosmology", Phys. Rev. D; arXiv:0806.2805

"f(R) cosmology from q-theory", Pis'ma ZhETF **88**, 339 (2008)

* **quantum vacuum as self-sustained Lorentz invariant medium**

* **conserved vacuum charge**

* **thermodynamics of relativistic quantum vacuum**

* **dynamics of relativistic quantum vacuum**

* **application to cosmology:**

relaxation of cosmological constant from Planck scale to present value

* **future application to black hole**

physics of BH singularity

action

$$S = \int d^4x (-g)^{1/2} [\varepsilon(q) + K(q)R] + S_{\text{matter}}$$

gravitational coupling $K(q)$ is determined by vacuum
and thus depends on vacuum variable q

dynamic equations

equation
for q

$$d\varepsilon/dq + R dK/dq = \mu \quad \text{integration constant}$$

Einstein
equations

$$K(Rg_{\mu\nu} - 2R_{\mu\nu}) + (\varepsilon - \mu q)g_{\mu\nu} - 2(\nabla_{\mu}\nabla_{\nu} - g_{\mu\nu}\nabla^{\lambda}\nabla_{\lambda})K = T_{\mu\nu}$$

gravitational
coupling

Einstein
tensor

cosmological
term

matter

$$\nabla_{\mu} T^{\mu\nu} = 0$$