# Optical Differential Reflectance Spectroscopy of Ultrathin Epitaxial Organic Films

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WS10-ETOLDs, Valencia

June 2, 2010



www.www.organics.uni-jena.de

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- 2. Differential Reflectance Spectroscopy (DRS)
- 3. Case I: From Single Molecule to Crystal
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- 6. Summary



## **Motivation**

Advantages of *in situ* optical spectroscopy on molecular films:

- Organic molecules have large absorption coefficient:  $\alpha \approx 10^7$  cm<sup>-1</sup> (cmp.: Si:  $\alpha \approx 10^3$  cm<sup>-1</sup>)
- Extremely sensitive, down to ~0.1 ML of aromatic molecules

 $\Rightarrow$  Substrate-film-interface can be probed

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  - ⇒ Substrate-film-interface can be probed
- Optical properties depend sensitively on molecular arrangement:



- Lateral vs. vertical packing (anisotropy)

$$\Delta E = \frac{2|\vec{\mu}|^2}{r^3} \left( \cos(\alpha) - 3\cos^2(\vartheta) \right)$$

 $\Delta E$  = exciton splitting energy

- Effects of physical "chain length" (exciton delocalization and confinement)
- Charge transfer (molecular ions)





[M. Kasha et al., Pure Appl. Chem. 11, 371 (1965).]

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- Effects of physical "chain length" (exciton delocalization and confinement)
- Charge transfer (molecular ions)
- $\Rightarrow$  Structural information can be deduced
- Non-destructive



[M. Kasha et al., Pure Appl. Chem. <u>11</u>, 371 (1965).]

$$\mathsf{DRS}(E,d) = \frac{R(E,d) - R(E,0)}{R(E,0)}$$

d = film thickness (sub-ML to several MLs)



#### Growth of molecular films:

- Base pressure ≈ mid 10<sup>-10</sup> mbar
- Deposition in UHV at ≈ 300 .. 500 ℃
- Rate: 0.1 .. 1 ML/min
- Substrates at room temperature

[H. Proehl et al., *Phys. Rev. B* <u>71</u>, 165207 (2005).] [R. Forker et al., *PCCP* <u>11</u>, 2142 (2009).]

Optical properties of molecules shall be described by energy-dependent dielectric function:

$$\hat{\varepsilon}_{\text{film}}(E) = \varepsilon'_{\text{film}}(E) - i \cdot \varepsilon''_{\text{film}}(E)$$





[A. B. Djurišić et al., Opt. Commun. 183, 123 (2000).]

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#### How are DRS and dielectric function related???



[A. B. Djurišić et al., Opt. Commun. 183, 123 (2000).]

$$\mathsf{DRS}(E,d) = \frac{R(E,d) - R(E,0)}{R(E,0)}$$

McIntyre-approximation for ultrathin films ( $d \ll \lambda$ ):

$$\mathsf{DRS} \approx -\frac{8\pi d}{\lambda} \cdot \left[ \mathbf{A} \cdot \boldsymbol{\varepsilon}_{\textit{film}}'' + \mathbf{B} \cdot \left( \boldsymbol{\varepsilon}_{\textit{film}}' - 1 \right) \right] \quad \Rightarrow \quad \boldsymbol{\varepsilon}_{\textit{film}}'' \approx -\frac{\lambda}{8\pi d} \cdot \frac{\mathsf{DRS}}{\mathsf{A}} - \frac{\mathsf{B}}{\mathsf{A}} \cdot \left( \boldsymbol{\varepsilon}_{\textit{film}}' - 1 \right)$$

- A and B depend solely on substrate's properties



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[J.D.E. McIntyre et al., Surf. Sci. 24, 417 (1971).]

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 $\rightarrow$ 

- A and B depend solely on substrate's properties
- on **transparent substrates** and some metals (low *E*) : *B* ≈ 0

$$\varepsilon_{\rm film}'' \approx -\frac{\lambda}{8\pi d} \cdot \frac{\rm DRS}{\rm A}$$



$$\mathsf{DRS}(E,d) = \frac{R(E,d) - R(E,0)}{R(E,0)}$$

Approximation for ultrathin films ( $d \ll \lambda$ ):

$$\mathsf{DRS} \approx -\frac{8\pi d}{\lambda} \cdot \left[ \mathbf{A} \cdot \boldsymbol{\varepsilon}_{\textit{film}}'' + \mathbf{B} \cdot \left( \boldsymbol{\varepsilon}_{\textit{film}}' - 1 \right) \right] \quad \Rightarrow \quad \boldsymbol{\varepsilon}_{\textit{film}}'' \approx -\frac{\lambda}{8\pi d} \cdot \frac{\mathsf{DRS}}{\mathsf{A}} - \frac{\mathsf{B}}{\mathsf{A}} \cdot \left( \boldsymbol{\varepsilon}_{\textit{film}}' - 1 \right)$$

- A and B depend solely on substrate's properties
- on **transparent substrates** and some metals (low *E*) : *B* ≈ 0 →

$$\varepsilon_{\text{film}}'' \approx -\frac{\lambda}{8\pi d} \cdot \frac{\text{DRS}}{A}$$

- on **non-transparent substrates**: numerical calculation required! →

DRS 
$$\xrightarrow{\text{numerical}} \varepsilon'_{\text{film}}$$
 and  $\varepsilon''_{\text{film}}$ 

- model-free Kramers-Kronig consistent algorithm
- no approximations made in thin film equations
- generally valid for all kinds of substrates





**Further Details: Review-Article** 

**Roman Forker und Torsten Fritz,** Phys. Chem. Chem. Phys. 11, 2142-2155 (2009).

![](_page_13_Picture_4.jpeg)

spectroscopy of ultrathin epitaxial organic films

electrocatalysis in polymeric iron oxyhydroxide films

![](_page_13_Picture_7.jpeg)

## 3. Case I: From Single Molecule to Crystal

## **PTCDA on Mica**

+

![](_page_14_Picture_2.jpeg)

![](_page_14_Picture_3.jpeg)

![](_page_14_Picture_4.jpeg)

## What is the Problem to be Solved?

![](_page_15_Figure_1.jpeg)

![](_page_15_Picture_2.jpeg)

## What is the Problem to be Solved?

![](_page_16_Figure_1.jpeg)

![](_page_16_Picture_2.jpeg)

## **PTCDA is a quasi-1-dimensional Material**

Why would one expect thickness depending properties?

**Crystal Structure of PTCDA:** 

![](_page_17_Figure_3.jpeg)

Two different dimers exist: A and B, but interaction in B very strong and in A very weak

- $\Rightarrow$  3D-crystals behave like 1D-crystals, and:
- ⇒ Monomer should become visible for  $d \rightarrow 1$  ML !

![](_page_17_Picture_7.jpeg)

![](_page_18_Figure_2.jpeg)

![](_page_18_Picture_3.jpeg)

[H. Pröhl, R. Nitsche, T. Dienel and TF, PRL <u>93</u>, 097403 (2004).]

![](_page_19_Figure_2.jpeg)

![](_page_19_Picture_3.jpeg)

[H. Pröhl, R. Nitsche, T. Dienel and TF, PRL <u>93</u>, 097403 (2004).]

![](_page_20_Figure_2.jpeg)

![](_page_20_Picture_3.jpeg)

[H. Pröhl, R. Nitsche, T. Dienel and TF, PRL <u>93</u>, 097403 (2004).]

![](_page_21_Figure_2.jpeg)

![](_page_21_Picture_3.jpeg)

[H. Pröhl, R. Nitsche, T. Dienel and TF, PRL <u>93</u>, 097403 (2004).]

![](_page_22_Figure_2.jpeg)

![](_page_22_Picture_3.jpeg)

[H. Pröhl, R. Nitsche, T. Dienel and TF, PRL <u>93</u>, 097403 (2004).]

![](_page_23_Figure_2.jpeg)

![](_page_23_Picture_3.jpeg)

[H. Pröhl, R. Nitsche, T. Dienel and TF, PRL <u>93</u>, 097403 (2004).]

![](_page_24_Figure_2.jpeg)

![](_page_24_Picture_3.jpeg)

[H. Pröhl, R. Nitsche, T. Dienel and TF, PRL <u>93</u>, 097403 (2004).]

![](_page_25_Figure_2.jpeg)

![](_page_25_Picture_3.jpeg)

[H. Pröhl, R. Nitsche, T. Dienel and TF, PRL <u>93</u>, 097403 (2004).]

## **Monomer-Dimer-Transition**

![](_page_26_Figure_2.jpeg)

[H. Pröhl, R. Nitsche, T. Dienel and TF, PRL <u>93</u>, 097403 (2004).]

## **Monomer-Oligomer-Transition**

![](_page_27_Figure_2.jpeg)

![](_page_27_Picture_3.jpeg)

[H. Pröhl, R. Nitsche, T. Dienel and TF, PRL <u>93</u>, 097403 (2004).]

## 4. Case II: Function Follows Form

## PTCDA on KCI(100)

+

![](_page_28_Picture_2.jpeg)

![](_page_28_Picture_3.jpeg)

![](_page_28_Picture_4.jpeg)

![](_page_29_Figure_1.jpeg)

• Monomer spectra for *d* < 1ML

![](_page_29_Picture_3.jpeg)

[T. Dienel, C. Loppacher, S. Mannsfeld, R. Forker and TF, Adv. Mat. 20, 959 (2008).]

## DRS

## **PTCDA**

## **Absorption**

![](_page_30_Figure_2.jpeg)

- Monomer spectra for d < 1ML</li>
- Extremely narrow bands in comparison to solution and on mica

# Strongly reduced inhomogeneous broadening Hints towards special layer structure

![](_page_30_Picture_6.jpeg)

**AFM** 

## Low temperature NC-AFM (Chr. Loppacher):

![](_page_31_Figure_3.jpeg)

PTCDA on KCI(100)

Model

- Commensurate structure causes similar environment on atomic level for every molecule
- Reason why inhomogeneous broadening is strongly reduced

## Increasing coverage to more than 1 ML renders structure instable:

![](_page_32_Figure_3.jpeg)

- Spectra become very broad
- Further ageing leads to spectrum similar to pc 3D-crystals
  - ➔ Spectroscopy tells us: recrystallization!

![](_page_32_Picture_7.jpeg)

**AFM** 

## Low temperature NC-AFM (Chr. Loppacher):

![](_page_33_Picture_3.jpeg)

2.5 ML PTCDA on KCI(100)

- As the substrate's influence is already lost in the 2<sup>nd</sup> layer, layer has recrystallized into bulk structure
- Optical properties depend strongly on physical layer structure

**Potential Energy Calculation** 

Recrystallization can be fully understood by potential energy calculations:

Calculated potential energy as function of layer number for PTCDA on KCI(100) Values are given per molecule						
Phase	Layer number	E <sub>inter</sub> (kcal/mol)	E <sub>intra</sub> (kcal/mol)	E <sub>total</sub> (kcal/mol)		
SQ	1.	-37.17	-4.62	-41.79		
SQ	2.	-32.30	-4.62	-36.92		
HB	1.	-22.74	-11.74	-34.48		
HB	2.	-38.70	-11.74	-50.44		
HB	3.	-40.26	-11.74	-52.00		

![](_page_34_Picture_4.jpeg)

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- 1<sup>st</sup> layer: square-phase (SQ) energetically superior to herringbonephase (HB)

![](_page_35_Picture_5.jpeg)

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Structures in organic-inorganic heteroepitaxie are determined by a delicate balance between substratemolecule-interaction and molecule-molecule-interaction

## 5. Case III: Charged Molecules on Metals

## PTCDA on Au(111) & pc-Al

+

![](_page_37_Picture_2.jpeg)

![](_page_37_Picture_3.jpeg)

![](_page_37_Picture_4.jpeg)

![](_page_38_Figure_1.jpeg)

![](_page_38_Figure_2.jpeg)

**Distinct spectral development observed:** 

- 1<sup>st</sup> ML: Very broad spectra → coupling to Au!
- 2<sup>nd</sup> ML: Two peaks resolvable
- **3<sup>rd</sup> ML** and thicker: Double-feature diminishes

[R. Forker, C. Golnik, G. Pizzi, T. Dienel and TF, Org. Electr. <u>10</u>, 1448 (2009).]

## DRS

![](_page_39_Figure_2.jpeg)

![](_page_40_Figure_1.jpeg)

![](_page_40_Picture_2.jpeg)

[R. Forker, C. Golnik, G. Pizzi, T. Dienel and TF, Org. Electr. <u>10</u>, 1448 (2009).]

![](_page_41_Figure_1.jpeg)

 $\rightarrow$  Clear monomeric signature for E > 2.2 eV (neutral PTCDA)

- → No physical dimerization with the 1<sup>st</sup> ML of PTCDA !!!
- $\rightarrow$  Decoupling of 2<sup>nd</sup> ML is justification of 2-layer-approach

![](_page_41_Picture_5.jpeg)

[R. Forker, C. Golnik, G. Pizzi, T. Dienel and TF, Org. Electr. 10, 1448 (2009).]

![](_page_42_Figure_1.jpeg)

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- → No physical dimerization with the 1<sup>st</sup> ML of PTCDA !!!
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#### $\rightarrow$ New feature @ 2.05 eV which is <u>no</u> aggregation effect

→ Assignment to PTCDA radical cations (PTCDA•+)

[R. Forker, C. Golnik, G. Pizzi, T. Dienel and TF, Org. Electr. <u>10</u>, 1448 (2009).]

![](_page_42_Picture_8.jpeg)

## cmp. to solvent spectra

![](_page_43_Figure_2.jpeg)

[R. Forker, C. Golnik, G. Pizzi, T. Dienel and TF, Org. Electr. 10, 1448 (2009).]

![](_page_43_Picture_4.jpeg)

## cmp. to solvent spectra

![](_page_44_Figure_2.jpeg)

<sup>[</sup>R. Forker, C. Golnik, G. Pizzi, T. Dienel and TF, Org. Electr. 10, 1448 (2009).]

## cmp. to solvent spectra

![](_page_45_Figure_2.jpeg)

[R. Forker, C. Golnik, G. Pizzi, T. Dienel and TF, Org. Electr. 10, 1448 (2009).]

## PTCDA on *pc*-Al

## cmp. to solvent spectra

![](_page_46_Figure_2.jpeg)

# 6. Summary

- By DRS the optical properties of ultrathin films can be measured reliably and with an excellent SNR
- DRS reflects directly changes in the physical structure of ultra thin films
- On inert substrates (i.e. mica): molecules show monomeric behavior
- On metals: first layer hybridizes with the metal; no distinct spectral features in 1<sup>st</sup> ML
- Depending on the sign of the surface dipole, charging may be observed

![](_page_47_Picture_6.jpeg)

## **Acknowledgment**

#### **Co-workers:**

Dr. Roman Forker Dr. Stefan Mannsfeld Dr. Holger Pröhl Dr. Thomas Dienel Dr. Robert Nitsche Dr. Michael Törker

Christian Wagner Rainer Jacob Andreas Krause Giovanni Pizzi

![](_page_48_Picture_4.jpeg)

TU Dresden, IAPP (head: Prof. Dr. Karl Leo)

**\$\$\$**:

DFG, DAAD, Leibniz-Price Karl Leo

![](_page_48_Picture_8.jpeg)

# THE END

![](_page_49_Picture_1.jpeg)

![](_page_50_Picture_0.jpeg)

![](_page_51_Picture_0.jpeg)