

Volume-preserving mean curvature flow of revolution hypersurfaces in a Rotationally Symmetric Space

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Abstract

In an ambient space with rotational symmetry around an axis (which include the Hyperbolic and Euclidean spaces), we study the evolution under the volume-preserving mean curvature flow of a revolution hypersurface M generated by a graph over the axis of revolution and with boundary in two totally geodesic hypersurfaces (tgh for short). Requiring that, for each time $t \geq 0$, the evolving hypersurface M_t meets such tgh ortogonally, we prove that: a) the flow exists while M_t does not touch the axis of rotation; b) throughout the time interval of existence, b1) the generating curve of M_t remains a graph, and b2) the averaged mean curvature is double side bounded by positive constants; c) the singularity set (if non-empty) is finite and lies on the axis; d) under a suitable hypothesis relating the enclosed volume to the n -volume of M , we achieve long time existence and convergence to a revolution hypersurface of constant mean curvature.

1 Introduction and definition of the ambient space

Let $\{X_t : M \longrightarrow \overline{M}\}_{t \in [0, T[}$, be a smooth family of immersions of a compact n -dimensional manifold (may be with boundary) M into a $(n+1)$ -dimensional Riemannian manifold \overline{M} . By “smooth family” we understand that the map $X : M \times [0, T[\longrightarrow \overline{M}$ defined by $X(p, t) = X_t(p)$ is smooth. Let N_t be the outward unit normal vector of the immersion X_t and H_t the trace of the Weingarten map L_{N_t} of X_t associated to N_t (with the convention that H_t is n times the usual mean curvature with the sign which makes positive the mean curvature of a round sphere in \mathbb{R}^{n+1}). By M_t we shall denote both the immersion $X_t : M \longrightarrow \overline{M}$ and the image $X_t(M)$, as well as the Riemannian manifold (M, g_t) with the metric g_t induced by the immersion. Analogous notation will be used when we have a single immersion $X : M \longrightarrow \overline{M}$. We are working with embeddings, then there is no confusion about considering objects on M or on $X_t(M)$.

The *volume preserving mean curvature flow* is defined as a solution of the equation

$$\frac{\partial X_t}{\partial t} = (\overline{H}_t - H_t) N_t, \quad (1.1)$$

where \overline{H}_t is the averaged mean curvature $\overline{H}_t = \frac{\int_{M_t} H_t dv_{g_t}}{\int_{M_t} dv_{g_t}}$, being dv_{g_t} the volume element on M_t . This flow decreases the area of M_t , but preserves the volume of the domain Ω_t enclosed by M_t (when such Ω_t exists).

A nearest relative of (1.1) is the mean curvature flow $\frac{\partial X_t}{\partial t} = -H_t N_t$, which has been extensively studied in Euclidean and non Euclidean ambient spaces. The work for the flow (1.1) is much shorter and mainly devoted to the evolution of hypersurfaces in the Euclidean space. Special attention has been paid, in the Euclidean or Minkowski ambient spaces, to the evolution of convex hypersurfaces ([15] and [3]) and hypersurfaces in the neighborhood of a round sphere ([12]). The results for the last situation were extended in [2] to general Riemannian manifolds whose scalar curvature has nondegenerate critical points. Results both for the evolution of h -convex hypersurfaces and hypersurfaces which are near a geodesic sphere were obtained in [7] for the Hyperbolic space.

As far as we know, few other evolutions under (1.1) have been studied. This makes the works [5] and [6] by M. Athanassenas specially interesting and valuable. Indeed, she deals with two extra difficulties: lack of convexity or sphere closeness assumptions and boundary conditions. In particular, she considers the following

Setting Euc ([5] and [6]). *Let M be a smoothly embedded hypersurface in \mathbb{R}^{n+1} contained in the domain $G = \{x \in \mathbb{R}^{n+1} : 0 \leq x^{n+1} \leq d; d > 0\}$, and with boundary $\partial M \subset \partial G$. Assume that M is a revolution hypersurface around the axis $z := x^{n+1}$ generated by the graph of a function $r(z)$, that it intersects ∂G orthogonally at ∂M and encloses (inside G) a $(n+1)$ -volume V .*

Athanassenas' papers study the behavior of the flow defined by (1.1) with M in the setting Euc as initial condition and the boundary condition that the boundary ∂M_t of the evolving surface M_t remains in ∂G and meets it orthogonally. Next we recall her main results:

- A1.** The flow exists for all time $t > 0$ and converges to the cylinder $C \subset G$ enclosing a volume V under the assumption $n\text{-volume}(M) \leq \frac{V}{d}$ (cf. [5]).
- A2.** The singular set (if not empty) is finite and is located along the axis of rotation (cf. [6]).

In the present paper, we face the challenge of extending the above statements to hypersurfaces living in a larger universe; in other words, our goal is to export the Euclidean results to more general ambient spaces. But not every Riemannian manifold is well suited to deal with the situation analogous to the setting Euc. Indeed, in order the concept of revolution hypersurface makes sense, we need that the ambient space also has some kind of rotational symmetry. This motivates the introduction of the spaces defined below.

Definition 1 *A rotationally symmetric space with respect to an axis z is a manifold $\overline{M} = J \times B_R^n$, with J an interval in \mathbb{R} and B_R^n the Euclidean open ball of radius R in \mathbb{R}^n , endowed with a C^∞ metric of the form*

$$\overline{g} := dr^2 + f^2(r) dz^2 + h^2(r) g_S, \quad (1.2)$$

in the cylindrical coordinates $r = |x|$, $z = z$, $u = x/|x|$ of $(z, x) \in \overline{M}$, where g_S is the standard metric of sectional curvature 1 on the sphere S^{n-1} .

On such a manifold there is a natural action of $SO(n)$ in the following way. Given a point $q = (z, x)$ and $\mathcal{R} \in SO(n)$, we define $\mathcal{R}q := (z, \mathcal{R}x) \equiv \exp_p \mathcal{R} \exp_p^{-1} q$, where let $p = (z, 0)$ and \exp_p is the exponential map in $(\overline{M}, \overline{g})$. The axis z is invariant under such action, and this is the reason for the name of these spaces. From now on, we shall call the axis z the **axis of rotation**.

Notice that the previous definition includes the three model spaces with constant sectional curvature. In fact,

- When $R = \infty$, $J = \mathbb{R}$, $\lambda < 0$, $f(r) = \cosh(\sqrt{|\lambda|} r)$, $h(r) = \frac{1}{\sqrt{|\lambda|}} \sinh(\sqrt{|\lambda|} r)$, $(\overline{M}, \overline{g})$ is the Hyperbolic space of sectional curvature λ .
- When $\lambda > 0$, $R = \frac{\pi}{2\sqrt{\lambda}}$, $J =] -\frac{\pi}{2\sqrt{\lambda}}, \frac{\pi}{2\sqrt{\lambda}}[$, $f(r) = \cos(\sqrt{\lambda} r)$, $h(r) = \frac{1}{\sqrt{\lambda}} \sin(\sqrt{\lambda} r)$, $(\overline{M}, \overline{g})$ is the open half sphere of sectional curvature λ .
- When $R = \infty$, $J = \mathbb{R}$, $f(r) = 1$, $h(r) = r$, $(\overline{M}, \overline{g})$ is the Euclidean Space.

According to the previous definition, a natural setting for the extension of the results A1 and A2 mentioned above is:

Setting RSS. *Let M be a smoothly embedded hypersurface of revolution around the axis of rotation in \overline{M} , generated by the graph of a function $r(z)$ and contained in the domain $G = \{(r, z, u) \in \overline{M} : a \leq z \leq b\}$, with boundary ∂M , which intersects ∂G orthogonally at the boundary and encloses a $(n+1)$ -volume V inside G .*

However, as we pointed out in [7], the negative ambient curvature seems to be a friendlier setting for the flow (1.1), so here we still need suitable hypotheses on the negativity of some sectional curvatures. In particular, our main results are valid for the following

Setting RSS2. *Let M be in the setting RSS and assume that either*

- (a) *the curvatures of \overline{M} satisfy $S_{zi} < 0$ and $S_{ri} \leq 0$, or*
- (b) *\overline{M} is the Euclidean space.*

Here (cf. notation of section 2), for each i , S_{zi} means the sectional curvature of the plane generated by the axis of rotation and a direction E_i orthogonal to the plane π_G containing the generating curve, and S_{ri} is the sectional curvature given by E_i and the direction in π_G orthogonal to the axis.

In this setting, flow M by (1.1) with the boundary condition that the evolving hypersurface

$$M_t \text{ intersects } G \text{ orthogonally at the boundary for every } t. \quad (1.3)$$

We shall prove that

B1. The flow exists while the evolving hypersurface M_t does not touch the axis of rotation (Theorem 6).

- B2.** As long as the flow exists, the generating curve remains a graph over the axis of rotation (Theorem 5).
- B3.** Throughout the same time, the averaged mean curvature \overline{H} is positive and double side bounded (Proposition 9).
- B4.** If the motion M_t of M has a first singularity at time T , the singular set is finite and lies on the axis of rotation (Theorem 10).
- B5.** If $\text{vol}(M)$ is small enough compared with the volume V enclosed by M inside G , then the solution of (1.1) is defined for all $t > 0$ and converges to a revolution hypersurface of constant mean curvature in \overline{M} (Theorem 12).

At this moment, we think that the non positivity of the sectional curvature is an essential hypothesis for B2, but we have no concluding counterexample. The other results may still be true in an ambient space with positive curvature, but this would require a different kind of argument, because our proofs rely strongly on B2.

About the techniques used to prove the above results, this paper follows the ideas introduced in [1] and [19] (which study the mean curvature flow for closed rotationally symmetric hypersurfaces) with the clever modifications proposed by M. Athanassenas to deal with the averaged mean curvature, and some methods learned from [9].

But these ideas are not enough to work in the setting RSS2, where we actually have to deal with an ambient space non (necessarily) flat and with sectional curvatures non (necessarily) constant. The key to overcome this further difficulty is the mixture of some ideas from the aforementioned works with a bunch of tricky technicalities which spread out all over our paper.

In addition, we had to face an extra and unexpected complication: as far as we are aware, there is a couple of mistakes in [5], [6] (in the last addend of the first displayed formula of the proof of Proposition 5 on page 63 of [5] and on the last displayed expression on page 10 of [6]); and, at least for the mistake in the computation of [6], when doing it properly, the arguments that follow fail. This makes necessary (and not superfluous by far) the inclusion of the Euclidean space within the setting RSS2. In short, we not only extend the desired results to a wider setting, but we also complete the proof in the Euclidean space.

The paper is organized as follows. We begin with the study of the main features of the spaces introduced in Definition 1 (section 2). After some remarks on short time existence and basic computations (section 3), we find bounds for the distance to the axis of rotation and for \overline{H} (section 4) which will be the key to prove B2 (section 5). Section 6 is devoted to the proof of B1, and in section 7 we obtain B3 together with finer bounds for \overline{H} . We shall use these finer bounds to prove B4 (section 8). Finally, B5 is proved in section 9.

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2 Rotationally symmetric spaces around an axis

To compute the covariant derivatives and curvatures corresponding to the metric (1.2), we shall use the orthonormal frame

$$E_0 := E_r = \frac{\partial}{\partial r}, \quad E_1 := E_z = \frac{1}{f(r)} \frac{\partial}{\partial z}, \quad E_i = \frac{1}{h(r)} e_i, \quad i = 2, \dots, n \quad (2.1)$$

where $\{e_i\}$ is an orthonormal frame of S^{n-1} with its standard metric, and the dual orthonormal frame

$$\theta^r = dr, \quad \theta^z = f(r) dz, \quad \theta^i = h(r) \delta^i, \quad (2.2)$$

being $\{\delta^i\}$ the dual frame of $\{e_i\}$ in S^{n-1} with its standard metric. In these frames, the Cartan connection forms ω_a^b (defined by $d\theta^b = -\sum_{a=0}^n \omega_a^b \wedge \theta^a$) are

$$\omega_r^z = \frac{f'(r)}{f(r)} \theta^z, \quad \omega_r^i = \frac{h'(r)}{h(r)} \theta^i, \quad \omega_z^i = 0, \quad \omega_j^i = {}^S\omega_j^i, \quad (2.3)$$

where ${}^S\omega_j^i$ are the connection 1-forms of S^{n-1} with its standard metric and $'$ denotes the derivative respect to r .

From the above expressions, the covariant derivatives of the orthonormal frame (given by $(\bar{\nabla}_X E_a = \sum_{b=0}^n \omega_a^b(X) E_b)$) are

$$\begin{aligned} \bar{\nabla}_{E_r} E_r &= 0, & \bar{\nabla}_{E_z} E_r &= \frac{f'(r)}{f(r)} E_z, & \bar{\nabla}_{E_i} E_r &= \frac{h'(r)}{h(r)} E_i, \\ \bar{\nabla}_{E_r} E_z &= 0, & \bar{\nabla}_{E_z} E_z &= -\frac{f'(r)}{f(r)} E_r, & \bar{\nabla}_{E_i} E_z &= 0, \\ \bar{\nabla}_{E_r} E_i &= 0, & \bar{\nabla}_{E_z} E_i &= 0, & \bar{\nabla}_{E_j} E_i &= -\frac{h'(r)}{h(r)} \delta_{ij} E_r + {}^S\omega_j^k(E_i) E_k. \end{aligned} \quad (2.4)$$

Also from (2.3) we can compute the curvature forms Ω_a^b (which satisfy $\Omega_a^b = d\omega_a^b - \sum_{c=0}^n \omega_a^c \wedge \omega_c^b$):

$$\begin{aligned} \Omega_r^z &= \frac{f''}{f} dr \wedge \theta^z, & \Omega_r^i &= \frac{h''}{h} dr \wedge \theta^i, \\ \Omega_z^i &= -\frac{h'f'}{hf} \theta^i \wedge \theta^z, & \Omega_j^i &= {}^S\Omega_j^i - \left(\frac{h'}{h}\right)^2 \theta^i \wedge \theta^j, \end{aligned} \quad (2.5)$$

where ${}^S\Omega_j^i$ are the curvature 2-forms of S^{n-1} with its standard metric.

Let us remark a fact which is necessary to have in mind when doing computations:

$${}^S\omega_j^i(E_k) = \frac{1}{h} {}^S\omega_j^i(e_k) \quad {}^S\Omega_j^i(E_k, E_\ell) = \frac{1}{h^2} {}^S\Omega_j^i(e_k, e_\ell).$$

This is interesting to recall because ${}^S\omega_j^i(e_k)$ and ${}^S\Omega_j^i(e_k, e_\ell)$ give the standard values in S^{n-1} with its standard metric.

Remark 1 In any Riemannian manifold \mathcal{M} , given a point $p \in \mathcal{M}$ and a geodesic Γ through p , the eigenvalues of the Weingarten map of the geodesic tube of radius r around Γ approach ∞ as $1/r$ when $r \rightarrow 0$ in the directions tangent to geodesic discs orthogonal to Γ (cf. [17]). Notice that in the metric (1.2) the Weingarten map of such a tube in the directions tangent to geodesic discs is given by $\bar{\nabla}_{E_i} E_r$. Then, by formulae (2.4), $\lim_{r \rightarrow 0} \frac{\frac{h'(r)}{h(r)}}{\frac{1}{r}} = 1$.

On the other hand, $h(r)$ is the norm of the vector e_i tangent to the geodesic sphere in the hyperplane $z = \text{constant}$ with radius r , then $h(0) := \lim_{r \rightarrow 0} h(r) = 0$ with $\lim_{r \rightarrow 0} \frac{h(r)}{r} = 1$, which combined with the above limit gives $h'(0) = 1$. Moreover, $f(r)$ is the norm of the vector tangent to the curve $z \mapsto (r_0, z, u_0)$, which is a geodesic when $r_0 = 0$. Therefore, $f(0) > 0$ and constant, and we can take, without loss of generality, $f(0) = 1$.

If we do not have any singularity in our manifold $\bar{M} = J \times B_R^n$ at the points $(z, 0)$, the limit of $\bar{\nabla}_{E_z} E_z$ must be independent on the direction we approach the point $(z, 0)$. But from (2.4) we have opposite vectors $\bar{\nabla}_{E_z} E_z$ of norms $|\frac{f'(0)}{f(0)}|$ when we approach $(z, 0)$ by opposite directions, then we must have $f'(0) = 0$.

From all these remarks we have that, if the metric \bar{g} given by (1.2) is a Riemannian metric, then

$$f(0) = 1, \quad f'(0) = 0, \quad h(0) = 0, \quad h'(0) = 1. \quad (2.6)$$

Now we use the expression $R_{abcd} = \Omega_b^a(E_c, E_d)$ of the curvature components and (2.5) to obtain the sectional curvatures:

$$\begin{aligned} S_{rz} = R_{r z r z} &= -\frac{f''}{f}, & S_{ri} = R_{r i r i} &= -\frac{h''}{h}, \\ S_{zi} = R_{z i z i} &= -\frac{h' f'}{h f}, & S_{ij} = R_{i j i j} &= \frac{1 - h'^2}{h^2}. \end{aligned} \quad (2.7)$$

In the next two remarks, we translate the curvature conditions of the setting RSS2 in terms of the functions f and h defining \bar{g} .

Remark 2 If $S_{ri} \leq 0$ the second formula in (2.7) implies $h'' \geq 0$, then h' is non-decreasing and, since $h'(0) = 1$, $h' > 0$. Thus h is increasing.

Henceforth, we shall denote by $z(\phi)$ the first positive zero of a function $\phi : \mathbb{R} \rightarrow \mathbb{R}$.

Remark 3 If $S_{zi} \leq 0$ (resp. $S_{zi} < 0$) the third formula in (2.7) implies $f' h'(r) \geq 0$ (resp. $f' h'(r) > 0$) for every r and, since $h'(r) > 0$ for $r \in [0, z(h')]$, $f'(r) \geq 0$ (resp. $f'(r) > 0$) for r in this interval. So f is non-decreasing, (resp. increasing) in particular, $f(r) \geq f(0) = 1$ (resp. $f(r) > f(0) = 1$). The property $f'(r) \geq 0$, together with $f'(0) = 0$, implies $f''(r) \geq 0$ (resp. $f''(r) > 0$) at some points of $[0, z(h')]$, and $S_{rz} \leq 0$ (resp. $S_{rz} < 0$) at those points.

If, moreover, $S_{ri} \leq 0$, Remark 2 implies $z(h') = \infty$, then the inequalities stated in the last paragraph for an interval now hold for every $r > 0$.

3 Short time existence and some formulae

Let us begin with a remark on the notation. Along all the paper we use many quantities (functions, vector and tensor fields, ...) depending on the evolved hypersurface M_t of M under the flow (1.1). When they are introduced for the first time, we write the subindex t to denote their dependence on t , but later we shall omit it unless it is not clear from the context that we are making reference to the function depending on t . Sometimes we also denote the dependence on t by (\cdot, t) .

Observe that, since M is invariant by the action of $SO(n)$ as a group of isometries of \overline{M} , also M_t will be. Then the unit normal vector N_t to M_t will be contained in the plane generated by E_r and E_z and can be written as

$$N = \langle N, E_r \rangle E_r + \langle N, E_z \rangle E_z, \quad (3.1)$$

in turn, the unit vector \mathbf{t}_t tangent to the generating curve will be

$$\mathbf{t} = -\langle N, E_z \rangle E_r + \langle N, E_r \rangle E_z. \quad (3.2)$$

We shall use the coordinates (r_t, z_t, u_t) for M_t . Since M is generated by the graph of a function $r(z)$, for small t a solution of (1.1) will still be generated by the graph of a function $r_t(z_t)$. Using this function, the vectors \mathbf{t} and N can be written as

$$\mathbf{t} = \frac{1}{\sqrt{\dot{r}^2 + f^2}}(\dot{r}E_r + fE_z), \quad N = \frac{1}{\sqrt{\dot{r}^2 + f^2}}(fE_r - \dot{r}E_z), \quad (3.3)$$

where \dot{r} denotes the derivative of r_t with respect to z_t .

A convenient orthonormal frame of M_t is given by $\mathbf{t}, E_2, \dots, E_n$. If we compute the mean curvature of M_t using this frame, and denote $\tilde{\mathbf{t}} = \sqrt{\dot{r}^2 + f^2} \mathbf{t}$, $\tilde{N} = \sqrt{\dot{r}^2 + f^2} N$, we get

$$H = k_1 + (n-1)k_2, \quad (3.4)$$

where

$$k_1 = -\frac{1}{(\sqrt{\dot{r}^2 + f^2})^3} \langle \nabla_{\tilde{\mathbf{t}}} \tilde{N} \rangle = \frac{1}{\sqrt{\dot{r}^2 + f^2}} \left(\frac{-\ddot{r}f + \dot{r}^2 f'}{\dot{r}^2 + f^2} + f' \right) \quad (3.5)$$

$$= \frac{1}{\sqrt{\dot{r}^2 + f^2}} \left(-\frac{d}{dz} \arctan \left(\frac{\dot{r}}{f} \right) + f' \right) \quad (3.6)$$

is the normal curvature of M_t in the direction of \mathbf{t} , and

$$k_2 = \frac{1}{\sqrt{\dot{r}^2 + f^2}} \langle \nabla_{E_2} \tilde{N}, E_2 \rangle = \frac{f}{\sqrt{\dot{r}^2 + f^2}} \frac{h'}{h} \quad (3.7)$$

is the normal curvature of M_t in the direction of E_i , $i = 2, \dots, n$.

It is well known (cf. [11]) that, up to tangential diffeomorphisms, equation (1.1) is equivalent to

$$\left\langle \frac{\partial X_t}{\partial t}, N_t \right\rangle = \overline{H} - H. \quad (3.8)$$

In this flow the variable z does not change with time, and formulae (3.3), (3.4), (3.5) and (3.7) are true just by this change. Using them, equation (3.8) becomes

$$\frac{\partial r}{\partial t} = \frac{\ddot{r}}{\dot{r}^2 + f^2} - \frac{f'(r)}{f(r)} \left(1 + \frac{\dot{r}^2}{\dot{r}^2 + f^2} \right) - (n-1) \frac{h'(r)}{h(r)} + \overline{H} \sqrt{1 + \frac{\dot{r}^2}{f^2}}. \quad (3.9)$$

Replacing \overline{H} in (3.9) by any $C^{1,\alpha/2}$ real valued function ψ such that $\psi(0) = \overline{H}(0)$, we obtain a parabolic equation which, at least for small t , has a unique solution satisfying $\dot{r}(a) = \dot{r}(b) = 0$. Now, using a routine fixed point argument (cf. [18]), we can establish short time existence also for (3.9) with the same boundary conditions.

4 Upper bounds for r and rough bounds for \overline{H}

In this section, we shall prove that if M_t is a maximal solution of (1.1) with initial condition in the setting RSS and satisfying the boundary condition (1.3), then r_t has a finite upper bound and \overline{H} is positive and bounded from above. Under additional curvature assumptions, which include both situations of the setting RSS2, we shall show that \overline{H} is actually bounded away from zero.

The aforementioned bounds will be the key to prove the conservation of the property of being a graph for the generating curve (cf. section 5). This and the results of section 6 will help us to obtain more accurate bounds for \overline{H} (cf. section 7).

Let us define the functions β and δ and the numbers $r_2 > r_1 > 0$ by

$$\beta(r) = \int_0^r f(r)h(r)^{n-1}dr, \quad r_1 = \beta^{-1} \left(\frac{V}{(b-a)\sigma} \right), \quad (4.1)$$

$$\delta(r) = \int_0^r h(r)^{n-1}dr, \quad r_2 = \delta^{-1} \left(\frac{\text{vol}(M)}{\sigma} + \delta(r_1) \right), \quad (4.2)$$

where σ denotes the volume of S^{n-1} with its standard metric. Let us remark that the functions β and δ are increasing, thus they have inverse and $r_2 > r_1$ as claimed. Moreover, r_1 and r_2 depend only on $b-a$, V , $\text{vol}(M)$ and the metric \overline{g} of the ambient space \overline{M} .

Proposition 1 *If M is in the setting RSS and $[0, T[$ is the maximal time interval where the flow (1.1) satisfying (1.3) is defined, then $r_t < r_2$ for every $t \in [0, T[$.*

Proof From the expression (1.2) of the metric of \overline{M} it follows that its volume element is

$$\overline{\mu} = f(r)h(r)^{n-1}dr \wedge dz \wedge {}^S\mu, \quad (4.3)$$

being ${}^S\mu$ the volume element on S^{n-1} defined by its standard metric.

Using expression (4.3), the volume V enclosed by M_t can be computed as

$$V = \sigma \int_a^b \int_0^{r(z)} f(r)h(r)^{n-1}dr \, dz, \quad (4.4)$$

and the definition of r_1 gives

$$V = \sigma \int_a^b \int_0^{r_1} f(r) h(r)^{n-1} dr \, dz. \quad (4.5)$$

Comparing these formulae, we deduce that $\min_{z \in [a,b]} r(z) =: r_m \leq r_1 \leq r_M := \max_{z \in [a,b]} r(z)$ and one of the inequalities is strict if and only if the other is.

The induced volume form μ_t on M_t is defined by the contraction

$$\mu_t = \iota_{N_t} \bar{\mu} = \sqrt{\dot{r}^2(z) + f^2(r(z))} \, h(r(z))^{n-1} dz \wedge {}^S \mu, \quad (4.6)$$

where we have used (2.2), (3.3) and (4.3). Hence

$$\begin{aligned} \text{vol}(M_t) &= \sigma \int_a^b \sqrt{\dot{r}^2(z) + f^2(r(z))} \, h(r(z))^{n-1} dz > \sigma \int_a^b |\dot{r}| \, h(r(z))^{n-1} dz \\ &\geq \sigma \int_{r_m}^{r_M} h(r)^{n-1} dr = \sigma \left(\int_{r_m}^{r_1} h(r)^{n-1} dr + \int_{r_1}^{r_M} h(r)^{n-1} dr \right). \end{aligned} \quad (4.7)$$

From this inequality we obtain

$$\text{vol}(M_t) > \sigma \int_{r_1}^{r_M} h(r)^{n-1} dr = \sigma(\delta(r_M) - \delta(r_1)). \quad (4.8)$$

Then, having into account that $\text{vol}(M_t) < \text{vol}(M)$, we reach

$$\delta(r_M) < \frac{\text{vol}(M)}{\sigma} + \delta(r_1),$$

from which the proposition follows. \square

Next, the goal is to bound the averaged mean curvature \bar{H}_t .

Proposition 2 *Let M_t be the solution of (1.1) with initial condition M in the setting RSS and satisfying (1.3). For every t such that*

$$0 < \rho \leq r_t \leq z_0 := \min\{z(h'), z((n-1)h'f + f'h)\} \quad \text{for some fixed } \rho,$$

there is a constant $h_2(V, \bar{g}, n, b-a, \text{vol}(M), \rho) > 0$ such that $0 \leq \bar{H} \leq h_2$.

Proof Before starting the proof, let us remark that

- By (2.6), we have $\min\{z(h'), z((n-1)h'f + f'h)\} > 0$.
- The hypothesis on the variation of r_t together with Proposition 1 imply that $r_t \leq \mathfrak{r} = \min\{r_2, z_0\}$.

From (3.4), (3.6) and (3.7) we can write

$$\overline{H}_t = \frac{1}{\text{vol}(M_t)} \int_{M_t} \frac{1}{\sqrt{\dot{r}^2 + f^2}} \left(-\frac{d}{dz} \arctan \left(\frac{\dot{r}}{f} \right) + (n-1) \frac{h'}{h} f + f' \right) \mu_t = I_1 + I_2, \quad (4.9)$$

where we have (by (4.6)) the following expressions for I_1 and I_2 :

$$I_1 = \frac{\sigma}{\text{vol}(M_t)} \int_a^b -\frac{d}{dz} \arctan \left(\frac{\dot{r}}{f} h(r(z))^{n-1} \right) dz \quad \text{and} \quad (4.10)$$

$$I_2 = \frac{1}{\text{vol}(M_t)} \int_{M_t} \frac{1}{\sqrt{\dot{r}^2 + f^2}} \frac{(n-1)h'f + hf'}{h} \mu_t. \quad (4.11)$$

Integrating by parts and realizing that the condition in the boundary gives $\dot{r}(b) = \dot{r}(a) = 0$, we get

$$I_1 = \frac{\sigma}{\text{vol}(M_t)} \int_a^b \arctan \left(\frac{\dot{r}}{f} \right) (h^{n-1})' \dot{r} dz \geq 0 \quad (4.12)$$

because $\arctan \left(\frac{\dot{r}}{f} \right) \dot{r} \geq 0$. Since $\arctan \left(\frac{\dot{r}}{f} \right) \dot{r} \leq \frac{\pi}{2} \sqrt{\dot{r}^2 + f^2}$,

$$I_1 \leq \frac{\sigma}{\text{vol}(M_t)} \frac{\pi}{2} \int_a^b \sqrt{\dot{r}^2 + f^2} (h^{n-1})' dz \quad (4.13)$$

$$\begin{aligned} &= \frac{\pi}{2\text{vol}(M_t)} \int_{S^{n-1}} \int_a^b \frac{(h^{n-1})'}{h^{n-1}} h^{n-1} \sqrt{\dot{r}^2 + f^2} dz S_\mu \\ &= \frac{\pi}{2\text{vol}(M_t)} \int_{M_t} \frac{(h^{n-1})'}{h^{n-1}} \mu_t \leq \frac{\pi}{2} \max_{r \in [\rho, \mathfrak{v}]} \frac{(h^{n-1})'}{h^{n-1}}(r) < \infty. \end{aligned} \quad (4.14)$$

From (4.6) and (4.11) we have

$$I_2 = \frac{\sigma}{\text{vol}(M_t)} \int_a^b ((n-1)h'f + f'h) h^{n-2} dz \geq 0. \quad (4.15)$$

On the other hand, (4.11) also implies that

$$I_2 \leq \frac{1}{\text{vol}(M_t)} \int_{M_t} \frac{1}{f} \frac{(n-1)h'f + f'h}{h} \mu_t \leq \max_{r \in [\rho, \mathfrak{v}]} \frac{(n-1)h'f + f'h}{fh}(r) < \infty. \quad (4.16)$$

In conclusion, the existence of the finite upper bound h_2 follows from (4.14) and (4.16). Moreover, the non-negativity of \overline{H} is due to (4.12) and (4.15), which in turn are true thanks to the assumption $r_t \leq z_0$. \square

Corollary 3 *Let M_t be the solution of (1.1) with initial condition M in the setting RSS and satisfying (1.3). Let us suppose that the sectional curvatures S_{z_i} and S_{r_i} of \overline{M} are non-positive. For every t such that $0 < \rho \leq r_t$ for some fixed ρ , there are constants $h_i(V, \overline{g}, n, b - a, \text{vol}(M), \rho) > 0$, $i = 1, 2$ such that $h_1 \leq \overline{H} \leq h_2$.*

Proof From Remark 3 it follows that the hypotheses on the curvature of \overline{M} imply $z_0 = \infty$. Then the upper bound is an immediate consequence of Proposition 2. Remarks 2 and 3 tell us that the hypotheses on the curvature imply that h, h' are non decreasing, $f \geq 1$ and $f'h \geq 0$, therefore, from (4.15) we obtain

$$I_2 \geq \frac{\sigma(b-a)}{\text{vol}(M_t)} \min_{r \in [\rho, r_2]} ((n-1)h'f + f'h) h^{n-2}(r) \geq \frac{\sigma(b-a)}{\text{vol}(M)} (n-1) (h'h^{n-2})(\rho) > 0, \quad (4.17)$$

which gives the positive lower bound for \overline{H} . \square

5 Preserving the property of being a graph for the generating curve

This section is devoted to prove that, for M in the setting RSS2, the evolving hypersurface remains a revolution hypersurface generated by a smooth graph. As we pointed out above, M_t is always a revolution hypersurface. Then the aim is to show that the generating curve remains a graph over the axis of rotation for all the time.

In this section we continue with the notation used in sections 2 and 3. Moreover, by α we represent the second fundamental form of M_t , with the sign convention $\alpha(X, Y) = \langle \overline{\nabla}_X N, Y \rangle$, and L denotes its Weingarten map (defined by $LX = \overline{\nabla}_X N$). ∇ means the intrinsic covariant derivative on M_t .

Recall that the generating curve is a graph if and only if $u := \langle N, E_r \rangle > 0$, which is equivalent to say $1 < v = \frac{1}{u} < \infty$. Therefore, our goal is to obtain an upper bound for v . To achieve this, we need the evolution equation for v .

Lemma 4 Under (1.1), $v = \langle N, E_r \rangle^{-1}$ evolves as

$$\begin{aligned} \frac{\partial v}{\partial t} = & \Delta v - \frac{2}{v} |dv|^2 - \left(\frac{1}{v} \frac{f'}{f} - k_1 \right)^2 v + \overline{H} \frac{f'}{f} (1 - v^2) \\ & - \left(\frac{f''}{f} - 2 \left(\frac{f'}{f} \right)^2 + (n-1) \left(\left(\frac{h'}{h} \right)' - \frac{h'f'}{hf} \right) \right) \left(1 - \frac{1}{v^2} \right) v. \end{aligned} \quad (5.1)$$

Proof First we compute Δu . To do so, we shall use an orthonormal frame of M of the form $\mathbf{t}, E_2, \dots, E_n$. It follows from (2.4), (3.1) and (3.2) (compare with (3.7)) that

$$\overline{\nabla}_{E_i} N = k_2 E_i = u \frac{h'}{h} E_i, \quad (5.2)$$

$$E_i u = 0, \quad \nabla_{\mathbf{t}} \mathbf{t} = 0, \quad (5.3)$$

$$\overline{\nabla}_{\mathbf{t}} E_r = u \frac{f'}{f} E_z, \quad \overline{\nabla}_{\mathbf{t}} N = k_1 \mathbf{t}. \quad (5.4)$$

From (5.3) we have

$$\Delta u = \mathfrak{t}u - \sum_{i=2}^n (\nabla_{E_i} E_i)u, \quad (5.5)$$

and, using formulae (5.4), (2.4), (3.2) and $\langle N, E_z \rangle^2 = 1 - u^2$, we obtain

$$\begin{aligned} \mathfrak{t}u &= \left(u \frac{f'}{f} - k_1 \right) \sqrt{1 - u^2}, \\ \mathfrak{t}u &= \left(u \frac{f'}{f} - k_1 \right) \frac{f'}{f} (1 - u^2) - \left(u \frac{f'}{f} - k_1 \right)^2 u \\ &\quad - u \left(\frac{f'}{f} \right)' (1 - u^2) - \mathfrak{t}k_1 \sqrt{1 - u^2}, \end{aligned} \quad (5.6)$$

$$(\nabla_{E_i} E_i)u = \langle \nabla_{E_i} E_i, \mathfrak{t} \rangle \mathfrak{t}u = \frac{h'}{h} (1 - u^2) \left(u \frac{f'}{f} - k_1 \right). \quad (5.7)$$

Joining (5.5), (5.6) and (5.7), we reach the desired expression for Δu :

$$\begin{aligned} \Delta u &= \left(u \frac{f'}{f} - k_1 \right) \frac{f'}{f} (1 - u^2) - \sqrt{1 - u^2} \mathfrak{t}k_1 - \left(u \frac{f'}{f} - k_1 \right)^2 u \\ &\quad - u \left(\frac{f'}{f} \right)' (1 - u^2) - (n - 1) \left(u \frac{f'}{f} - k_1 \right) \frac{h'}{h} (1 - u^2). \end{aligned} \quad (5.8)$$

From the standard evolution formula $\frac{\overline{\nabla} N}{\partial t} = \text{grad} H$ for (1.1), using again (2.4) and (3.2), we get

$$\frac{\partial}{\partial t} u = \langle \text{grad} H, E_r \rangle + (\overline{H} - H) \langle N, E_z \rangle^2 \frac{f'}{f}$$

but

$$\begin{aligned} \langle \text{grad} H, E_r \rangle &= (\mathfrak{t}k_1 + (n - 1)\mathfrak{t}k_2) \langle \mathfrak{t}, E_r \rangle \\ &= -\sqrt{1 - u^2} \mathfrak{t}k_1 + (n - 1)u(1 - u^2) \left(\frac{h'}{h} \right)' \\ &\quad - (n - 1)(1 - u^2) \left(u \frac{f'}{f} - k_1 \right) \frac{h'}{h}. \end{aligned} \quad (5.9)$$

Then, combining (5.8) with (5.9) and rearranging terms, we arrive at the formula

$$\begin{aligned} \frac{\partial u}{\partial t} &= \Delta u + \left(u \frac{f'}{f} - k_1 \right)^2 u + u(1 - u^2) \left(\frac{f''}{f} - 2 \left(\frac{f'}{f} \right)^2 + (n - 1) \left(\left(\frac{h'}{h} \right)' - \frac{h' f'}{h f} \right) \right) \\ &\quad + \overline{H} \frac{f'}{f} (1 - u^2) \end{aligned} \quad (5.10)$$

Now the definition of v leads to

$$dv = -\frac{1}{u^2}du, \quad \frac{\partial v}{\partial t} = -\frac{1}{u^2} \frac{\partial u}{\partial t}, \quad \Delta v = -\frac{1}{u^2} \Delta u + \frac{2}{u^3} |du|^2. \quad (5.11)$$

Finally, (5.1) follows from (5.10) and (5.11). \square

Theorem 5 *Let M_t be the solution of (1.1) defined on a maximal time interval $[0, T[$, with initial condition M in the setting RSS2 and satisfying (1.3). The generating curve of the solution M_t of (1.1) remains a graph over the axis of revolution for every $t \in [0, T[$.*

Proof The case of \overline{M} being the Euclidean space was proved in [5]. Hence we shall suppose that $S_{zi} < 0$ and $S_{ri} \leq 0$.

For every $t_0 \in [0, T[$, let $\rho(t_0) = \min_{(x,t) \in M \times [0, t_0]} r_t(x) > 0$. Then we can apply Corollary 3 to conclude that

$$0 < h_1 \leq \overline{H}_t \leq h_2 \quad \text{for } t \in [0, t_0], \quad (5.12)$$

where $h_i := h_i(V, \overline{g}, n, b - a, \text{vol}(M), \rho(t_0))$.

While the generating curve is a graph, one has $v \geq 1$ and so

$$\left(1 - \frac{1}{v^2}\right) v \leq v \quad \text{and} \quad \overline{H} \frac{f'}{f} \leq \overline{H} \frac{f'}{f} v. \quad (5.13)$$

Putting the above inequalities in (5.1) and forgetting about the negative addends, we reach the inequality:

$$\frac{\partial v}{\partial t} \leq \Delta v - \frac{2}{v} |dv|^2 - \left(\frac{1}{v} \frac{f'}{f} - k_1\right)^2 v - K v^2 + C v \leq \Delta v - K v^2 + C v, \quad (5.14)$$

where

$$K = K(V, \overline{g}, n, b - a, \text{vol}(M), \rho(t_0)) = h_1 \min_{r \in [\rho(t_0), r_2]} \frac{f'}{f}, \quad (5.15)$$

$$\begin{aligned} C &= C(V, \overline{g}, n, b - a, \text{vol}(M), \rho(t_0)) \\ &= \max_{r \in [\rho(t_0), r_2]} \left| \frac{f''}{f} - 2 \left(\frac{f'}{f}\right)^2 + (n-1) \left(\left(\frac{h'}{h}\right)' - \frac{f' h'}{f h} \right) \right| + h_2 \max_{r \in [\rho(t_0), r_2]} \frac{f'}{f}. \end{aligned} \quad (5.16)$$

Let $t_1 \in [0, t_0]$ be any time such that the generating curve of M_t is a graph for every $t \in [0, t_1]$. Let $(x_2, t_2) \in M_{t_2} \times [0, t_1]$ be the space-time point where v attains its maximum for $t \leq t_1$. If $v_2 := v(x_2, t_2) > 1$ and $t_2 \neq 0$, then x_2 is an interior point of M_{t_2} (because at the boundary $v = 1$) and it has to satisfy $\frac{\partial v}{\partial t} \geq 0$ as well as $\Delta v \leq 0$. Using these inequalities in (5.14), we can write $0 \leq -K v_2^2 + C v_2 = v_2(-K v_2 + C)$, and so

$$v_2 \leq C_1 = \max\{\max_{x \in M} v(x), C/K\}.$$

Notice that C_1 does not depend on t_1 , thus v cannot become ∞ in the interval $[0, t_0]$. Since t_0 is arbitrary, the same conclusion is true for $[0, T[$. \square

6 The first singularities of the motion are produced at the axis of revolution

This section is devoted to the proof of the following result, which assures long time existence unless the evolving hypersurface reaches the axis of rotation.

Theorem 6 *Let M_t be the maximal solution of (1.1), defined on $[0, T[$, with initial condition M in the setting RSS2 and satisfying (1.3). Then either*

$$T = \infty \quad \text{or} \quad \lim_{t \rightarrow T} \min_{x \in M_t} r_t(x) = 0.$$

Proof First, let us recall that $r_t > 0$ for every $t \in [0, T[$. Hence the continuous function $t \mapsto \beta(t) := \min_{x \in M_t} r_t(x)$ satisfies $\beta(t) > 0$ for every $t \in [0, T[$. If $\lim_{t \rightarrow T} \beta(t) > 0$, then there is a $\rho > 0$ such that $\beta(t) \geq \rho$ for every $t \in [0, T[$. We shall prove that, if this is the case, then the solution of the flow can be prolonged after T , which is a contradiction.

The key for proving that the solution can be extended if $r_t \geq \rho > 0$ is to show that, under this condition, $|L|^2$ is bounded. To do so, we need the evolution equation of $|L|^2$ under (1.1) for hypersurfaces in arbitrary Riemannian manifolds. Previously, it is necessary to know the evolution equations for the metric g_t of M_t , its dual metric g_t^\flat and the second fundamental form α_t :

$$\frac{\partial g}{\partial t} = 2(\bar{H} - H)\alpha, \quad \frac{\partial g^\flat}{\partial t} = -2(\bar{H} - H)\alpha^\flat,$$

$$\begin{aligned} \frac{\partial \alpha}{\partial t} = & \Delta \alpha + (\bar{H} - 2H)\alpha L - \bar{H}\bar{R} \cdot N \cdot N + (|L|^2 + \bar{Ric}(N, N))\alpha \\ & - \bar{Ric}(\cdot, L \cdot) - \bar{Ric}(L \cdot, \cdot) + \bar{R}(N, L \cdot, N, \cdot) + \bar{R}(N, \cdot, N, L \cdot) \\ & + 2RiL - \bar{\nabla} \cdot \bar{Ric}(\cdot, N) - \bar{\delta} \bar{R} \cdot N, \end{aligned}$$

where $RiL(Z, Y) = \sum_{i=1}^n \bar{R}_{e_i, Z, Le_i, Y}$ and $\bar{\delta} \bar{R} = \sum_{i=1}^n \bar{\nabla}_{e_i}(\bar{R})_{e_i}$ for some local orthonormal frame $\{e_i\}$ of M_t and αL is defined by $\alpha L(X, Y) = \alpha(LX, Y)$.

From these equations we obtain the evolution formula for the squared-norm of the shape operator $|L|^2$:

$$\begin{aligned} \frac{\partial |L|^2}{\partial t} = & \Delta |L|^2 - 2|\nabla L|^2 + 2|L|^2 (|L|^2 + \bar{Ric}(N, N)) \\ & - 2\bar{H} \left(\text{tr} L^3 + \sum_j \bar{R}_{Le_j Ne_j N} \right) - 4 \sum_{i,j} (\bar{R}_{Le_i e_j Le_i e_j} - \bar{R}_{Le_i e_j e_i Le_j}) \\ & - 2 \sum_{i,j} (\bar{\nabla}_{Le_i} \bar{R}_{Ne_j e_i e_j} + \bar{\nabla}_{e_j} \bar{R}_{Ne_i Le_i e_j}). \end{aligned} \quad (6.1)$$

In our case, using the orthonormal local frame $N, \mathbf{t} = E_1, E_2, \dots, E_n$, we get

$$\begin{aligned} \frac{\partial |L|^2}{\partial t} = & \Delta |L|^2 - 2|\nabla L|^2 + 2|L|^4 - 2\bar{H} \text{tr} L^3 + 2|L|^2 (\mathcal{T} + (n-1)\mathcal{J}) \\ & - 2\bar{H} (k_1 \mathcal{T} + (n-1)k_2 \mathcal{J}) - 4(k_1 - k_2)^2 (n-1) \mathcal{Y} - 2 \langle \alpha, \tilde{\delta} \bar{R}_N \rangle, \end{aligned} \quad (6.2)$$

where $\tilde{\delta}\bar{R}_N(X, Y) := \sum_i (\bar{\nabla}_X \bar{R}_{NE_i Y E_i} + \bar{\nabla}_{E_i} \bar{R}_{N Y X E_i})$, \mathcal{T} , \mathcal{J} and \mathcal{Y} are the sectional curvatures of the planes generated by $\{\mathfrak{t}, N\}$, $\{E_i, N\}$ and $\{E_i, \mathfrak{t}\}$, respectively.

We cannot deduce any bound for $|L|$ directly from (6.2). Then, as in [5], we shall study the evolution of $g = (\varphi \circ v)|L|^2$, where φ is a real function to be determined later. The evolution of g is given by those of φ and $|L|^2$ according to the formula:

$$\begin{aligned} \left(\frac{\partial}{\partial t} - \Delta\right) g &= |L|^2 \left(\frac{\partial}{\partial t} - \Delta\right) \varphi + \varphi \left(\frac{\partial}{\partial t} - \Delta\right) |L|^2 - 2 \langle d\varphi, d|L|^2 \rangle \\ &= |L|^2 \varphi' \left(\frac{\partial}{\partial t} - \Delta\right) v - |L|^2 \varphi'' |dv|^2 + \varphi \left(\frac{\partial}{\partial t} - \Delta\right) |L|^2 - 2 \langle d\varphi, d|L|^2 \rangle. \end{aligned} \quad (6.3)$$

For the last addend in (6.3), we use the following inequality (cf. [10])

$$-2 \langle d\varphi, d|L|^2 \rangle \leq -\frac{1}{\varphi} \langle dg, d\varphi \rangle + 2\varphi |d|L|^2| + \frac{3}{2\varphi} |L|^2 |d\varphi|^2 \quad (6.4)$$

and Kato's inequality

$$|d|L|^2| \leq |\nabla L|^2. \quad (6.5)$$

If $\varphi' > 0$, from (5.14) to (5.16) and (6.2) to (6.5), we conclude that there are constants $K = K(V, \bar{g}, n, b - a, \text{vol}(M), \rho)$ and $C = C(V, \bar{g}, n, b - a, \text{vol}(M), \rho)$ such that

$$\begin{aligned} \left(\frac{\partial}{\partial t} - \Delta\right) g &\leq |L|^2 \varphi' \left(-\frac{2}{v} |dv|^2 - \left(\frac{f'}{vf} - k_1 \right)^2 v - Kv^2 + Cv \right) - |L|^2 \varphi'' |dv|^2 \\ &\quad + 2\varphi |L|^4 - 2\varphi \bar{H} \text{tr} L^3 + 2\varphi |L|^2 (\mathcal{T} + (n-1)\mathcal{J}) \\ &\quad - 2\varphi \bar{H} (k_1 \mathcal{T} + (n-1)k_2 \mathcal{J}) - 4\varphi (k_1 - k_2)^2 (n-1) \mathcal{Y} \\ &\quad - 2 \langle \alpha, \tilde{\delta}\bar{R}_N \rangle \varphi - \frac{1}{\varphi} \langle dg, d\varphi \rangle + \frac{3}{2\varphi} |L|^2 |d\varphi|^2. \end{aligned} \quad (6.6)$$

Next, let us bound and/or rearrange the different addends in (6.6).

$$-\frac{2}{v} |L|^2 \varphi' |dv|^2 - |L|^2 \varphi'' |dv|^2 + \frac{3}{2\varphi} |L|^2 |d\varphi|^2 = \left(-\frac{2}{v\varphi'} - \frac{\varphi''}{(\varphi')^2} + \frac{3}{2\varphi} \right) |L|^2 |d\varphi|^2. \quad (6.7)$$

$$\begin{aligned} 2\varphi |L|^4 - |L|^2 \varphi' v \left(\frac{f'}{vf} - k_1 \right)^2 &= g \frac{\varphi'}{\varphi} v \left(2\frac{k_1 f'}{vf} - \frac{(f')^2}{v^2 f^2} - |L|^2 + (n-1)k_2^2 \right) + 2\varphi |L|^4 \\ &= \left(-\frac{\varphi' v}{\varphi^2} + \frac{2}{\varphi} \right) g^2 + (n-1)k_2^2 \frac{\varphi' v}{\varphi} g + 2k_1 \frac{f' \varphi'}{f \varphi} g - \left(\frac{f'}{f} \right)^2 \frac{\varphi'}{v \varphi} g. \end{aligned} \quad (6.8)$$

Using Young's inequality $xy \leq \varepsilon x^2 + \frac{1}{4\varepsilon} y^2$ with $x = g$, $y = \frac{2f' \varphi' |k_1|}{f \varphi}$ and $\varepsilon = k/2$, for some $k > 0$ that we shall determine later, we have

$$2k_1 \frac{f' \varphi'}{f \varphi} g \leq \frac{k}{2} g^2 + \frac{1}{2k} \left(\frac{2k_1 f' \varphi'}{f \varphi} \right)^2 \leq \frac{k}{2} g^2 + \frac{2}{k \varphi} \left(\frac{\varphi' f'}{\varphi f} \right)^2 g. \quad (6.9)$$

If φ is bounded from below and φ' is bounded from above by some positive constants, there is $0 < K_1 = C \sup \left(\frac{\varphi'}{\varphi} v \right)$ satisfying

$$|L|^2 \varphi' (-Kv^2 + Cv) \leq K_1 g. \quad (6.10)$$

Using $|\operatorname{tr} L^3| \leq |L|^3$ and Young's inequality with $x = |L|$, $y = 2\bar{H}$ and $\varepsilon = k\varphi$:

$$-2\varphi \bar{H} \operatorname{tr} L^3 \leq 2\varphi \bar{H} |L|^3 = 2\bar{H} |L| g \leq \left(k\varphi |L|^2 + \frac{1}{4k\varphi} 4\bar{H}^2 \right) g = kg^2 + \frac{\bar{H}^2}{k\varphi} g. \quad (6.11)$$

To bound the rest of the addends in (6.6), we recall that, since $\rho \leq r \leq r_2$, the curvatures \mathcal{T} , \mathcal{J} , \mathcal{Y} and $|\tilde{\delta} \bar{R}_N|$ are bounded; moreover, as

$$k_1 \leq |L|, \quad k_2 \leq |L| \quad \text{and} \quad (k_1 - k_2)^2 \leq 3|L|^2,$$

there are positive constants C_1, C_2, C_3 and C_4 such that

$$\begin{aligned} & 2\varphi |L|^2 (\mathcal{T} + (n-1)\mathcal{J}) + 2\varphi \bar{H} (k_1 \mathcal{T} + (n-1)k_2 \mathcal{J}) \\ & - 4\varphi (k_1 - k_2)^2 (n-1) \mathcal{Y} - 2 \langle \alpha, \tilde{\delta} \bar{R}_N \rangle \varphi \\ & \leq C_1 g + C_2 \sqrt{\varphi} \sqrt{g} + C_3 g + C_4 \sqrt{\varphi} \sqrt{g}. \end{aligned} \quad (6.12)$$

Now, in order that $\varphi' \circ v$ and $\varphi \circ v$ be double side bounded, and $\varphi > 1$ (conditions that we have used before), we choose the function φ and the constant k to be

$$\varphi(v) := \frac{v^2}{1 - kv^2}, \quad k := \frac{1}{2 \max v^2}. \quad (6.13)$$

With this election, φ also satisfies

$$\frac{3}{2\varphi} - \frac{\varphi''}{\varphi^2} - \frac{2}{\varphi'v} = \frac{v^2 - \varphi}{2\varphi^2} < 0 \quad \text{and} \quad \frac{2}{\varphi} - \frac{\varphi'}{\varphi^2} v + k + \frac{k}{2} = -\frac{k}{2}. \quad (6.14)$$

Thanks to the expressions from (6.6) to (6.14), we reach

$$\left(\frac{\partial}{\partial t} - \Delta \right) g \leq -\frac{k}{2} g^2 + K_2 g + K_3 \sqrt{g} - \frac{1}{\varphi} \langle dg, d\varphi \rangle, \quad (6.15)$$

for two positive constants K_2 and K_3 .

Let $t_1 \in [0, T[$. Let us suppose that (x_0, t_0) is the point where g attains its maximum for $t \leq t_1$. We further assume that x_0 is an interior point of M_{t_0} and $t_0 \neq 0$. Observe that this assumption can be made without loss of generality:

- i. If $z(x_0) = b$, then we can use the map $\iota_1 : (r, z, u) \mapsto (r, 2b - z, u)$, which is an isometry of \bar{M} , to extend the solution symmetrically onto $[a, 2b - a]$. In fact, by applying ι_1 to M_t , we construct a new solution $M_t \cup \iota_1(M_t)$ of (1.1), which is orthogonal to $z = a$ and $z = 2b - a$ at the boundary, such that x_0 still is a maximum of g , but it is not at the spatial boundary.

- ii. If $z(x_0) = a$, a similar argument with the isometry $\iota_2 : (r, z, u) \mapsto (r, 2a - z, u)$, allows us to regard x_0 as an interior point of M_{t_0} .

If $g_0 := g(x_0, t_0) > 1$, then it has to satisfy $0 \leq -\frac{k}{2}g_0^2 + (K_2 + K_3)g_0$, so

$$g_0 \leq \max\{\max_M g, 2(K_2 + K_3)/k, 1\}. \quad (6.16)$$

Since φ and g are bounded independently of T , $|L|^2$ is bounded on $[0, T[$ with a bound independent of t .

Once we achieve the upper bound for $|L|^2$, it follows, like in [13] and [14] (see [15], [5] or [7] for the volume-preserving version), that $|\nabla^j L|$ is bounded for every $j \geq 1$. If $T < \infty$, these bounds imply (cf. [13] pages 257, ff.) that X_t converges (as $t \rightarrow T$, in the C^∞ -topology) to a unique smooth limit X_T . Now we can apply the short time existence theorem to continue the solution after T , contradicting the maximality of $[0, T[$. \square

7 Finer bounds for \overline{H}

If M is a revolution surface with boundary in the setting RSS2, theorems 5 and 6 tell us that the flow M_t of M under (1.1), subject to the boundary conditions (1.3), exists as long as M_t does not touch the axis of revolution, and the generating curve of each M_t is a graph over the same axis. This fact will allow us to improve the bounds for \overline{H} obtained in section 4 by using a Sturmian theorem (cf. [4]), as it was done in [6] when \overline{M} is the Euclidean space.

Before going on, we want to stress a couple of remarks which is necessary to have in mind from now on.

Remark 4 *After the results of sections 4 and 5, it is clear that r and v remain bounded on $[0, t]$ for every $t \in [0, T[$. Then, from (3.3) and the definition of v , it follows that \dot{r} is also bounded. Now, from (3.5) and (3.7), we conclude that all the derivatives of r are bounded if and only if $|\nabla^j L|$ is bounded for every $j \geq 0$, which was established within the proof of Theorem 6.*

Remark 5 *Since the solutions of equations (1.1) and (3.8) differ only on the parametrization of M_t , the functions $r_t(z_t)$ and $r_t(z)$ corresponding to the solutions of (1.1) and (3.8), respectively, will have the same number of zeroes for their first derivatives (with respect to z) at the same values of z , and the same bounds for all their derivatives. Therefore, whenever we are interested on these magnitudes, we can use either equation (1.1) or (3.8) at our convenience.*

Lemma 7 *Let M_t , with $t \in [0, T[$, be a maximal solution of (1.1) with initial condition M in the setting RSS2 and satisfying (1.3). For each $t \in [0, T[$, the set $Z_t = \{z \in [a, b]; \dot{r}_t = 0\}$ is finite, the function $t \mapsto N(t) := \sharp(Z_t)$ is non increasing and, at the points (z_0, t_0) satisfying $0 = \dot{r}_{t_0}(z_0) = \ddot{r}_{t_0}(z_0)$ there is a neighborhood where the number of zeroes decreases.*

Proof Thanks to Remark 5, we can use equation (3.8) for this proof. In order to apply the Sturmian theorem (cf. Theorem 1.1 in [6]) to $\psi_t := \dot{r}_t$, we compute $\frac{\partial \psi}{\partial t}$ by taking the derivative of (3.9) with respect to z . The result is

$$\begin{aligned} \frac{\partial \psi}{\partial t} = & \frac{\ddot{\psi}}{\dot{r}^2 + f^2} + \left(\frac{\overline{H}(f\ddot{r} - f'\dot{r}^2)}{f^2\sqrt{\dot{r}^2 + f^2}} - \left(\frac{f'}{f} + (n-1)\frac{h'}{h} \right)' - 2\frac{f'\ddot{r}}{f(\dot{r}^2 + f^2)} - 2\ddot{r}\frac{\ddot{r} + ff'}{(\dot{r}^2 + f^2)^2} \right. \\ & \left. - \frac{\dot{r}^2}{\dot{r}^2 + f^2} \left(\frac{f'}{f} \right)' + 2\dot{r}^2 \frac{f'}{f} \frac{\ddot{r} + ff'}{(\dot{r}^2 + f^2)^2} \right) \psi. \end{aligned} \quad (7.1)$$

For every $t_1 < T$, let $\rho(t_1) = \min_{(x,t) \in M \times [0,t_1]} r_t(x) > 0$. As follows from Remark 4, the coefficients of (7.1) satisfy the hypotheses of the Sturmian theorem on $[0, t_1]$. Then, applying such theorem, we conclude that the statement is true for $t \in [0, t_1]$ and, since t_1 is arbitrary, also for all $t \in [0, T[$. \square

A first consequence, deduced from the previous lemma, is that we can find an upper bound for the length of the generating curve of the evolving hypersurface. Indeed,

Lemma 8 *Let M_t , with $t \in [0, T[$, be a maximal solution of (1.1) with initial condition M in the setting RSS2 and satisfying (1.3). There is a constant $c(N(0), r_2, \bar{g}, b - a)$ such that, for every $t \in [0, T[$, the length of the generating curve of M_t is lower than c .*

Proof From Lemma 7, $Z_t = \{a = z_1, \dots, b = z_{N(t)}\}$ and $N(t) \leq N(0)$. The computation of the length of the generating curve γ_t gives

$$L(\gamma_t) = \sum_{j=1}^{N(t)-1} \int_{z_j}^{z_{j+1}} \sqrt{\dot{r}(z)^2 + f(r(z))^2} dz \leq f(r_2)(b - a) + (N(0) - 1)r_2 =: c.$$

\square

Now we are in position to use the previous lemmas in order to achieve finer bounds for \overline{H} . With “finer” we mean that, unlike Corollary 3, here we obtain a double side bound for \overline{H} independent of the lower bound for r . In fact,

Proposition 9 *Let M_t , with $t \in [0, T[$, be a maximal solution of (1.1) with initial condition M in the setting RSS2 and satisfying (1.3). Then we can find constants $h_i(V, \bar{g}, n, b - a, \text{vol}(M), \mathcal{A}_V) > 0$, $i = 1, 2$ such that $h_1 \leq \overline{H} \leq h_2$, where \mathcal{A}_V is the infimum of the n -volume of the hypersurfaces satisfying the constraint of the setting RSS2 and enclosing a volume V .*

Proof As in the proof of Proposition 2 we have $I_1 \geq 0$ and, starting with formula (4.13) for I_1 , we can write:

$$I_1 \leq \sigma \frac{(h^{n-1})'(r_2)}{\text{vol}(M_t)} \frac{\pi}{2} L(\gamma_t) \leq \sigma \frac{(h^{n-1})'(r_2)}{\mathcal{A}_V} \frac{\pi}{2} c, \quad (7.2)$$

where the last inequality follows by Lemma 8. On the other hand, from the expression (4.15) for I_2 , we obtain

$$I_2 \leq \frac{\sigma}{\mathcal{A}_V}(b-a) \max_{r \in [0, r_2]} ((n-1)h'f + f'h)h^{n-2}(r). \quad (7.3)$$

To prove that I_2 has a positive uniform lower bound on $[0, T[$, let us consider

$$r_3 = \beta^{-1} \left(\frac{V}{2(b-a)\sigma} \right). \quad (7.4)$$

Observe that, by the above definition, $0 < r_3 < r_1 < r_2$. Let L_3 be the measure of the set $\mathcal{L}_3 = \{z \in [a, b]; r_t(z) \geq r_3\}$. From (7.4), (4.1) and (4.4), we have

$$\begin{aligned} \frac{V}{2\sigma} &= \int_a^b \int_0^{r_t(z)} f h^{n-1} dr dz - \int_a^b \int_0^{r_3} f h^{n-1} dr dz = \int_a^b \int_{r_3}^{r_t(z)} f h^{n-1} dr dz \\ &\leq \int_{\mathcal{L}_3} \int_{r_3}^{r_t(z)} f h^{n-1} dr dz \leq f h^{n-1}(r_2) L_3 (r_2 - r_3), \end{aligned} \quad (7.5)$$

which implies that $L_3 \geq \rho_3 > 0$ for some ρ_3 independent of t . Now, we can use again (4.15) to get

$$I_2 \geq \frac{\sigma}{\text{vol}(M_t)} \int_{\mathcal{L}_3} ((n-1)h'f + f'h) h^{n-2} dz \geq \frac{\sigma \rho_3}{\text{vol}(M)} \min_{r \in [r_3, r_2]} ((n-1)h'f + f'h)h^{n-2}(r).$$

This gives the desired positive lower bound for I_2 , which, together with the other inequalities obtained for I_1 and I_2 , finishes the proof of the proposition. \square

8 Discreteness of the first singularity set

According to Theorem 6, at the points z where one has a singularity at time T one must have $\lim_{t \rightarrow T} r_t(z) = 0$. It seems that these points must be limits of zeroes of \dot{r}_t . We shall show that this is the case, which, thanks to Lemma 7, gives the discreteness of the set of singular points at time T .

By Lemma 7, $N(t)$ is finite and does not increase with time. Moreover, if $0 = \dot{r}_{t_0}(z_0) = \ddot{r}_{t_0}(z_0)$, there is a neighborhood where the number of zeroes decreases. Then, we can find some time $t_1 \in [0, T[$ such that for every $t \in [t_1, T[$, $N(t)$ is constant and $0 \neq \ddot{r}_t(z)$ for every $z \in Z_t$. Thus all the critical points of r_t are local maxima or minima. Hereafter $\xi_i(t)$ denotes the minima and $\eta_i(t)$, the maxima.

Notice that, if the number of critical points for r_t on $[a, b]$ is $2k + 2$, we may have the following situations

$$\begin{aligned} a &< \xi_1(t) < \eta_1(t) < \dots < \xi_k(t) < \eta_k(t) < b \\ a &< \eta_1(t) < \xi_1(t) < \dots < \eta_k(t) < \xi_k(t) < b \end{aligned}$$

and, if the number of critical points for r_t is $2k + 3$,

$$\begin{aligned} a &< \xi_1(t) < \eta_1(t) < \dots < \xi_k(t) < \eta_k(t) < \xi_{k+1}(t) < b \\ a &< \eta_1(t) < \xi_1(t) < \dots < \eta_k(t) < \xi_k(t) < \eta_{k+1}(t) < b. \end{aligned}$$

On the other hand, as $\ddot{r}_t(\xi_i(t)) \neq 0 \neq \ddot{r}_t(\eta_i(t))$, by the implicit function theorem, the solutions $\xi_i(t)$, $\eta_i(t)$ of the equation $\dot{r}_t(z) = 0$ are smooth as functions of time. In addition, since the map $(r, z, u) \mapsto (r, 2z_0 - z, u)$ is an isometry of \overline{M} for every $z_0 \in \mathbb{R}$, the same arguments given in Lemma 2.3 of [6] (together with those of 5.1 in [1]) imply the existence of the limits

$$\xi_i(T) = \lim_{t \rightarrow T} \xi_i(t) \quad \text{and} \quad \eta_i(T) = \lim_{t \rightarrow T} \eta_i(t).$$

We can now state the main theorem of the present section.

Theorem 10 *Let M_t , with $t \in [0, T[$, be a maximal solution of (1.1) with initial condition M in the setting RSS2 and satisfying (1.3). For any closed interval $[c, d]$ not containing any of the points $a, \xi_i(T), \eta_i(T), b$, there exist $\delta > 0$ and $t_2 \in [0, T[$ such that $r_t(z) \geq \delta$ for $z \in [c, d]$, $t \in]t_2, T[$.*

Proof First, assume that $[c, d] \subset]\xi_j(T), \eta_j(T)[$ (the other cases can be proved in a similar way, in a half of them with the inequalities reversed and using upper right Dini derivative instead of the lower one). Let a', b' be real numbers satisfying $[c, d] \subset]a', b'[$ and $[a', b'] \subset]\xi_j(T), \eta_j(T)[$. By the definition of t_1 (at the beginning of this section), $\dot{r}_t(z) > 0$ for $\xi_j(t) < z < \eta_j(t)$ and $t \in [t_1, T[$. Moreover, by the continuity of $\xi_j(t)$ and $\eta_j(t)$ on $[t_1, T]$, there is a $t_2 \in]t_1, T[$ such that $\dot{r}_t(z) > 0$ for $(z, t) \in [a', b'] \times]t_2, T[$. On this domain, by discarding the positive addends in (7.1), we have

$$\frac{\partial \psi}{\partial t} \geq \frac{\ddot{\psi}}{\dot{r}^2 + f^2} + T(\psi) \psi - \frac{f' \overline{H}}{f^2 \sqrt{\dot{r}^2 + f^2}} \psi^3, \quad (8.1)$$

being

$$\begin{aligned} T(\psi) = & \frac{\overline{H} \ddot{r}}{f \sqrt{\dot{r}^2 + f^2}} - \frac{f''}{f} - 2 \frac{f' \ddot{r}}{f(\dot{r}^2 + f^2)} - 2 \frac{\dot{r}^2 + f f' \ddot{r}}{(\dot{r}^2 + f^2)^2} \\ & - (n-1) \frac{h''}{h} + \frac{\dot{r}^2}{f(\dot{r}^2 + f^2)} \left(-f'' + \frac{2 \ddot{r} f'}{\dot{r}^2 + f^2} \right). \end{aligned}$$

Next, taking A and M_1 two constants to be determined later, we define

$$w(z, t) = e^{-At} \sin(\mu(z - a')), \quad \text{where } \mu = \frac{\pi}{b' - a'} \quad \text{and} \quad A \geq \mu^2 + M_1. \quad (8.2)$$

Observe that w satisfies

$$\frac{\partial w}{\partial t} \leq -(\mu^2 + M_1) w \leq -\frac{\mu^2}{f^2 + \dot{r}^2} w - M_1 w = \frac{\ddot{w}}{f^2 + \dot{r}^2} - M_1 w. \quad (8.3)$$

Let us consider the difference $\mathfrak{w} = \psi - w$. From (8.1) and (8.3), we obtain the estimate:

$$\frac{\partial \mathfrak{w}}{\partial t} \geq \frac{\ddot{\mathfrak{w}}}{\dot{r}^2 + f^2} + \left(T(\psi) - \frac{f' \overline{H}}{f^2 \sqrt{\dot{r}^2 + f^2}} \right) \psi + M_1 w + \frac{f' \overline{H}}{f^2 \sqrt{\dot{r}^2 + f^2}} (\psi - \psi^3). \quad (8.4)$$

Let us define the function $\mathfrak{w}_m(t) := \min_{z \in [a', b']} \mathfrak{w}(z, t)$. For any $t' \in [t_2, T[$, let $I_{t'} = \{z \in [a', b']; \mathfrak{w}_m(t') = \mathfrak{w}(z, t')\}$. We have two possibilities for any fixed $\bar{t} \in [t_2, T[$: either $\{a', b'\} \cap I_{\bar{t}} \neq \emptyset$ or $\{a', b'\} \cap I_{\bar{t}} = \emptyset$.

$$\begin{aligned} \text{If } a' \text{ or } b' \text{ are in } I_{\bar{t}}, \text{ then} \quad & \mathfrak{w}(z, \bar{t}) \geq \mathfrak{w}_m(\bar{t}) = \psi_{\bar{t}}(a') > 0 \text{ or} \\ & \mathfrak{w}(z, \bar{t}) \geq \mathfrak{w}_m(\bar{t}) = \psi_{\bar{t}}(b') > 0. \end{aligned} \quad (8.5)$$

In the other case (that is, if $\{a', b'\} \cap I_{\bar{t}} = \emptyset$), let us call $]t_3, t_4[$ the maximal open interval in $[t_2, T]$ such that $\bar{t} \in]t_3, t_4[$ and $\{a', b'\} \cap I_t = \emptyset$ for every $t \in]t_3, t_4[$ (notice that $t_3 = t_2$ or $\{a', b'\} \cap I_{t_3} \neq \emptyset$ and that $t_4 = T$ or $\{a', b'\} \cap I_{t_4} \neq \emptyset$).

On this interval we can write

$$\frac{\partial \mathfrak{w}_m}{\partial t}(t) \geq \min_{z \in I_t} \left(\left(T(\psi) - \frac{f' \overline{H}}{f^2 \sqrt{\dot{r}^2 + f^2}} \right) \psi + M_1 w \right) + \min_{z \in I_t} \frac{f' \overline{H}}{f^2 \sqrt{\dot{r}^2 + f^2}} (\psi - \psi^3), \quad (8.6)$$

where $\frac{\partial \mathfrak{w}_m}{\partial t}(t)$ must be understood in the sense of a lower right Dini derivative (cf. page 160 of [8]).

Using the function $\Psi(t) = \min_{z \in I_t} (\psi - \psi^3)(z, t)$ (defined on $[t_3, t_4[$), we decompose

$$[t_3, t_4[= J_1 \cup J_2, \text{ where } J_1 = \Psi^{-1}(]0, \infty[) \text{ and } J_2 = \Psi^{-1}(]-\infty, 0]).$$

For $t \in J_2$, there is a $z_t \in I_t$ such that $\psi(z_t, t) \geq 1$. Hence, for every $z \in [a', b']$,

$$\mathfrak{w}(z, t) \geq \mathfrak{w}_m(t) = \psi(z_t, t) - w(z_t, t) \geq 1 - w(z_t, t) > 0. \quad (8.7)$$

If $t_0 \in J_1$, let $(\tau, \eta[$ be the maximal interval containing t_0 and contained in J_1 (this interval is $] \tau, \eta[$ if $\tau > t_3$ and it may be $[t_3, \eta[$ when $\tau = t_3$). Notice that $\Psi(t) > 0$ on $] \tau, \eta[$ and the function $t \mapsto \min_{z \in I_t} \frac{f' \overline{H}}{f^2 \sqrt{\dot{r}^2 + f^2}} (\psi - \psi^3)(z, t)$ is positive at the same t as $\Psi(t)$.

Therefore, by (8.6), we have

$$\frac{\partial \mathfrak{w}_m}{\partial t}(t) \geq \min_{z \in I_t} \left(\left(T(\psi) - \frac{f' \overline{H}}{f^2 \sqrt{\dot{r}^2 + f^2}} \right) \psi + M_1 w \right) (z, t) \quad \text{on }] \tau, \eta[. \quad (8.8)$$

On the other hand, at $z \in I_t$,

$$0 = \dot{\mathfrak{w}} = \dot{\psi} - \dot{w}, \quad \text{that is,} \quad \dot{\psi} = \dot{w} = e^{-At} \mu \cos(\mu(z - a')).$$

Thus $|\dot{\psi}| = |\dot{r}| \leq \mu$ on I_t , which can be used to bound

$$\begin{aligned} & T(\psi) - \frac{f' \overline{H}}{f^2 \sqrt{\dot{r}^2 + f^2}} \\ & \geq - \sup_{r \in [0, r_2]} \left(h_2(\mu + f') + 2|f''| + 4\mu f' + 2\mu^2 + 2ff'\mu + (n-1)\frac{h''}{h} \right) = -C_5, \end{aligned} \quad (8.9)$$

where $C_5 < \infty$ because $-\frac{h''}{h}$ is a sectional curvature of the ambient space (cf. (2.7)) and f has all its derivatives bounded on a closed interval. If $\bar{z} \in I_t$ is the point where the minimum in the right hand side of (8.8) is attained, then taking $M_1 = C_5$ and substituting (8.9) in (8.8), we get

$$\frac{\partial \mathfrak{w}_m}{\partial t}(t) \geq -C_5 (\psi - w)(\bar{z}, t) = -C_5 \mathfrak{w}_m(t) \quad \text{on }]\tau, \eta[. \quad (8.10)$$

From (8.10) we conclude that, on $]\tau, \eta[$,

$$\mathfrak{w}_m(t) \geq e^{-C_5(t-\tau)} \mathfrak{w}_m(\tau) \geq e^{-C_5(T-t_2)} \mathfrak{w}_m(\tau). \quad (8.11)$$

For every $z \in I_\tau$,

$$\mathfrak{w}_m(\tau) = \psi(z, \tau) - e^{-A\tau} \sin(\mu(z - a')) \geq \psi(z, \tau) - e^{-At_2}. \quad (8.12)$$

We have two possibilities for τ . Either $\tau = t_3$ or $\Psi(\tau) = 0$.

In the first case ($\tau = t_3$), as remarked after (8.5), we have again two possibilities:

- $\{a', b'\} \cap I_{t_3} \neq \emptyset$, then $\mathfrak{w}_m(\tau) = \mathfrak{w}_m(t_3) = \begin{cases} \psi(a', t_3) > 0 \\ \text{or} \\ \psi(b', t_3) > 0 \end{cases}$.
- $t_3 = t_2$, then $\Psi(t_2) \geq 0$ and, for $z \in I_{t_2}$,

$$m_1 := \min_{[a', b']} \psi(\cdot, t_2) \leq \psi(z, t_2) \leq 1.$$

Then, from (8.12), $\mathfrak{w}_m(\tau) > 0$ if $A \geq -\frac{1}{t_2} \ln m_1$, where $m_1 = \min_{z \in [a', b']} \dot{r}(z, t_2)$.

In the second case, there is a $z_\tau \in I_\tau$ such that $1 = \psi(z_\tau, \tau)$, thus $\mathfrak{w}_m(\tau) > 0$ if $A > 0$.

Since A has also to satisfy (8.2), we define $A = \max\{-\frac{1}{t_2} \ln m_1, C_5 + \mu^2\}$. With this election, having also in mind the inequalities (8.5), (8.7) and (8.11), we obtain $\mathfrak{w} \geq 0$ on $[a', b'] \times [t_2, T[$, i.e. $\dot{r}(z, t) \geq e^{-At} \sin(\mu(z - a'))$ and, integrating with respect to z between a' and z , for $z \in [c, d]$,

$$r(z, t) \geq \frac{e^{-At}}{\mu} \left(1 - \cos \left(\frac{\pi(z - a')}{b' - a'} \right) \right) \geq \frac{e^{-AT}}{\mu} \left(1 - \cos \left(\frac{\pi(c - a')}{b' - a'} \right) \right) = \delta > 0,$$

which finishes the proof of the theorem. \square

9 Small n-volume for M implies long time existence and convergence.

Lemma 11 *In the setting RSS, if $\text{vol}(M) \leq \frac{V}{(b-a)} \frac{\delta(r_1)}{\beta(r_1)}$, then there is a $r_0 > 0$ such that $r_t > r_0$ for every t such that the flow M_t exists.*

Proof If there is some t_0 and x_0 for which $r_{t_0}(x_0) = 0$, using (4.7) and (4.5), we obtain

$$\text{vol}(M) \geq \text{vol}(M_{t_0}) > \sigma \int_0^{r_1} h(r)^{n-1} dr \geq \frac{V\delta(r_1)}{(b-a)\beta(r_1)}, \quad (9.1)$$

in contradiction with the hypothesis. \square

Remark 6 Notice that $\frac{\delta(r_1)}{\beta(r_1)} \leq 1$ with equality if $f(r) \equiv 1$. So the hypothesis on the n -volume of M when \overline{M} is the Euclidean space coincides with that of [5]. For $n = 2$ and \overline{M} a space of constant curvature λ , the explicit value of the upper bound in Lemma 11 is $\frac{2\pi}{-\lambda} \left(-1 + \sqrt{1 - \frac{\lambda V}{\pi(b-a)}} \right)$.

Theorem 12 If M belongs to the setting RSS2 and has $\text{vol}(M) \leq \frac{V}{(b-a)} \frac{\delta(r_1)}{\beta(r_1)}$, then the solution of (1.1) satisfying (1.3) is defined for all $t > 0$ and converges to a revolution hypersurface of constant mean curvature in \overline{M} .

Proof By Lemma 11 and Theorem 6, we know that the solution of (1.1) is defined on $[0, \infty[$. Using again Lemma 11, together with Proposition 1, we conclude that r is double side bounded uniformly on $[0, \infty[$. After the results of section 5, it is clear that \dot{r} remains bounded all the time. In addition, the proof of Theorem 6 shows that $|\nabla^j L|^2$ is uniformly bounded for every $j \geq 0$. Once we have all these bounds, it follows from (3.5) and (3.7) that all the derivatives of r are bounded on $[0, \infty[$.

We are now in position to apply Arzelà-Ascoli Theorem to assure the existence of a sequence of maps r_{t_i} satisfying (3.9) which C^∞ -converges to a smooth map $r_\infty : [a, b] \rightarrow \mathbb{R}^+$ satisfying (3.9). A standard argument like in [7] proves that $M_\infty = (z, r_\infty(z), u)$ has constant mean curvature.

On the other hand, for any $0 < \tau_1 < \tau_2$, from (3.9), we have

$$\begin{aligned} \max_M |r(x, \tau_2) - r(x, \tau_1)| &\leq \max_M \int_{\tau_1}^{\tau_2} \left| \frac{\partial r}{\partial t}(x, t) \right| dt \\ &\leq \max_M \int_{\tau_1}^{\tau_2} |\overline{H} - H| \sqrt{\dot{r}^2 + f^2} dt \leq C\varepsilon(t) (\tau_2 - \tau_1), \end{aligned}$$

where $\varepsilon(t) \rightarrow 0$ as $t \rightarrow \infty$. Then, arguing as in [16] (end of section 2), we conclude that r_t C^∞ -converges to r_∞ . Therefore, $M_\infty = (z, r_\infty(z), u)$ is, up to tangential diffeomorphisms, the limit of M_t . \square

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